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Two new computer-aided testing techniques based on the impulse frequency response approach were developed. Small beam specimens were excited in either flexural or extensional vibration by an electromagnetic hammer with a force transducer, while specimen response was measured with an eddy current displacement transducer or an accelerometer. A desktop computer/Fast Fourier Transform analyzer system was then used for rapid data acquisition and computation of the complex modulus by curve-fitting to the frequency response function. Specimens of graphite/epoxy, boron/epoxy and Kevlar aramid/epoxy were fabricated by using an autoclave-style press cure which was developed specifically for this program.

Although a number of parameters were studied, the emphasis was on the influence of fiber length, fiber orientation and fiber material on damping of polymer composites. The experimental results show that, as predicted, very low fiber aspect ratios are required to produce significant improvements in damping. Of the three fiber types tested, the Kevlar aramid fiber composite was found to have much better damping than graphite or boron fiber composites. Measurements and predictions also indicate that the control of fiber orientation in a continuous fiber reinforced laminate may be a better approach to the improvement of damping than the control of fiber aspect ratio.

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# IMPROVEMENT AND OPTIMIZATION OF INTERNAL DAMPING IN FIBER REINFORCED COMPOSITE MATERIALS

by

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#### ABSTRACT

This report describes two new techniques for characterization of damping in fiber reinforced composite materials and the application of these techniques to the study of improvement and optimization of damping. Flexural and extensional vibration tests are used for determination of complex moduli of aligned discontinuous fiber composite and off-axis fiber composite specimens, respectively. The results from tests on graphite/epoxy, Kevlar/epoxy and boron/epoxy composites are then compared with theoretical predictions from micromechanics models based on a single fiber. The analytical model is fitted to the experimental results by varying certain parameters that have uncertainties associated with them. It is shown that improved damping can be obtained with very low fiber aspect ratios, and that even better damping properties are possible with off-axis fibers.

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#### LIST OF SYMBOLS

a	- geometrical dimension defined in Fig. 2.9
a(x <sub>o</sub> )	- displacement amplitude at location of probe
a <sub>i</sub>	- amplitude of the input
a <sub>r</sub>	- amplitude of the response
ä <sub>b</sub>	- base acceleration amplitude of beam
A	- cross sectional area of the beam
b	- geometrical dimension defined in Fig. 2.9
c <sub>f</sub>	- final crucible weight
c	- initial crucible weight
$c_n$	- constant depending on resonant mode number n
đ	- diameter of the fiber
D	- energy dissipated per cycle
E*	- complex moduli
E', E"	- storage and loss moduli
E*	- complex extensional modulus for the fiber
EtT	- complex transverse modulus for the fiber
E*L	- composite extensional complex modulus
E*m	- complex modulus for the matrix
E* <sub>T</sub>	- composite transverse complex modulus
E*x	- composite complex modulus along the x axis
Ef', Ef"	- extensional storage and loss moduli for the fiber
E'fT' E"T	- transverse storage and loss moduli for the fiber

E', E'	<ul><li>extensional</li></ul>	storage	and loss	moduli	for	the
	composite					

 $E_m'$ ,  $E_m''$  - storage and loss moduli for the matrix

 $\mathbf{E_{T}}'$ ,  $\mathbf{E_{T}}''$  - transverse storage and loss moduli for the composite

 $\mathbf{E}_{\mathbf{x}}$ ',  $\mathbf{E}_{\mathbf{x}}$ " - composite storage and loss moduli along the x axis

f - frequency

f - frequency above resonance where the real part of the compliance reaches a peak (Fig. 2.6)

f frequency below resonance where the real part of the compliance reaches a peak of opposite sign (Fig. 2.6)

 $f_n$  - resonant frequency of mode n

f<sub>1</sub>, f<sub>2</sub> - frequencies at each end of the diameter parallel
to the real axes in the Nyquist plot

G<sub>f</sub> - shear modulus for the isotropic fiber

 ${f G}_{{f fLT}}$  - in-plane shear modulus for the anisotropic fiber

 ${f G}_{{f LT}}$  - in-plane shear modulus for the composite

 $G_{\mathrm{fLT}}^{\star}$  - complex in-plane shear modulus for the anisotropic fiber

 $\mathsf{G}_{\mathrm{LT}}^{\, lacktrell}$  - complex in-plane shear modulus for the composite

 $G_{m}^{*}$  - complex shear modulus for the matrix

 $G_f$ ',  $G_f$ " - real and imaginary components of the shear modulus for the isotropic fiber

 $G_{\mathrm{fLT}}^{"}$   $G_{\mathrm{fLT}}^{"}$  - real and imaginary components of the shear modulus for the anisotropic fiber

 ${\sf G_{LT}}^{\sf T}$  ,  ${\sf G_{LT}}^{\sf T}$  - real and imaginary components of the in-plane shear modulus for the composite

 $G_{\overline{m}}$ ,  $G_{\overline{m}}$  - real and imaginary components of the shear modulus for the matrix

G<sub>v</sub>(f) - power spectrum of the excitation

 $G_{xy}$  (f) - cross power spectrum of the excitation

 $G_{V}^{}(f)$  - power spectrum of the response

I - area moment of inertia of the beam about the neutral axes

K - bulk modulus

K<sub>f</sub> - fiber bulk modulus

 $K_{m}$  - matrix bulk modulus

l - fiber length

(l/d) - fiber aspect ratio

 $(\chi/d)_{eff}$  - effective fiber aspect ratio

L - subscript denoting longitudinal direction along the fiber

M - corresponding composite modulus:  $E_T$  or  $G_{L,T}$ 

 $\mathtt{M}_{\mathrm{f}}$  - corresponding fiber modulus  $\mathtt{E}_{\mathrm{f}}$  or  $\mathtt{G}_{\mathrm{f}}$ 

 $\mathbf{M}_{\!m}$  — corresponding matrix modulus  $\mathbf{E}_{\!m}$  or  $\mathbf{G}_{\!m}$ 

M<sub>1</sub>, M<sub>2</sub> - masses at each end of the free-free beam

n - mode number (1, 2, 3, ...)

Q - amplification factor

r - radial distance from center of the fiber

r - fiber radius

R - radius of representative volume element, or one

half of center-to-center fiber spacing

s - Laplace plane

span of the beam

t - time

- subscript denoting transverse direction, perpen-

dicular to the fiber

u - displacement of the beam in the x direction

U - stored energy per unit volume

 $v_c$  - composite volume fraction

 $v_f$  - fiber volume fraction

v<sub>m</sub> - matrix volume fraction

 $v_v$  - void volume fraction

w<sub>C</sub> - composite weight

 $w_d$  - dry composite weight

 $w_f$  - fiber weight

w<sub>m</sub> - matrix weight

 $w_w$  - wet weight of composite

W <sub>f</sub>	- strain energy stored in a volume v of fiber
<b>w</b> <sub>m</sub>	- strain energy stored in a volume v of matrix
×	- distance along span of beam
× <sub>n</sub>	- amplitude after n cycles have elapsed in free
	vibration decay
x <sub>o</sub>	- first amplitude considered in free vibration
	decay
x(t)	- excitation time history
X(f)	- Fourier transform of excitation time history
X (s)	- Laplace transform of excitaiton time history
X(f)	- complex conjugate of X(f)
У	- deflection of the beam
y(t)	- response time history
Y(f)	- Fourier transform of response time history
Y(s)	- Laplace transform of response time history
z	- curve fitting parameter
α	- parameter defined in Eq. 2.10
8	- parameter defined in Eq. 2.7
Yc	- specific weight of composite
Ϋ́f	- specific weight of fiber
Υm	- specific weight of matrix
δ	- logarithmic decrement
ε	- unit strain

 $\epsilon_{\rm L}$  - extensional strain

 $\epsilon_{T}$  - transverse strain

ζ - damping factor

n - loss factor

 $n_f$  - fiber loss factor

 $n_h$ ,  $n_{h1}$ ,  $n_{h2}$  - parameters defined in Eqs. 2.24, 2.29 and 2.30, respectively

n<sub>m</sub> - matrix loss factor

angle of the fibers with respect to the direction
 of the applied load

 $\lambda_n$  - eigenvalue for the nth mode

ν - Poisson ratio

v<sub>r</sub> - fiber Poisson ratio

 $v_{I,T}^{*}$  - composite complex major Poisson ratio

 $\nu_{LT}^{\prime}, \nu_{LT}^{\prime\prime}$  - real and imaginary components of the composite major Poisson ratio

 $\xi, \xi_1, \xi_2$  - Halpin-Tsai parameters

p - density of specimen

 $\rho_a$  - density of air

 $\rho_{c}$  - density of composite

 $\rho_{\mathbf{w}}$  - density of water

σ - applied stress

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of fiber extensional stress

 $\overline{\sigma}_{f}$  - average extensional stress in the fiber

 $\overline{\sigma}_{\mathbf{L}}$  — average extensional stress in the composite

 $\overline{\sigma}_{\underline{\phantom{a}}}$  - average extensional stress in the matrix

- interfacial shear stress

- mode shape function

 $\phi_n\left(x_o\right),\phi_n\left(s\right)$  - mode shape functions of the nth mode at location of the probe or at the end of the beam length, respectively

 $\omega$  - angular frequency

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 $\omega_{\rm p}$  - angular resonant frequency of mode n

#### 1.0 INTRODUCTION

"A composite material is a material system composed of a mixture or combination of two or more macroconstituents differing in form and/or material composition that are essentially insoluble in each other" [1]. One of the constituents is a continuous phase known as the matrix and the other(s) is discontinuous and usually stronger and harder than the continuous phase. It is called the reinforcement or reinforcing material [2]. Matrix materials vary from very soft organic foams to glasses, ceramics and metallic matrices. Reinforcing agents range from elastomers and cotton to extremely hard intermetallic whiskers with shapes varying from very small, round or angular particles to continuous, very long fibers. The composite materials exhibit the best qualities of their constituents.

Composite materials can be of two types: fiber reinforced or particle reinforced. The fiber reinforced composite materials can be single layered laminae or multilayered laminates. The layers can have continuous or discontinuous fiber reinforcement with preferred orientation (off-axis) or unidirectional fibers (aligned). Hybrid composites having either mixed fiber materials or mixed continuous/discontinuous fibers are also common.

The numbers in square brackets denote entries in the list of References.

Advanced fiber reinforced composite materials [3] with fibers such as boron or graphite brought improvements in strength and stiffness. Organic fiber reinforcements such as aramid improves damping characteristics, which means even more versatility for composite materials. The use of advanced composite materials is greatly increasing day by day. The range of applications goes from aerospace/aircraft/automotive structures to printed circuit boards, prosthetic devices, golf clubs, and tennis racquets [4].

Widespread use of composite materials is possible because they allow the designer to use their anisotropy to tailor the properties to reach specific requirements of the final product. This includes high strength applications, high stiffness applications, low thermal conductivity and thermal expansion applications, corrosion-resistant and dent-resistant applications, and designs where weight reductions or part consolidation are required.

A lot of work has been done on improvement and optimization of the strength and stiffness of composite materials [1,2,5], but many of the applications already outlined also require good vibration damping properties [4,6,7].

Damping is a measure of the total energy dissipated in any vibrating structure, so it is important to mention the primary sources of energy dissipation in fiber reinforced polymeric matrix materials.

It appears that the viscoelastic behavior of the bulk matrix and the interface, and the friction at the interface caused by relative motion between matrix and fiber are the primary sources for the good damping characteristics of these materials. Both effects are significant in discontinuous fiber composites, since high shear stresses are developed at the fiber matrix interface when the material is subjected to a cyclic strain. The matrix near the end of the short fibers undergoes a high cyclic shear strain that produces a significant viscoelastic energy loss.

The shear stress concentration may induce plastic effects as well as partial debonding at the fiber matrix interface that could result in slip between fiber and matrix with corresponding frictional losses. It is desirable to have a strong interfacial bond such that slip at the interface can be avoided, since this adversely affects the strength and stiffness of the composite.

In conclusion, the most viable mechanisms of enhanced energy dissipation appears to be the shear deformation in the matrix caused by shear stress concentration near the fiber ends [6], and the shear stress due to off-axis coupling effects, as the experimental results show.

Based on the stress transfer mechanism between fiber and the matrix, internal damping can be improved and optimized by varying several parameters: fiber aspect ratio (length/diameter), fiber matrix modulus ratio  $E'_f/E'_m$ , fiber

volume fraction,  $v_f$ , and fiber direction. Even more, random fiber orientation and hybridization appear to cause interlaminar shear stresses that may be additional sources of energy dissipation [8,9,10].

The damping properties of continuous fiber reinforced composites have been studied by several authors [11,12,13,14], but, as mentioned earlier, the possibility of improving and even optimizing the damping characteristics of composites using short reinforcing fibers [6,15,16], the lack of valid data on damping and the interest and support of the U.S. Air Force Office of Scientific Research, gave the author the motivation for this research.

The research described here is constrained to the improvement and optimization of internal damping of fiber reinforced composite materials by variation of fiber aspect ratio, type of fiber reinforcement and fiber direction. The present dissertation includes the development of techniques for fabrication of composite material laminated plates by using an autoclave style press cure and the development of a faster and more accurate damping measurement technique using an impulsive force, as described in Chapter 3.

The measured static and dynamic properties of bare resin specimens that are used in the micromechanical model of a single round fiber surrounded by a cylindrical matrix, as explained in Chapter 2, are included in Chapter 4.

The experimental data and corresponding correlation with the theoretical model for three advanced composite materials: graphite/epoxy, Kevlar / epoxy and boron/epoxy in aligned continuous and discontinuous fiber configurations are given in Chapter 5.

Finally, results from tests done on off-axis fiber specimens of graphite/epoxy and Kevlar/epoxy are compared with predictions in Chapter 6.

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Revlar is a proprietary aramid fiber produced by E.I. duPont de Nemours, Inc., Wilmington, Delaware.

#### 2.0 BACKGROUND

This chapter is written with the objective of reviewing the theory involved in the development of the present research.

# 2.1 PREDICTION OF COMPLEX MODULI FOR COMPARISON WITH DAMPING MEASUREMENTS

The purpose of this section is to present a theoretical approach to the prediction of the composite complex moduli from those of the constituent materials (matrix and fibers) for comparison, in later chapters, with measured composite properties. The analysis of internal damping and dynamic stiffness for aligned discontinuous fiber reinforced composite materials was based on the complex moduli as established in [6].

The complex modulus of a viscoelastic material can be defined in terms of the storage and loss moduli [17].

$$E^{\star} = E' + iE'' \tag{2.1}$$

Where the storage modulus, E', is related to the stored energy per unit volume, U, by Eq. 2.2,

$$U = \frac{1}{2} \epsilon^2 E' \tag{2.2}$$

and the loss modulus, E" is related to the specific damping energy, D, by Eq. 2.3.

$$D = \pi \varepsilon^2 E^n \tag{2.3}$$

The damping property, or loss factor, of materials is defined by the ratio of the energy dissipated per cycle or damping energy to strain energy, Eq. 2.4.

$$\eta = \frac{D}{2\pi U} = \frac{\pi \epsilon^2 E''}{2\pi (\epsilon^2 E'/2)} = \frac{E''}{E'}$$
(2.4)

The complex modulus generally depends on frequency. The real part, or elastic modulus, is generally less frequency dependent than is the imaginary part [18,19]. Experimental results shown in Chapter 4 for neat resin specimens, and in Chapter 5 for the composite material specimens tested, verify this prediction. In the same way, the complex modulus is temperature dependent [20] and care was taken in performing all the experimental work at the same ambient conditions.

The analytical model was based on the Cox stress distribution [21] which assumes:

- a.- A round fiber surrounded by a cylindrical matrix under extensional load, (Fig. 2.1).
- b.- Fiber and matrix are isotropic. In order to fit the experimental data to the analytical model, this assumption had

to be modified as shown later in this chapter, so that the fibers are considered to be anisotropic.

- c.- A perfect bond exists between the fiber and the matrix.
- d.- There is no load transfer through the ends of the fiber, and,
- e.- The transfer of load from the matrix to the fiber depends upon the difference between the actual displacement at a point on the interface and the displacement that would exist if the fiber were absent [22,23,24].

The interfacial shear stress,  $^{\tau}$ , and longitudinal fiber stress,  $^{\sigma}f$ , obtained by Cox were used by Gibson, et.al [6] as a starting point to find the solution for the complex modulus by two approaches, the Energy Method and the Force-Balance method.

The shear stress is:

$$\tau = \frac{r_0^2 E_f^1 \epsilon \beta \sinh[\beta(\ell/2 - x)]}{2r \cosh(\beta \ell/2)}$$
(2.5)

And the longitudinal stress is:

$$\sigma_{f} = \varepsilon E_{f}^{*} \left\{ 1 - \frac{\cosh[\beta(\ell/2 - x)]}{\cosh(\beta \ell/2)} \right\}$$
 (2.6)

Where  $\beta$  is:

$$\beta^{2} = \frac{G'_{m}}{E'_{f}} \frac{2}{r_{o}^{2} \ln(R/r_{o})}$$
 (2.7)

#### 2.1.1 ENERGY APPROACH

For the first approach, the stresses from Eqs. 2.5 and 2.6 are used to find the energy stored in the matrix and fiber and the resulting longitudinal complex modulus in terms of its component storage and loss moduli [17].

It is assumed that the matrix is linearly viscoelastic and that the interface has the same viscoelastic properties as that of the bulk matrix. It is stated that the strain energy is stored in both constituent materials but the dissipation of energy takes place only in the matrix, (i.e.,  $\eta_f=0$ ). Thus,

$$\eta = \frac{\eta_f W_f + \eta_m W_m}{W_f + W_m}$$

Since  $\eta_f = 0$ , the :

$$\eta = \frac{\eta_{m} W_{m}}{W_{f} + W_{m}} \tag{2.8}$$

After some manipulation [6], the longitudinal storage modulus is given by:

$$\frac{E_L'}{E_m'} = \frac{E_f'}{E_m'} v_f + v_m + \frac{E_f'}{E_m'} \left\{ \frac{v_f}{\beta \ell} - \frac{v_m}{\beta \ell \left[ (R^2/r_0^2) - 1 \right]} \right\} \left[ \frac{\alpha \sinh(\beta \ell) + \beta \ell}{1 + \cosh \beta \ell} \right]$$
 (2.9)

Where

$$\alpha = \frac{1 - v_f (3R^2/r_0^2 - 2)}{v_f (R^2/r_0^2) - 1}$$
 (2.10)

and the longitudinal loss modulus is given by

$$\frac{E_{L}^{"}}{E_{m}^{"}} = v_{m} \left\{ \frac{E_{f}^{'}}{E_{m}^{'}} \frac{1}{2[(R^{2}/r_{0}^{2}) - 1]} \left[ \frac{\tanh(\beta \ell/2)}{\beta \ell/2} - \frac{1}{\cosh^{2}(\beta \ell/2)} \right] + 1 \right\}$$
 (2.11)

The ratio  $R/r_0$  may be expressed in terms of the fiber volume fraction for a specified packing array.

For a hexagonal array:

$$\left(\frac{R}{r_0}\right)^2 = \frac{\pi}{2\sqrt{3} v_f}$$
 (2.12)

For a square array:

$$\left(\frac{R}{r_0}\right)^2 = \frac{\pi}{4 v_f}$$
 (2.13)

#### 2.1.2 FORCE-BALANCE APPROACH

In this approach, static equilibrium, geometric compatibility, and the elastic-viscoelastic correspondence principle [18] are used to obtain the complex moduli. In this derivation, the fibers are assumed to contribute to the energy dissipation and the fiber loss factor is non-zero.

Based on the fiber stress distribution (Eq. 2.6) the average fiber stress can be written as:

$$\overline{\sigma}_{f} = \frac{1}{\ell/2} \int_{0}^{\ell/2} \sigma_{f} d_{\chi} = \varepsilon E_{f}^{i} \left[ 1 - \frac{\tanh(\beta \ell/2)}{\beta \ell/2} \right]$$
 (2.14)

Considering static equilibrium, the total longitudinal stress applied to the composite is:

$$\vec{\sigma}_{L} = E_{L}^{\dagger} \epsilon = \vec{\sigma}_{f} v_{f} + \vec{\sigma}_{m} v_{m}$$
 (2.15)

The extensional strain for composite, fiber and matrix are assumed to be equal. The longitudinal storage modulus and loss modulus are then:

$$E_{L}^{i} = E_{f}^{i} v_{f} \left[ 1 - \frac{\tanh(\beta \ell/2)}{(\beta \ell/2)} - \frac{E_{f}^{m} v_{f}}{2} (n_{Gm} - n_{f}) \right]$$

$$\left[ \frac{\tanh(\beta \ell/2)}{(\beta \ell/2)} - \frac{1}{\cosh^{2}(\beta \ell/2)} \right] + E_{m}^{i} v_{m}$$
(2.16)

and

$$E_{L}^{"} = E_{f}^{"} v_{f} \left[ 1 - \frac{\tanh(\beta \ell/2)}{(\beta \ell/2)} \right] + \frac{E_{f}^{'} v_{f}}{2} (\eta_{GM} - \eta_{f})$$

$$\left[ \frac{\tanh(\beta \ell/2)}{(\beta \ell/2)} - \frac{1}{\cosh^{2}(\beta \ell/2)} \right] + E_{m}^{"} v_{m}$$
(2.17)

From these equations it can be seen that when the fiber loss factor is equal to zero, Eqs. 2.16 and 2.17 will be similar to Eqs. 2.9 and 2.11 obtained with the energy method. The force balance approach was used for the analysis of damping of composite materials as explained later in chapters 5 and 6.

Both analytical methods predict an optimum fiber aspect ratio (ratio of the length to diameter of the fiber) for maximum damping [6]. This prediction was one of the main reasons for this research, because this analytical result had to be checked against experimental values.

#### 2.2 OFF-AXIS ANALYSIS

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The previous analysis was used for the optimization of internal damping of fiber reinforced polymer composite materials when the applied load is parallel to the fiber direction. The following analysis will consider the variation and possible optimization of internal damping when the applied load is not parallel to the fiber direction.

Again the Force-Balance approach [6] and the elastic-viscoelastic correspondence principle [18] are used to obtain the equation for storage modulus, loss factor and loss modulus along the fiber direction [25].

The stress-strain relations for a lamina of arbitrary orientation under constant environmental conditions are given by

$$\{\varepsilon\}_{xy} = \overline{[S]}\{\sigma\}_{xy}$$
 (2.18)

Where

$$- T = [S] = [T] [S][T]$$
 (2.19)

is the off-axis compliance matrix, [S] is the compliance matrix along the fiber direction, and

$$[T] = \begin{bmatrix} \cos^2\theta & \sin^2\theta & 2\cos\theta\sin\theta \\ \sin^2\theta & \cos^2\theta & -2\cos\theta\sin\theta \\ -\sin\theta\cos\theta & \sin\theta\cos\theta & \cos^2\theta-\sin^2\theta \end{bmatrix}$$
 (2.20)

is the matrix that accounts for the tensor transformation of stresses and strains from the x axis to the principal axes (Fig. 2.2).

 $[T]^{T}$  is the matrix transpose of [T].

Each of the terms  $S_{ij}$  of the compliance matrix is a function of the engineering constants [3]. For the case of isotropic materials there are two such constants, the elastic modulus E and Poisson's ratio . For an orthotropic material in plane stress there are four independent material properties: the longitudinal modulus,  $E_{L}$ , the transverse modulus,  $E_{T}$ , the in-plane shear modulus,  $G_{T}$ , and the major Poisson's ratio,

The apparent modulus,  $E_{\chi}$ , along the direction of the applied stress is given in Eq. 2.21 as a function of the four independent properties already given and the angle of direction of the fibers with respect to the direction of the applied load,  $\theta$ , [3].

$$\frac{1}{E_{x}} = \frac{1}{E_{L}} \cos^{4}\theta + (\frac{1}{G_{LT}} - \frac{2v_{LT}}{E_{L}}) \sin^{2}\theta\cos^{2}\theta + \frac{1}{E_{T}} \sin^{4}\theta \qquad (2.21)$$

By using either of the models described in the previous section, the longitudinal modulus can be written as:

$$E_L = E_f \left[1 - \frac{\tanh(\beta \ell/2)}{(\beta \ell/2)}\right] v_f + E_m v_m$$
 (2.22)

Where  $\beta$  is given by equation (2.7).

The transverse modulus and the in-plane shear modulus can be obtained empirically by using the Halpin-Tsai equations [3], which are approximate representations of more complicated

micromechanics results and can be equally applied to fiber, ribbon or particulate composites:

$$\frac{M}{M_{m}} = \frac{1 + \xi \eta_{h} v_{f}}{1 - \eta_{h} v_{f}}$$
 (2.23)

Where

$$n_h = \frac{(M_f/M_m) - 1}{(M_f/M_m) + \xi}$$
 (2.24)

 $\xi$  is a measure of fiber reinforcement which depends on the fiber geometry, packing array and load conditions. For circular fibers in a square array, good results were obtained using  $\xi$  = 2 for the transverse modulus and  $\xi$  = 1 for the in-plane shear modulus [3].

The major Poisson's ratio can be obtained by using the rule of mixtures [3]:

$$^{\vee}LT = {}^{\vee}f^{\vee}f + {}^{\vee}m^{\vee}m \qquad (2.25)$$

At this point, the elastic-viscoelastic correspondence principle [26] can be used to transform the elastic constants to viscoelastic constants:

$$E_X^* = E_X^* + iE_X^*$$

$$E_f^* = E_f^1 + iE_f^0$$

$$E_{T}^{*} = E_{L}^{i} + iE_{T}^{i}$$

$$E_{T}^{*} + E_{T}^{i} + iE_{T}^{i}$$

$$G_{T}^{*} = G_{T}^{i} + iG_{T}^{i}$$

$$G_{M}^{*} = G_{M}^{i} + iG_{M}^{i}$$

$$G_{LT}^{*} = G_{LT}^{i} + iG_{LT}^{i}$$

$$v_{LT}^{*} = v_{LT}^{i} + iv_{LT}^{i}$$

$$(2.26)$$

After introducing the complex elastic constants in Eq. 2.21, the off-axis composite storage moduli and the loss moduli can be obtained by separating the real and imaginary parts and neglecting higher order terms in the loss factor (since  $\eta \ll 1$ ). These equations are very long, so only the general form is presented here. The complete mathematical manipulation is given in Appendix A of Reference [53].

In general terms,

(a) 
$$\frac{E_{x}'}{E_{m}'} = \phi_{1} \left[ E_{f}', E_{m}', v_{f}', v_{V}', (\ell/d), \eta_{f}', \eta_{m}', v_{f}', v_{m}', \theta, packing array \right]$$

(b) 
$$n = \phi_2 \left[ E_f', E_m', v_f, v_v, (e/d), n_f, n_m, v_f, v_m, \theta, packing array \right]$$
 (2.27)

(c) 
$$\frac{E_{m}^{"}}{E_{m}^{"}} = \phi_{3} \left[ E_{f}^{"}, E_{m}^{"}, v_{f}, v_{v}, (\ell/d), \eta_{f}, \eta_{m}, v_{f}, v_{m}, \theta, \text{packing array} \right]$$

The analysis up to this point is based on the assumption that matrix and fibers are each isotropic, so the properties of the fiber are the same in all directions and the shear modulus can be obtained using the well known elastic relation given by Eq. 2.28 [27].

$$G_f' = \frac{E_f'}{2(1 + v_f)}$$
 (2.28)

In the present research, graphite, Kevlar and boron fibers were used in the experimental part. These fibers are essentially anisotropic [26,28] which means that the elastic properties vary with direction. Whitney [29] showed that by treating the fiber as an orthotropic, transversely isotropic material, the micromechanics equations can be adapted easily to take care of its anisotropy. Thus, the composite transverse modulus is given by

$$E_{T} = E_{m} \left[ \frac{1 + \epsilon \eta_{h1} v_{f}}{1 - \eta_{h1} v_{f}} \right]$$
 (2.29)

Where:

$$\eta_{h1} = \frac{(E_{fT}/E_m) - 1}{(E_{fT}/E_m) + \epsilon_1}$$

and the in-plane shear modulus is given by

$$G_{LT} = G_{m} \frac{1 + \epsilon_{2} \eta_{h2} v_{f}}{1 - \eta_{h2} v_{f}}$$
 (2.30)

Where:

$$n_{h2} = \frac{(G_{fLT}/G_m) - 1}{(G_{fLT}/G_m) + \xi_2}$$

#### 2.3 VIBRATION OF A CONTINUOUS SYSTEM

For the development of the experimental technique, small cantilever beam specimens were used for transverse modes, and free-free with masses attached at both ends for longitudinal modes.

Cantilever beam specimens of aligned discontinuous fiber reinforced composites were used for the fiber/length studies. Under flexural vibration, no mixed coupling effects are present and the complex moduli can be obtained from measured resonant frequencies and the frequency equation for a cantilever beam. Since the laminate properties do not vary through the thickness, the measured flexural modulus is equal to the extensional modulus.

For the case of off-axis tests used for fiber direction studies, the flexural vibration technique cannot be used; since flexural-torsional coupling effects exist. These specimens were therefore tested in longitudinal vibration with masses attached at both ends. The masses were needed in order to adjust the resonant frequency of the first mode to the range of interest for this research. Again the complex modulus was obtained from

the measured resonant frequency and the corresponding frequency equation for a mass-mass beam.

## 2.3.1 TRANSVERSE VIBRATION OF A CONTINUOUS BEAM

The elastic moduli of aligned discontinuous fiber composites were found, as mentioned in the previous section, from cantilever beam specimens by exciting the flexural or transverse modes and measuring the modal frequencies.

The equation of motion for the cantilever beam of Fig. 2.3 is given by Eq. 2.31

$$EI \frac{\partial^4 y}{\partial x^4} + \rho A \frac{\partial^2 y}{\partial t^2} = 0$$
 (2.31)

The boundary conditions for a cantilever beam are:

At the free end

$$\frac{\partial^2 y}{\partial x^2} = \frac{\partial^3 y}{\partial x^3} = 0 \tag{2.32}$$

At the clamped end

$$y = \frac{\partial y}{\partial x} = 0 \tag{2.33}$$

Using Eqs. 2.31, 2.32 and 2.33 the resonant frequencies for each mode are found to be:

$$f_n = \frac{\lambda_n^2}{2\pi S^2} \left(\frac{EI}{\rho A}\right)^{1/2} \tag{2.34}$$

Where n = 1, 2, 3, ...

The eigenvalues  $\lambda_{n}$  for the cantilever beam are found by solving the transcendental Equation 2.35

$$\cos \lambda_{n} \cosh \lambda_{n} + 1 = 0 \tag{2.35}$$

Values of the eigenvalues for all the modes are given in [30].

#### 2.3.2 LONGITUDINAL VIBRATION OF A CONTINUOUS BEAM

The free-free beam with masses attached at its ends (Fig. 2.4), was used to obtain the elastic moduli under longitudinal vibrations.

The differential equation of motion is given in Eq. 2.36 [31].

$$\frac{\partial^2 \mathbf{u}}{\partial \mathbf{x}^2} = \frac{\rho}{\mathbf{E}} \frac{\partial^2 \mathbf{u}}{\partial \mathbf{t}^2} \tag{2.36}$$

The boundary conditions are:

At 
$$x = 0$$

AE 
$$\left(\frac{au}{ax}\right) = M_1 \left(\frac{a^2u}{at^2}\right)$$
 (2.37)

At x = S

$$AE\left(\frac{\partial u}{\partial x}\right) = -M_2\left(\frac{\partial^2 u}{\partial t^2}\right) \tag{2.38}$$

The natural frequencies are given by Eq. 2.39

$$f_n = \frac{\lambda_n}{2\pi S} \left(\frac{E}{\rho}\right)^{1/2} \tag{2.39}$$

Where n = 1, 2, 3, ...

The eigenvalues  $\lambda_{\mbox{\scriptsize n}}$  are obtained from the following transcendental equation:

$$tan\lambda_{n} = \frac{(M_{1} + M_{2}) AS_{\rho}\lambda_{n}}{M_{1}M_{2}\lambda_{n}^{2} - (A_{\rho}S)^{2}}$$
 (2.40)

The complete mathematical development of the equations for the transverse and longitudinal vibrations are given in Appendix B of Reference [53].

#### 2.4 MODAL ANALYSIS

The principal objective of this research was to experimentally characterize the dynamic properties of composite materials. The use of modal analysis helps by identifying the modes of vibration. Each mode of vibration has a specific

natural frequency, damping factor and mode shape that can be identified from any point of the structure [32].

Modal analysis is useful for the following reasons:

- a.- It allows verification and/or changes in the mathematical model that describe the structure under testing.
- b.- It helps to identify the weak points of the structure.
- c.- It provides information about unwanted noise or vibration.

#### 2.4.1 FREQUENCY RESPONSE TRANSFER FUNCTION

Modal analysis is based on the measurement of the frequency response function, or transfer function.

There are six useful forms of the transfer function. These are:

- a.- Compliance (displacement/force)
- b.- Dynamic or apparent stiffness (force/displacement)
- c.- Mobility (velocity/force)
- d.- Impedance (force/velocity)
- e.- Inertance (acceleration/force)
- f.- Dynamic or apparent mass (force/acceleration)

Computing techniques like the Fast Fourier Transform in combination with powerful instruments like Digital Fourier Analyzers provide fast and accurate determination of the frequency spectrum of a time-domain signal. In order to facilitate the implementation of modal analysis on a Digital Fourier

Analyzer, the relationships between the time, frequency and Laplace domain are greatly used. The mathematical tools are the Fourier and Laplace transforms.

In the time domain the system behavior may be determined in terms of the system impulse response h(t) using the convolution integral [33,34].

$$y(t) = \int_0^t x(\tau)h(t - \tau)d\tau \qquad (2.41)$$

The transfer function that relates the impulse to the output is expressed in the Laplace domain by:

$$H(s) = \frac{Y(s)}{X(s)}$$
 (2.42)

Where:

$$X(s) = \int_{0}^{\infty} e^{-st}x(t)dt \qquad (2.43)$$

$$Y(s) = \int_{0}^{\infty} e^{-st} y(t) dt$$
 (2.44)

$$H(s) = \int_{0}^{\infty} e^{-st}h(t)dt \qquad (2.45)$$

The s-plane is complex and the magnitude of any function such as the transfer function H(s) can be plotted as a three

dimensional surface above the complex s-plane [32]. The transfer function evaluated along the frequency axes ( $s=i\omega$ ) is the Fourier Transform, or the system frequency response function.

Since these variables cannot be measured in the Laplace domain, the Fourier transform is of great importance for computation [35]. The frequency domain transfer function is:

$$H(f) = \frac{Y(f)}{X(f)} = \frac{Y(f)X(f)}{|X(f)|^2}$$
 (2.46)

Where:

$$H(f) = \int_{-\infty}^{\infty} h(t)e^{-i2\pi ft}dt$$
 (2.47)

$$X(f) = \int_{-\infty}^{\infty} x(t)e^{-i2\pi ft}dt$$
 (2.48)

$$Y(f) = \int_{-\infty}^{\infty} y(t)e^{-i2\pi ft}dt \qquad (2.49)$$

and  $\overline{X}(f)$  is the complex conjugate of X(f).

The expanded numerator and denominator of Eq. 2.46 are called the cross spectrum between the input and output,  $G_{\chi y}$  (f), and the power spectrum of the excitation or input,  $G_{\chi}$  (f), respectively.

$$Y(f)X(f) = G_{XY}(f)$$
 (2.50)

$$\chi(f)\overline{\chi}(f) = |\chi(f)|^2 = G_{\chi}(f)$$
 (2.51)

There is an inherent bias error in measuring the cross spectrum. The magnitude of the bias error is inversely proportional to the number of averages made in the computation. The greater the noise, the greater the number of averages that should be done in order to approximate the value of the cross spectrum [35].

One of the most important advantages in computing the cross spectrum is that magnitude and phase can be computed at the same time, since both pieces of information are contained in it. Another advantage is that by using this formulation, any wave form that is Fourier transformable can be used, so it is not limited to just sinusoidal signals.

#### 2.4.2 COHERENCE FUNCTION

Another advantage of computing the cross spectrum is that it allows the computation of the coherence function between the input and output signals, defined by Eq. 2.52

$$\gamma_{xy}^{2}(f) = \frac{|G_{xy}(f)|^{2}}{G_{x}(f)G_{y}(f)}$$
 (2.52)

The coherence function will be 1 if there is no noise in the measurements and if the system is linear. It will be 0 when there is no correlation between the two signals. The coherence function is a good check of the influence of noise in the signals. It is the monitor of the quality of the transfer function.

## 2.4.3 MEASUREMENT OF DAMPING

Damping is a measure of the total energy dissipated in any vibrating structure. Its value is of great importance when designing dynamic structures. This section is written with the objective of reviewing the most commonly used methods of damping measurements. Most of the methods are based on the single-degree-of-freedom analysis of a single mode. This is appropriate, since analytical modal analysis shows that a continuous system can be modeled as an infinite number of single degree of freedom systems [34].

The damping properties in resonance are related to the amplification factor Q as follows:

$$Q = \frac{1}{2\varsigma} \tag{2.53}$$

The loss factor  $\eta$  is the inverse of the amplification factor. This is the parameter that will be used for the characterization of damping in this research.

$$\eta = \frac{1}{Q} \tag{2.54}$$

The loss factor can be measured from the frequency response function by several methods and some of the most popular are described below [36].

a.- Ratio of the static compliance to the compliance magnitude at the natural frequency.

$$n = \frac{\text{static compliance}}{\text{compliance magnitude at the natural frequency}}$$
 (2.55)

b.- Bandwidth technique (See Fig. 2.5).

$$\eta = \frac{\Delta f}{f_n} \tag{2.56}$$

c.- From the real component of the transfer function (see Fig. 2.6).

$$\eta = \frac{(f_a/f_b)^2 - 1}{(f_a/f_b)^2 + 1}$$
 (2.57)

d.- Coincident-quadrature or Nyquist plot (See Fig. 2.7).

$$\eta = \frac{f_2 - f_1}{f_n} \tag{2.58}$$

e.- Curve fitting based on a model of a single-degree-of-freedom or a multidegree-of-freedom system. In the first case the model is given by Eq. 2.59, where the variable for curve-fitting is the damping factor \$\zeta\$.

$$y(t) = \frac{x(t)}{\left\{ [1 - (\omega/\omega_n)^2]^2 + [2\zeta(\omega/\omega_n)]^2 \right\}}$$
 (2.59)

For a multidegree of freedom model, several modes are considered at the same time and exponential equations, polynomial equations or other equations can be used to fit the transfer function of the system.

There are other methods to obtain damping which are not based on the frequency transfer function. The most popular are:

a.- Free vibration decay, based on the decay of the oscillation in the time domain (See Fig. 2.8).

The loss factor is:

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$$n = \frac{k}{\pi} \tag{2.60}$$

Where the logarithmic decrement is:

$$\delta = \frac{1}{n} \operatorname{Ln} \frac{x_0}{x_n} \tag{2.61}$$

b.- Resonant dwell, which is based on the ratio of the amplitude of the input ,a<sub>i</sub>, to the amplitude of the response ,a<sub>r</sub> , at resonance under sinusoidal excitation.

$$\eta = K \frac{a_i}{a_r} \tag{2.62}$$

Where K is a factor which depends on boundary conditions and mode number.

c.- Hysteresis loop in a force-displacement curve (see Fig. 2.9) [38].

$$= \frac{D}{2\pi U} = \frac{a}{b}$$
 (2.63)

The choice of the technique depends on the properties of the material to be tested, the ranges of frequency and stress of interest, the type of excitation available (forced-sinusoidal, random or impulsive force), and apparatus available. In all cases, parasitic losses such as support friction and air drag must be minimized, since erroneous measurements of damping can be obtained.

#### 3.0 EXPERIMENTAL TECHNIQUE

From References [6,38,39] it can be concluded that, theoretically, discontinuous fiber reinforcement can be used to improve internal damping in fiber reinforced polymer composite materials. As shown in Reference 6, there is a theoretical optimum value for the fiber aspect ratio for which the loss modulus is maximum, but this is only obtained with very low fiber aspect ratios. In addition, theoretical optimum off-axis angles are shown in [25].

In this research, specimens with fiber length ranging from continuous to 1/16" (1.588 mm) were tested in order to verify predicted trends at low fiber aspect ratios, while continuous and 1/4" (6.35 mm) fiber length specimens were tested for different fiber orientations. The effects of the elastic modulus ratio (E'<sub>f</sub>/E'<sub>m</sub>) also were considered in these tests by using different kinds of fiber reinforcement: graphite, boron and Kevlar.

This chapter will cover the fabrication procedure, the development of the experimental technique for damping measurements, the experimental procedures for determination of fiber and void volume fraction and the Poisson's ratio for the different mposite materials used.

#### 3.1 FABRICATION OF THE SPECIMENS

The specimens used were small cantilever beams with global dimensions of 3/4" (19.05 mm) wide by 8" (203.2 mm) to 11" (279.4 mm) long. The thickness depended on the type of composite material under consideration, but twelve plies were used to fabricate all the specimens. The specimens were fabricated in the laboratories of the Mechanical Engineering Department at the University of Idaho. The procedure was similar to that reported in reference [6], except that only 12 plies were used and the autoclave curing process was simulated in a laminating press.

The composite materials are available in prepreg tape rolls of different width. For this research, three different composite materials were used:

- a.- Fiberite Hy-E1034C graphite/epoxy pre-preg with non-woven continuous T300 graphite filament and Fiberite 934 resin in a roll of 12" (304.8 mm) width.
- b.- Fiberite Hy-E1734A2 aramid/epoxy pre-preg with non-woven continuous Kevlar 49 filament and Fiberite 934 resin in a roll of 12" (304.8 mm) width.

Fiberite Corporation, 501 W. 3rd St., Winona, Minnesota 55987

c.- AVCO 5505/4 boron/epoxy pre-preg with continuous boron 4 mil filament and 5505 resin in a roll with three separated strips of 1" (25.4 mm) wide each.

As stated previously, References [6,38] predict that internal damping of composite materials can be improved with reinforcement having low fiber aspect ratio (ratio of fiber length to diameter of the fiber). Since the objective of this research is the optimization of internal damping of fiber reinforced composite materials, it was necessary to develop a suitable technique for the fabrication of plates of various fiber length. This includes cutting the pre-preg, curing the plates and machining the specimens to the actual size. Plates of graphite/epoxy, Kevlar/epoxy and boron/epoxy with fiber length ranging from continuous to 1/16" (1.588 mm) were fabricated with excellent results.

It was decided that the number of plies used to fabricate the plates were to be 12 after preliminary tests done by the author using 3M Scotchply Sp 322 T300 graphite/epoxy pre-preg tape. The thickness of the plate with 12 plies after curing gives the specimens enough rigidity for the tests to be done and the frequencies for the principal modes fall inside the range of interest for the present work (10-1000 Hz for flexural vibration, and up to 1300 Hz for longitudinal vibration).

AVCO, Specialty Materials Division, 2 Industrial Ave., Lowell, MA 01851

#### 3.1.1 CUTTING THE PREPREG TAPE

Two techniques were developed to cut the pre-preg tape at the desired fiber length. The first one was used on graphite/epoxy and Kevlar/epoxy while the second technique was used on boron/epoxy prepreg tape [39,40].

# A.- Graphite/Epoxy and Kevlar/Epoxy

These materials were cut in laminae having dimensions of 12" (304.8 mm)  $\times$  12" (304.8 mm) from a roll previously stored at 0 °F (-18 °C) with a standard utility knife with stainless steel blades. After warming up to room temperature, the tacky surface was pressed against a urethane surface that holds the ply to be cut in strips at the desired fiber length, while leaving the backing paper attached to the free surface. The measurements for the required fiber length were marked on the backing paper by using a pair of dividers. All the strips were sliced across the fibers beginning and ending within the edges of the pre-preg tape (Fig. 3.1) to prevent the edge fibers from being pulled out of the pre-preg. A contact paper was attached to the backing paper tape so that the strips would be aligned end-to-end. The single ply then was moved to the uncured laminate (Fig. 3.2). The alternating plies were layed up such that the fiber ends were staggered, with a final uncured laminate of

12" (304.8 mm)  $\times$  12" (304.8 mm). The contact and backing papers were removed from each ply either as one piece or separately.

The maximum number of specimens to be tested for each fiber aspect ratio was decided to be six, so, only laminated plates of 6" (152.4 mm) wide were required. Two laminated plates of 6" (152.4 mm) wide, each one of different fiber length, were placed side by side for curing purposes, since the aluminum mold was designed for plates of dimensions 12" (304.8 mm) x 12" (304.8 mm).

# B.- Boron/Epoxy

For boron/epoxy composite a different technique was used. After allowing the pre-preg to warm up to room temperature, the tape was arranged into a 3" (76.2 mm) wide strip and then it was cut in 12" (304. mm) lengths. The boron fibers were then sandwiched between the backing paper and a transparent contact paper previously marked at the desired fiber length on a thin strip of white paper. The 3" (76.2 mm) x 12" (304.8 mm) ply was then fractured at each mark using a guillotine device with a 3" (76.2 mm) carbide blade as shown in Fig. 3.3 The fibers were sheared against a urethane cutting surface. Two of these plies were placed side by side in order to increase ply width to 6" (152.4 mm). The backing paper was removed and the ply was pressed against the uncured laminate. The alternating plies were layed up such that the fiber ends were staggered. Then the

contact paper was carefully removed. Two of these uncured laminates, each one of 12 plies and dimensions 6" (152.4 mm) x 12" (304.8 mm) were finally placed side by side in the mold for the cure process.

#### 3.1.2 CURE PROCESS

Two different processes [39,40] were followed; one for curing the composite plates with specific variations for each one of the materials used, and the second process to cure neat resin plates. Specimens of pure matrix had to be tested in order to obtain properties that were applied in the micromechanics model.

## A.- Composite material cure process

The 12" (304.8 mm) x 12" (304.8 mm) laminated plates were cured using an autoclave style press cure with a specially designed vacuum mold as shown in Fig. 3.4. Table 3.1 shows a listing of the components used in the lay-up of the mold. Fig 3.5 shows the general set-up for the curing process.

The general procedure is as follows:

a.- Spray the mold, pressure plate, vacuum bag, rubber dam and

spacer plate with release agent Frekote 33. It requires at least 15 minutes to set.

- b.- Place the rubber dam and the uncured laminated plate (which was previously warmed up to room temperature) in the mold.
- c.- Place the porous release fabric and two layers of bleeder cloth (cut to oversized dimensions) over the uncured laminate.
- d.- Place the pressure plate in the mold.
- e.- Fold the excess bleeder cloth over the rubber dam and press into the manifold along each edge.
- f.- Start the vacuum pump and place the vacuum bag over the top of the mold. Press along the edge of the mount to seal O-ring against mold.
- g.- Put the spacer plate over the pressure plate with the vacuum bag between them, then place the mold in the laminating press.
- h.- Set and start the data acquisition system for temperature recording with a thermocouple located in a small hole just below the vacuum fitting.
- i.- Follow the cure cycle as described in Table 3.2 or 3.3, depending on material used.

The quality of the plate obtained using the process just described is comparable to that obtained from an autoclave. Two

Frekote Inc, 170 W. Spanish River Blvd, Boca Raton, FL 33431

identical 12 ply unidirectional continuous fiber laminates were layed up using Fiberite Hy-E1034C graphite/epoxy pre-preg. One plate was cured in a laboratory autoclave at the University of Florida and the other was cured by the autoclave style press cure at the University of Idaho. The density, fiber and void volume fractions were determined using the process described later in this chapter from four locations in each plate. The two plates were nearly identical, as seen in Table 3.4.

Micrograph sections of the cured laminates also show the good quality and low void content of the composite materials obtained using the autoclave style press cure. Micrographs of graphite/epoxy composite obtained from an autoclave, and graphite/epoxy, Kevlar/epoxy and boron/epoxy composites produced by an autoclave style press cure may be seen in Figs. 3.6, 3.7, 3.8 and 3.9, respectively.

#### B.- Resin Cure Process

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Resin castings of Fiberite 934 and AVCO 5505 having dimensions 12" (304.8 mm) x 12" (304.8 mm) x 1/8" (3.18 mm) were produced in an aluminum mold following the cure cycles outlined in Tables 3.5 and 3.6, respectively 40. The mold and the equipment required to produce neat resin plates are shown in Figs. 3.10 and 3.11, respectively. These specimens were needed since the micromechanics model required input data on the dynamic properties of the matrix material [6].

#### 3.1.3 MACHINING OF THE SPECIMENS

The specimens were machined in the Machine Shop of the Mechanical Engineering Department using a precision reciprocating grinder. The cured laminated plate was fixed with double sided tape on 1" (25.4 mm) thick styrofoam, then placed on the working table of the grinder. Graphite/epoxy and Kevlar/epoxy laminated plates were cut with an abrasive cut-off wheel while the boron/epoxy plates were cut with a diamond cut-off wheel Fig. 3.12.

The machining of the graphite/epoxy and boron/epoxy specimens did not give any major complications, but the machining of Kevlar/epoxy specimens was problematic since delamination occurred at the edges of the specimens. Water was used as lubricant to machine this material. More recently, Kevlar/epoxy plates were machined with an electric band saw at high speed with the saw blade mounted backward (Fig. 3.13), followed by a trimming process in a vertical rotor with an opposed helix router tool.

In addition to the specimens for the vibration tests, small specimens of 1.5" (38.1 mm) x 3/4" (19.05 mm) were machined for determination of material density and fiber volume fraction. In order to minimize the effect due to delamination in Kevlar/Epoxy specimens, the length was increased to 2" (50.8 mm). Four specimens were machined for each fiber length plate for these type of tests.

The total number of plates fabricated were:

- a.- 13 plates (118 small beam specimens) of graphite/epoxy,
- b.- 8 plates (66 small beam specimens) of Kevlar/epoxy,
- c.- 4 plates (42 small beam specimens) of boron/epoxy,
- d.- 4 plates (6 small beam specimens) of neat Fiberite 934 resin, and,
- e.- 2 plates (6 small beam specimens) of AVCO 5505 neat resin.

# 3.2 DIGESTION TEST

This test was performed with the objective to determine the fiber volume fraction and void volume fraction of the cured laminated plates by digestion of the matrix resin in a liquid medium which does not attack the fibers excessively. It is also necessary to obtain the density of each plate.

The specimens used in these tests were of 3/4" (19.05 mm) x 1.5" (38.1 mm) for graphite/epoxy and boron/epoxy specimens and 3/4" (19.05 mm) x 2" (50.8 mm) for Kevlar/epoxy as specified previously.

# 3.2.1 DENSITY DETERMINATION

The American Society for Testing and Materials (ASTM) Standard Test Method for Specific Gravity and Density of Plastics by Displacement, D-792-66, was used to determine the density of the laminated plates. Room temperature and pressure

at the time of the tests were recorded. The weight of the specimens in the air and in the water then were used in Eq. 3.1 to determine the required density value. The procedure followed is outlined in Table 3.7 [39].

The equipment needed includes an analytical balance, beaker, thermometer, and a metal clip suspended by a nylon wire from a glass rod (Fig. 3.14). The density of the specimen is:

$$\rho_{C} = \frac{w_{C}(\rho_{W} - \rho_{a})}{w_{C} - w_{W}}$$
(3.1)

#### 3.2.2 FIBER AND VOID VOLUME FRACTION DETERMINATION

The American Society for Testing and Materials (ASTM) Standard Test Method for Fiber Content of Reinforced Resin Composites, D3171-76, Procedure A, was used with graphite/epoxy. For Kevlar/epoxy composite material specimens, the same standard was slightly modified. The digestion was obtained using nitric acid assay 70.6 % at 167 °F (75 °C) for 5 hours for graphite/epoxy and 1 hour for Kevlar/epoxy The temperature was obtained with water and a circulating pump heater as shown in Fig. 3.15 [39].

It was noted that bare Kevlar fibers digested for 1 hour increased their weight by approximately 2.5 %. The fiber weight in the composite was then corrected for this value in the respective calculations.

The digestion for boron/epoxy was based on the procedure outlined in reference [41]. The digestion was carried out using sulfuric acid assay 96.3 % at 400 °F (204.4 °C) for 16 hours. This temperature was reached with a hot plate (Fig. 3.16).

The calculations for all materials were based on the equation of the rule of mixtures for specific weight [2].

$$Y_C = Y_f V_f + Y_m V_m$$
 (3.2)

The fiber volume fraction is:

$$V_{f} = \frac{W_{f}}{\rho_{f}} \frac{\rho_{c}}{W_{c}}$$
 (3.3)

and the void volume fraction is:

$$v_v = 1 - \frac{w_f^{\rho}c}{w_c^{\rho}f} - \frac{w_m^{\rho}c}{w_c^{\rho}m}$$
 (3.4)

The calculations were done using the HP 9836 desk computer. The detailed procedure is given in Table 3.8 for graphite/epoxy and Kevlar/epoxy composite materials and in Table 3.9 for boron/epoxy.

## 3.3 DEVELOPMENT OF THE TECHNIQUE FOR DAMPING MEASUREMENTS

The number of tests to be done and the lack of valid data on internal damping of composite materials motivated the author to find an accurate and fast damping measurement technique that was easy to use and had the potential for making in-situ measurements for non destructive testing [42].

Gibson and Plunkett [43] developed a forced sinusoidal vibration technique for measuring damping and stiffness in composite materials, which was improved later by Gibson, et al [44]. The block diagram for this technique is shown in Fig. 3.17. This technique has been used successfully on numerous composite materials [19,20]. This technique was not used in the present research because it is too slow. Only one frequency can be excited at a time, and tuning the resonant frequency to obtain the Lissajous pattern on an oscilloscope is not easy. The other techniques described here excite all the modes at the same time, under random or impulsive excitation, and the transfer function in the frequency domain is generated in real time. The sinusoidal technique was used for comparison of results during the development of the new technique.

A 2024-T351 aluminum alloy specimen was used for calibration tests since the Zener Thermoelastic theory is quite accurate in predicting damping for structural metals under flexural vibration [43,45,46]. That is, experimental damping

values were compared with well-accepted values from thermoelastic theory.

Unreinforced epoxy resin specimens and composite material specimens such as continuous fiber graphite/epoxy (12 plies) and short fiber E-glass/polyester were also used during calibration tests. These materials were selected to give a range of damping from low (aluminum) to intermediate (composite) to high (epoxy).

#### 3.3.1 RANDOM TECHNIQUE

Fig. 3.18 shows the block diagram for this technique. This diagram is similar to the one used for the forced sinusoidal technique, except for the shaker input. The transfer function in the frequency domain is obtained with the Fast Fourier Transform (FFT) analyzer while the specimen is continuously excited with an electromagnetic shaker with a random input. The random signal (actually pseudo-random) is generated by the noise source of the FFT analyzer. This type of signal minimizes the leakage effects of a non periodic signal [47], which is the energy from the non periodic part that leaks on to the periodic part giving less accurate results.

The pseudo-random waveform provides the fastest means for finding statistically accurate transfer functions, which then are used to determine the internal damping of the material by means of the half-power bandwidth method as described later. A

typical transfer function showing several resonant peaks is shown in Fig. 3.19. Use of the zoom feature of the FFT analyzer makes it possible to analyze the peaks more accurately, as shown in Fig. 3.20.

The specimen used is a double cantilever beam, which is the same as that used for the forced sinusoidal technique 44 (Fig. 3.21). The specimen response is sensed by a non contacting eddy current proximicy transducer, whose calibration curve is shown in Fig. 3.22. The input is measured with a piezoelectric accelerometer.

Excitation and response signals are fed into the FFT analyzer and the desired transfer function on the desired frequency span is displayed on the screen in real time.

The principal characteristics of this excitation method are [47]:

- a. The force level is easily controlled (noise source knob and power amplifier knob).
- b. The transfer function spectrum can easily be shaped at resonant peaks with ensemble averaging, which removes extraneous noises, nonlinearities and distortion effects (RMS averaging and Hanning window function).
- c. No leakage errors are present
- d. A well designed fixture and exciter system are required. This means that this technique may not be suitable for in-situ testing of engineering structures.

In this case, the transfer function for the specimen is found by tapping the specimen with a hammer which has a force transducer attached at its head. Two specimens were used in this technique. A cantilever beam specimen (same as Fig. 3.21) was fixed in a vise at the resin shoulder. A free-free beam specimen was supported at the nodal points for the fundamental mode with nylon wire (Fig. 3.23).

The signals from the force transducer and from the motion transducer (non-contacting proximity probe) were fed into the FFT analyzer, which displayed the desired transfer function. The block diagram for this technique is shown in Fig. 3.24.

The principal characteristics of the impulse technique are [47]:

- a.\_ The force level cannot be accurately controlled.
- b. The transfer function can be shaped with small difficulty at resonant peaks using ensemble averaging which removes extraneous noises and, to some degree, nonlinearities and distortion effects (RMS averaging and uniform window function).
- c. Leakage errors may be present.

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d.\_ A special fixture and exciter system are not required, so that in-situ testing of engineering structures is possible. For both random and impulse techniques, care must be taken to avoid measurement of response near nodal points for the modes to be tested. Such measurements would consist primarily of noise, since the actual response is very small near nodal points. In addition, the amplitude of vibration for the response must be kept smaller than the thickness of the specimen in order to minimize air damping and support friction damping.

Since the frequency response function is in terms of the cross- spectrum [35], the coherence function, which is the ratio of the response power caused by the applied input to the measured response power [32], can be computed at the same time. The value of the coherence function is equal to 1 when the measured response is caused totally by the measured input. The coherence value is less than 1 when the measured response is greater than that due to the measured input, which means that extraneous noise or nonlinearities are contributing to the output power. Every transfer function should be accompanied by a coherence check.

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#### 3.3.3 DATA REDUCTION

## A.- Forced Sinusoidal Vibration Technique

The forced sinusoidal technique was used to check the results from the new techniques. Data reduction for this test was carried out by using the HP-85 micro-computer.

The basic equations used to find the loss factor and the storage modulus are based on the ratio of the base to tip amplitudes at resonance [48] and on the frequency equation for a cantilever beam vibrating in its nth mode [30,49], respectively:

The flexural loss factor is found from [48,50]:

$$\eta = \frac{C_n \phi_n(x_0) a_b}{\phi_n(S) \omega_n^2 a(x_0)}$$
(3.5)

The flexural storage modulus is;

$$E' = \frac{4\pi^2 S^4 AP f_n^2}{\lambda_n^4 I}$$
 (3.6)

The flexural loss modulus is then:

$$E'' = nE'$$
 (3.7)

## B.- Random and Impulse Techniques

The loss factor at resonant frequency  $f_{n}$  is obtained with the half-power bandwidth technique:

$$\eta = \frac{\Delta f}{f_n} \tag{3.8}$$

The half power points are found at 3 dB below the peak value of the transfer function for a specific mode when a logarithmic scale is used, or at 0.707 of this peak value when a linear scale is used. The data reduction is carried out by using a Hewlett Packard HP-85 microcomputer which is connected to the FFT analyzer using the HP-IB interface. After the desired number of ensemble averages, the desired transfer function is displayed on the screen of the FFT. The computer program "DAMP" (Reference [53]) reads the binary values from the memory of the FFT, makes the corresponding transformations to the current scale, finds the points on either side of the half-power points, and finds the half-power points by interpolation. The resonant frequency and the half-power bandwidth are then used in Eq. 3.8 to find the loss factor, while the flexural storage modulus and loss modulus are obtained using Eqs. 3.6 and 3.7 respectively. A different frequency equation is required for extensional vibration, however. The coherence value at resonance is recorded in each measurement.

#### 3.3.4 RESULTS

The descriptions of the specimens used in developing the technique are given in Table 3.10. Double cantilever beam specimens of aluminum, graphite/epoxy<sup>1</sup> (12 plies, continuous), chopped glass fiber reinforced polyester<sup>2</sup> and unreinforced epoxy<sup>3</sup> were used for the three techniques. Free-free beam specimens of aluminum (lowest loss factor) and glass/polyester (high loss factor) were used to perform another variation of the impulse technique.

The results are presented in figures for each one of the materials used. Fig. 3.25 shows the variation of loss factor of the aluminum cantilever specimen at four resonant frequencies for the three techniques compared with the theoretical curve of the Zener thermoelastic theory. This is a severe test of the apparatus, since composite damping is much greater than that of aluminum. In all three techniques, acceptable values for the first three modes were obtained, and this reinforces the validity of the forced sinusoidal technique as reference. In the fourth mode, the values are consistently greater than the predicted values, which probably means that another physical mechanism is present in addition to the thermoelastic mechanism is present in addition to the thermoelastic mechanism.

Paul, Minnesota 55101.

PPG-SMC-R65. Manufactured by PPG Industries, Fiber Glass Division, Pittsburg, PA 15222

Sp 322 T300. Manufactured by 3M Company, 3M Center, St. Paul, Minnesota 55101.

<sup>&</sup>lt;sup>3</sup> Epoxy for 3M Sp 322 graphite/epoxy prepreg tape, manufactured by 3M Company.

anism. Other possible mechanisms are grain boundary viscosity, point defect relaxations, eddy current effects, stress-induced ordering, and electron effects [17].

In Figs. 3.26, 3.27 and 3.28, the results for the forced sinusoidal, random and impulse techniques for glass/polyester, graphite/epoxy and unreinforced epoxy specimens, respectively are shown. The values obtained for each mode are comparable between the three techniques with acceptable data scatter.

Additional data were obtained by testing aluminum and glass/polyester free-free beam specimens using the impulse technique (Fig. 3.29). For very low damping (aluminum), the results are not acceptable. Experimental values, which are high with respect to the predicted values, indicate the presence of nonlinearities (probably from the nylon wires). For high material damping, this technique gives reasonable values with the additional advantage that it can be used to analyze more than five modes.

The coherence values recorded after ensemble averages for the random and impulse techniques had a large scatter, going from as low as 0.4 to 1 depending on the mode analyzed and the material of the specimen tested. The low modes generally had lower coherence compared with higher modes. In the same way, low damping materials had lower coherence than the high damping materials. Since the damping values so obtained show good agreement with those obtained using the technique of reference,

the relationship between the value of the coherence function and validity of the damping measurements is not very clear.

## 3.3.5 CONCLUSIONS ON CALIBRATION TESTS

- a.- The new random and impulse techniques give acceptable experimental results based on comparison with the previously developed sinusoidal technique.
- b.- Both new techniques are faster and easier to use than the previously developed sinusoidal technique.
- c.- The random technique requires equipment that may not be available for in-situ tests of structures, whereas the impulse technique can be easily used for such tests.
- d.- The amplitude of vibration must be kept below the thickness of the specimen in order to minimize air damping and support damping. This is more easily accomplished by using the impulse technique than by using either the random or forced sinusoidal techniques.
- e.- All things considered, the impulse technique with cantilever beam specimens was selected for the following tests of this research.

# 3.3.6 ADDITIONAL IMPROVEMENTS ON THE IMPULSE TECHNIQUE

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a.- The hand-held hammer used to excite the specimen produced some variability in successive measurements. This problem

was solved by replacing the hand-held hammer with a solenoid-type electromagnetic exciter (Fig. 3.30), which gives a more reproducible impulse, while the nonlinearities and extraneous noises were minimized (the coherence value was improved).

- b.- The amplitudes of vibration were sufficiently low with the electromagnetic exciter that the epoxy shoulders on the single cantilever beams or in the double cantilever beams specimens were not necessary. The results of the experimental values obtained using specimens with and without shoulders are shown in Fig. 3.31.
- c.- The input using the electromagnetic-hammer is nearly constant so that the response spectrum is nearly the same for repeated tests. The shape of the transfer function spectrum after ensemble averaging tends to approach that of the single test response spectrum in the vicinity of the resonant peaks. In addition, the loss factors obtained using the bandwidth method on the response spectrum do not vary with respect to those values obtained using the transfer function (Fig. 3.32). The advantages of using the response spectrum can be summarized as follows:
  - 1.- Data scatter is reduced, and,

2.- The experimental procedure is faster, since it is not necessary to perform ensemble averages at all.

#### 3.3.7 GENERAL TEST PROCEDURE FOR DAMPING MEASUREMENTS

The impulse technique as described previously was extended for extensional vibration tests for the off-axis fiber composite specimens.

Under transverse or flexural vibration, the beam flexes perpendicular to its own axis to alternately store potential energy in the elastic bending of the beam and then release it into the kinetic energy of transverse motion [30]. The normal stress parallel to the beam axis is the only stress acting on the beam. When the specimens used are aligned fiber composite materials, complex moduli can be determined directly from this test. In the case of the off-axis specimens, the normal stress can induce both normal and shear strain and correction for various coupling effects, such as flexural-torsional effects, would be necessary before comparing with theoretical prediction. Thus the isolation of damping due only to flexural deformation appeared to be very difficult. Extensional vibrations eliminate flexure-torsion coupling effects, but both normal and shear strains are still present. The tensor transformations in the theoretical off-axis analysis in Section 2.2 take this into account, however. Thus, the complex moduli obtained by this method can be compared directly with predicted values.

The following two sections describe the general procedure for each one of these techniques as they were used in this research.

#### A.- Flexural vibration tests

The block diagram of the flexural vibration apparatus is shown in Fig. 3.33. Small beam specimens were tested under transverse vibration. An aluminum foil target of 3/8" (9.53 mm) diameter was placed on each specimen at a distance equivalent to 70% of the effective length of the specimen from the clamped end. The effective length in a cantilever beam is the length from the clamp to the free end. The aluminum target is needed because the eddy current probe used to sense the response works only with metallic surfaces.

The signals from the response transducer and the force transducer attached to the tip of the electromagnetic hammer are fed into the FFT Analyzer as explained previously in this chapter. The computer program "DAMP" (Reference [53]), was used for data acquisition and reduction. The input data to run this program are the length, width and thickness of the specimen, the density obtained from the ASTM standard test D-792-66, the temperature at the time of the test, and the fiber and void volume fraction. The next step is to set the front panel of the FFT Analyzer for the impulse technique: sensitivity of channels

A and B, frequency mode, frequency span, type of spectrum, window function and scale.

After the frequencies for the first four modes were determined, the zoom feature of the FFT Analyzer was used to isolate each mode one by one. Three to four measurements were done for each mode of every specimen. The computer program makes the corresponding calculations and the hard copy of the output shows the resonant frequency, the loss factor, and the storage and loss modulus. The coherence function is printed only if the transfer function is used to calculate the loss factor. The storage modulus is calculated by using the resonant frequency for the mode under consideration and the frequency equation (Eq. 3.6) for a cantilever beam.

The set up used for this test is shown in Fig. 3.34.

#### B.- Extensional vibration tests

The block diagram of the extensional vibration apparatus for longitudinal vibrations is shown in Fig. 3.35. Small beam specimens with masses attached at each end were used with this technique (Fig. 2.4). Changes of length of the beam and the masses are necessary in order to maintain the frequency of the first mode inside the range of interest for this research. This is accomplished with the help of a computer program written in BASIC for the Hewlett Packard HP-85 micro-computer. This program calculates the eigenvalue for the mode of interest, and

by varying the masses or the length of the beam, the estimated value for the resonant frequency is obtained. This calculation is repeated until the estimated value is close to the pre-defined resonant frequency to be used in the tests.

This technique was used for characterization of damping as a function of fiber orientation. The impulse technique was used again but the beam-mass assembly was suspended from a steel bracket with nylon wires attached at each concentrated mass. The specimens were excited longitudinally with the electromagnetic hammer, while the response was measured with a piezoelectric accelerometer.

Data acquisition and reduction were the same as for the flexural vibration, except that Eq. 3.9 was used to calculate the apparent storage modulus along the axis of the beam, and the values for each one of the masses were part of the input for the program "DAMP".

$$E' = \frac{4\pi^2 S^2 \rho f_n^2}{\lambda_n^2}$$
 (3.9)

The set-up of the equipment used in this test is shown in Fig. 3.36.

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#### 3.4 POISSON RATIO DETERMINATION

The Poisson's ratios for the matrix and the composite materials are needed for the micromechanics model used for comparison with the experimental results.

Poisson's ratios for the bare resin specimens and for the composite materials were obtained by static loading the same small beam specimens of dimensions 8" (203.2 mm) long x 0.75" (19.05 mm) wide and 12 plies thick for the composite materials or 1/8" (3.175 mm) for the resin specimens.

SR-4 strain gages type FAET-12A-12513L<sup>1</sup> were installed on the specimens in two quarter Wheatstone bridge circuits. A Hounsfield tensometer was used to apply load to the specimens. A HP 3947A data acquisition system interfaced with a HP 9836 desk-computer was used to sense the imbalance in the bridges when the load was applied. Using a computer program, the longitudinal and transverse strains were obtained and Poisson's ratio was calculated from the slope of the line formed by plotting longitudinal versus transverse strains.

Equation 3.10 is the mathematical expression for Poisson's ratio [27] which is the negative of the ratio of the transverse to the longitudinal strains.

$$v = -\frac{\varepsilon_{\uparrow}}{\varepsilon_{\downarrow}} \tag{3.10}$$

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The results are plotted in Fig. 3.37 for Fiberite 934 resin, Fig. 3.38 for AVCO 5505 resin, Fig. 3.39 for graphite/epoxy composite, Fig. 3.40 for Kevlar/epoxy composite and in Fig. 3.41 for boron/epoxy composite.

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#### 4.0 TESTS ON NEAT RESIN SPECIMENS

The dynamic properties of neat resin are needed, as stated before, as input data for the micromechanics model used to predict the dynamic properties of composite materials. The matrix is assumed to be linearly viscoelastic in shear but linearly elastic in dilatation, thus the bulk modulus is real, non-dissipative and frequency-independent while the shear modulus is complex, dissipative and frequency-dependent [49]. Under these assumptions, the complex modulus for the matrix as a function of frequency was needed for the analysis of composite specimens at different resonant frequencies.

Two types of epoxy resin are used in the pre-preg tapes used in the fabrication of the composite specimens. Fiberite 934 resin is a component of the graphite/epoxy and kevlar/epoxy pre-preg tapes and 5505 AVCO resin is in the boron/epoxy pre-preg tape. Resin castings were produced as described in section 3.1.

Six small cantilever beam specimens for each reserverse tested under flexural vibration, as explained in section 3.3.7 part A. The geometric and physical properties of the specimens are given in Table 4.1, where each dimension is the average value of all specimens of the same group and material.

With the flexural test, dynamic properties of the resin for the first four modes were obtained. The resonant frequency for the fourth mode had a maximum value of 900 Hz, which was enough for analysis of the aligned fiber composites. The resonant frequencies for the first modes in the tests done with off-axis fiber composite specimens under extensional vibration, were larger than 900 Hz, so the dynamic properties of the matrix in this range were also needed. Extensional vibration tests, as described in section 3.3.7 part B, were used with Fiberite 934 resin specimens (Table 4.2), since graphite/epoxy and Kevlar/epoxy are the only off-axis composite specimens tested.

Since the resin is highly hygroscopic, the tests were done as soon as the specimens were fabricated with the least possible time between operations (curing, cutting and testing). Thus, moisture absorption should have no effect on the results.

## 4.1 FIBERITE 934 RESIN CASTING

The tests were done at an average temperature of  $70^{\circ}$  F  $\pm 2$  °F (21 °C  $\pm$  1 °C) and a humidity of  $40\% \pm 2\%$ . Sample outputs of the tests are included in Appendix C of Reference [53].

Figs. 4.1 and 4.2 show the variation of storage modulus and loss factor with frequency and the respective curve-fitting done using the computer program "CURVE" from the standard pack of the HP-85 micro-computer. The experimental data for the storage modulus were fitted with a linear regression given by Eq. 4.1:

$$E_m' = 571252.737 + 55.647 \times f \text{ (psi)}$$
 (4.1)

Where f is the frequency in cycles per second (Hz).

For high frequencies, care must be taken since the small beam specimens show a drop in the effective storage modulus due to the presence of shear deformation. The analysis leading to the frequency equation in section 2.3 neglects shear, however.

The logarithmic regression found for the loss factor is given by Eq. 4.2.

$$\eta_{\rm m} = 0.004 + 0.006 \times Ln \text{ (f)}$$
 (4.2)

The loss factor as seen in Fig. 4.2 shows a large rate of change for low frequencies followed by a small rate of change up to the largest frequency of the test, 1400 Hz.

It is important to point out that both tests used with this resin: (flexural and extensional vibrations) show consistent results due to the overlap obtained with storage modulus and loss factor data in Figs. 4.1 and 4.2, respectively.

#### 4.2 AVCO 5505 RESIN CASTING

Only the flexural vibration was performed to find the dynamic properties for this resin. Six specimens (Table 4.1) were tested at an average temperature of 72°F (22°C) and a humidity of 41% following the same procedure as for the Fiberite 934 resin castings.

Sample output of the data for these tests is given in Reference [53]. Figs. 4.3 and 4.4 show the experimental data for storage modulus and loss factor with its respective curve-fit.

The curve fitting of the data with "CURVE" gives a linear equation for the storage modulus, Eq. 4.3 and an exponential equation for the loss factor, Eq. 4.4.

$$E'_{m} = 570140.406 + 14.943 \times f \text{ (psi)}$$
 (4.3)

$$\eta_{\rm m} = 0.019 \times f^{0.096}$$
 (4.4)

Where f is the frequency in cycles per second (Hz)

## 5.0 TESTS ON DISCONTINUOUS ALIGNED FIBER REINFORCED COMPOSITES

In this chapter, the experimental results will be presented and compared with predictions from the micromechanics analysis [5,14] as shown in Chapter 2. Flexural vibration of small cantilever beam specimens (Fig. 3.33) was the technique used to test the discontinuous aligned fiber reinforced composites of three types: graphite/epoxy, Kevlar/epoxy and boron/epoxy. The geometry and physical properties for the specimens tested are given in Table 4.1. Data reduction was carried out using the computer program "DAMP", while for comparison purposes and curve fitting needed to find the best agreement between the experimental results and the predicted curves, the computer program "OFFTRI" was used (Reference [53]).

#### 5.1 EXPERIMENTAL RESULTS

Three tests were done for each mode on every specimen. Six specimens for each fiber aspect ratio (fiber length/fiber diameter) for each composite material: graphite/epoxy, Kevlar/epoxy and boron/epoxy were tested under a room temperature fo  $70^{\circ}$  F +  $2^{\circ}$ F (21 °C + 1°C) and humidity of  $40^{\circ}$ 8 + 2%. Sample outputs for each test are shown in Appendix C of Reference [53].

Tables 5.1, 5.2 and 5.3 give the experimental results for the first three flexural modes for graphite/epoxy, Kevlar/epoxy

and boron/epoxy, respectively. Each value in this table is the average of 18 data values (6 specimens and 3 tests done on each one).

In Figs. 5.1, 5.2 and 5.3 the storage modulus for each composite material is plotted against frequency while in Figs. 5.4, 5.5 and 5.6 the loss factors vs frequency are presented.

For comparison with experimental data, Eqs. 2.16 and 2.17 are used to generate the predicted curves. As predicted [6,51] the storage modulus is essentially independent of frequency, while it decreases with decreasing fiber length. The loss factor increases with increasing frequency and decreasing fiber length.

The experimental data for the first resonant frequency was manipulated in order to present dimensionless ratios of the composite storage modulus E' to the matrix storage modulus E' and the composite loss modulus E' to the matrix loss modulus E' . Respective data scatter is also shown in Tables 5.4, 5.5 and 5.6. The matrix properties are estimated from Eqs. 4.1 and 4.2 for the case of graphite/epoxy and Kevlar/epoxy and from Eqs. 4.3 and 4.4 for the case of boron/epoxy. The frequencies used for these calculations correspond to the average of all first resonant frequencies of the specimens tested.

Figs. 5.7, 5.8 and 5.9 show the variation of the ratio  $E_L^*$  with fiber aspect ratio experimentally and theoretically.

Figs. 5.10, 5.11 and 5.12 show the variation of loss factor with fiber aspect ratio experimentally and theoretically.

Figs. 5.13, 5.14 and 5.15 show the variation of the estimated values of  $E_L^{"}/E_M^{"}$  obtained experimentally with respect to fiber aspect ratio in comparison with the predicted theoretical curves.

The predicted curves were plotted using, as stated before, the computer program "OFFTRI" with the following input:

a.- Type of resin:

Fiberite 934 resin for graphite/epoxy and Kevlar/epoxy AVCO 5505 resin for boron/epoxy

b.- Frequency of the first resonant frequency:

54 Hz for graphite/epoxy

38.75 Hz for Kevlar/epoxy

51.62 Hz for boron/epoxy

c.- Type of fiber packing array considered:

Square array for all cases. As shown in [6], there is almost no difference in the predicted results using square or hexagonal packing array.

d.- Storage modulus for the fiber (values given by the manufacturers):

For T300 graphite filament: 33 x 10 psi (227.37 GPa)

For Kevlar filament: 18 x 10 psi (124.02 GPa)

For boron filament: 58 x 10 psi (399.62 GPa)

e.- Bulk modulus for matrix. These values are real and frequency independent [49,52].

$$K_{\rm m} = \frac{E_{\rm m}}{3(1 - 2v_{\rm m})}$$
 (5.1)

Where:  $K_m = bulk modulus for matrix$ 

 $E_m = Static elastic modulus for matrix$ 

 $v_m$  = Poisson ratio for matrix

First Eq. 5.1 was used to find the real and frequency independent bulk modulus for the matrix using the elastic modulus and Poisson's ratio at zero frequency (static properties). Eqs. 4.1 or 4.3 (Fiberite 934 or AVCO 5505 resins), were used for a frequency of 0 Hz. The Poisson's ratio is obtained using the rule of mixtures with the experimental values from chapter 3. Secondly, the frequency-dependent properties as  $v_m$  and  $G_m^*$  were found using again E' from Eq. 4.1 or Eq. 4.3 at the resonant frequency of interest and Eq. 5.1, for the case of the Poisson's ratio.  $G_m^*$  was found from Eq. 5.2.

$$G_{m}^{i} = \frac{3K_{m}E_{m}^{i}}{9K_{m} - E_{m}^{i}}$$
 (5.2)

The complete mathematical manipulation is given in Appendix A of Reference [53].

The values were:

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For graphite/epoxy composite:

$$E'_{m} = 5.713 \times 10^{5} \text{ psi (3.936 GPa) (dynamic value)}$$

 $v_f = 0.16$  (static value)

 $v_m = 0.345$  (static value)

 $k_{m} = 6.143 \times 10^{5} \text{ psi (4.233 GPa) (static value)}$ 

For Kevlar/epoxy composite:

 $E_{m}' = 5.713 \times 10^{5} \text{ psi } (3.936 \text{ GPa}) \text{ (dynamic value)}$ 

 $v_f = 0.376$  (static value)

= 0.345 (static value)

 $k_{\rm m} = 6.143 \times 10^5 \text{ psi (4.233 GPa) (static value)}$ 

For boron/epoxy composite:

 $E_{m}' = 5.701 \times 10^{5} \text{ psi } (3.928 \text{ GPa}) \text{ (dynamic value)}$ 

 $v_f = 0.21$  (static value)

 $v_{\rm m}$  = 0.318 (static value)

 $k_{m} = 5.221 \times 10^{5} \text{ psi (3.597 GPa) (static value)}$ 

f.- Fiber and void volume fraction (  $v_f$ ,  $v_v$ )

graphite/epoxy: 0.654, 0.0036

Kevlar/epoxy: 0.661, 0.0302

boron/epoxy: 0.607, 0.003

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The fiber and void volume fractions shown above are the average values of the experimental results obtained by digestion tests of the cured laminate plates fabricated for each composite material (Chapter 3).

g.- Fiber loss factor equal to zero for all three composites since no information about this factor is available (this will be used as a curve-fitting parameter as shown later, however). As can be seen in Figs. 5.7 through 5.15, the experimental values do not fit the predicted curves. While the experimental values for  $E_{L}^{\prime}$  / $E_{m}^{\prime}$  are smaller and shifted to the right, the experimental values for the loss factor are greater and again shifted to the right.

The next step in the analysis is to find the necessary parameters to best fit the predicted curves to the experimental values.

#### 5.2 CURVE-FITTING

The parameters which had the greatest uncertainties associated with them were used as curve-fitting parameters. For example, due to uncertainties in the values for the elastic modulus of the fiber given by the manufacturers, this parameter was used as a curve-fitting parameter. The values given by the manufacturers are based on bare fibers and apparently there is a reduction in the fiber properties in the pre-preg material. This parameter shifts the curves in Figs. 5.7 to 5.9 up and down.

The fibers do contribute damping, but, since no information on fiber loss factors is available, this is another parameter used for curve-fitting. The fiber loss factor was assumed to be independent of frequency, however. Variation of the fiber loss factor shifts the curves in Figs. 5.10 to 5.12 up and down.

The Force-balance model is based on a single fiber 6. However, the influence of contiguous fibers in fiber bundles is not known, so an "effective fiber aspect ratio",  $(1/d)_{eff}$ , is introduced as given in Eq. 5.3.

$$(\ell/d)_{eff.} = (\ell/d)Z$$
 (5.3)

Where Z is a curve fitting parameter that shifts the curves from Figs. 5.7 through 5.12 to the left or to the right.

Varying these parameters one at a time and using the computer program "OFFTRI" a good fit was obtained with the following values:

For graphite/epoxy composite:

 $E_f^* = 25.51 \times 10^6 \text{ psi } (175.76 \text{ GPa})$ 

Manufacturer's value: 33 x 10<sup>6</sup> psi (227.36 GPa)

 $\eta_{F} = 0.0015$ 

z = 0.03

For Kevlar/epoxy composite:

 $E_f' = 14.48 \times 10^6 \text{ psi } (99.767 \text{ GPa})$ 

Manufacturer's value: 18 x 10<sup>6</sup> psi (124.02 GPa)

 $\eta_{f} = 0.011$ 

z = 0.03

For boron/epoxy composite:

 $E_{f}^{*} = 55.43 \times 10^{6} \text{ psi } (381.913 \text{ GPa})$ 

Manufacturer's value: 58 x 10<sup>6</sup> psi (399.62 GPa)

 $\eta_{\epsilon} = 0.0019$ 

z = 0.4

Where the values for Ef were obtained by using the rule of mixtures with the experimental results of composite and matrix storage modulus for a continuous fiber [6].

The curves generated by using these values are shown in Figs. 5.16 through 5.27.

Figs. 5.25, 5.26 and 5.27 are the theoretical tridimensional plots for  $E_L^{\bullet}$  / $E_M^{\bullet}$ , loss factor and  $E_L^{\bullet}$  / $E_M^{\bullet}$  as function of fiber aspect ratio and frequency for graphite/epoxy. These curves are presented for an overall view of the results obtained with these tests.

The predicted loss factor had an optimum value for a very low fiber aspect ratio based on a single fiber (See Figs. 5.19, 5.20 and 5.21), which is not attainable in practice by using discontinuous fibers. Microfibers and whiskers may be in this range, presenting interesting perspectives on optimization of damping. The predicted loss factor increases asymptotically with increasing frequency.

The predicted stiffness is practically frequency independent but decreases with decreasing fiber aspect ratio.

The estimated loss modulus (product of storage modulus and loss factor) has a maximum value at an optimum fiber aspect ratio.

These results should have interesting implications, especially in structures subjected to high frequency vibrations.

The combination of low fiber aspect ratio and high frequencies produces large increases in the damping properties and a comparatively small reduction in stiffness.

Even though good agreement was obtained between the experimental data and the predicted curves after curve fitting, some of the limitations of the model should be mentioned. The model does not consider fiber interaction, since it is based on a single fiber. In addition, it does not consider interface effects that can have substantial influence on the shear strains.

#### 6.0 TESTS ON OFF-AXIS FIBER REINFORCED COMPOSITES

These tests were done in order to prove whether, as predicted [25], there is maximum damping for an optimum fiber direction.

Flexural vibration was not used to characterize damping in off axis fiber specimens since mixed modes would have been present due to coupling of torsional and shear deformation and extensional deformation due to flexure. Shear deformation has a substantial effect on the natural frequency of the beam, and, of course on the damping [30]. In addition, the isolation of damping due only to the flexural deformation appeared to be very difficult. The solution chosen was to excite the longitudinal modes with an impulsive force at one end of a free-free beam with masses attached at each end, as described in section 3.3.7, part B.

In Ref. [25] it is predicted that for fiber aspect ratios greater than 100, the damping of the off-axis composite becomes independent of the fiber aspect ratio, except for angles below 15. The lowest fiber aspect ratio, (based on a single fiber) practically attainable from a pre-preg tape of aligned fibers was 226.45 for graphite/epoxy and 132.98 for Kevlar/epoxy. In order to check the prediction above, it was decided to test specimens of 0°, 2.5°, 7.5°, 12.5°, 20°, 45°, 60°, 75° and 90° for two different fiber lengths (continuous and 1/4" (6.35 mm))

using graphite/epoxy composite. Continuous fiber Kevlar/epoxy composite was also tested, since this material had the best damping characteristics (Chapter 5). The geometry and physical properties for the specimens tested are shown in Table 4.2.

The computer program "DAMP" was used for data reduction while the program "OFFTRI" was used to generate the predicted curves.

#### 6.1 EXPERIMENTAL RESULTS

Three tests were done on every specimen under room temperature of  $66\,^{\circ}$  F  $\pm 2\,^{\circ}$  F (18  $^{\circ}$  C  $\pm 1\,^{\circ}$  C) and humidity of  $468\pm\,$  2%. Sample outputs for each test are shown in Reference [53].

Tables 6.1, 6.2 and 6.3 show the experimental results as functions of the direction of the fibers and the first resonant frequency. The storage modulus was obtained using Eqs. 2.39 and 2.40, the loss factor as explained in chapter 3, using the halfpower points and the bandwidth technique, and the loss modulus by simple multiplication of the storage modulus and the loss factor.

For comparison with the predicted curves, the dimensionless ratio of  $E_X''/E_M''$ , the loss factor and the dimensionless ratio of  $E_X'''/E_M'''$  are plotted against direction of the fibers in Figs. 6.1 through 6.9 for continuous and discontinuous fiber graphite/epoxy and continuous fiber Kevlar/epoxy specimens.

The predicted curves are plotted using the real and imaginary parts of Eq. 2.21 (Reference [53]) and the computer program "OFFTRI" with the following input:

- a.- Type of resin: Fiberite 934 resin
- b.- Frequency of the tests:

Graphite/epoxy: continuous fiber 1152.5 Hz

Discontinuous fiber 1206.15 Hz

Kevlar/epoxy: 895.1 Hz

- c.- Type of fiber packing array: Square array in all cases
- d.- Storage modulus for the fiber:

  For T300 graphite filament:  $33 \times 10^6$  psi ( 227.37 GPA)

  For Kevlar filament:  $18 \times 10^6$  psi ( 124.02 GPa)
- e.- Bulk modulus for matrix:

  Graphite/epoxy composite:  $K_m = 6.143 \times 10^5$  psi (4.233 GPa)

  Kevlar/epoxy composite:  $K_m = 6.143 \times 10^5$  psi (4.233 GPa)
- f.- Fiber and Void volume fraction (v<sub>f</sub>, v<sub>V</sub>)
   Graphite/epoxy composite: continuous fiber 0.675, 0.002
   Discontinuous fiber 0.678, 0.0067

Kevlar/epoxy composite 0.715, 0.0138

- g.- Fiber loss factor equal to zero in all cases
- h.- Halpin-Tsai parameters:
  - 2 for the transverse modulus
  - 1 for the shear modulus [3]

(76.2 mm) fiber length.

i.- Fiber aspect ratio:
 Graphite/epoxy continuous fiber: 10,870 corresponding to 3"

Graphite/epoxy discontinuous fiber: 905.8 corresponding to 1/4" (6.35 mm) fiber length.

Kevlar/epoxy continuous fiber: 3191.5 corresponding to 1.5"
(38.1 mm) fiber length.

j.- Parameter that changes the fiber aspect ratio: Z=1

As seen in Figs. 6.1 through 6.9, the experimental values do not fit the predicted curves. For  $E_X^{\prime}$  / $E_M^{\prime}$  the experimental values are smaller and shifted to the left. The loss factors are smaller from 90° down to 5° where they cross over the predicted curves. The necessity of using parameters that can be varied until the predicted curves fit with the experimental values is evident.

#### 6.2 CURVE-FITTING

The first parameters introduced were those that best fit the predicted curves in the previous case of aligned discontinuous fiber (fiber angle  $\theta$ = 0°) as given in Chapter 5. These values are:

a. - Fiber storage modulus:

For graphite/epoxy:  $25.51 \times 10^6$  psi (175.76 GPa)

For Kevlar/epoxy: 14.48 x 10<sup>6</sup> psi (99.767 GPa)

b.- Fiber loss factor:

For T300 Graphite filaments: 0.0015

For Kevlar filaments: 0.011

c.- Parameter for the effective fiber aspect ratio, Z equal to 0.03 for both composite materials.

These parameters corrected the curves for the 0 fiber direction. The remainder of this section deals with curve-fitting for other fiber angles up to  $90^{\circ}$ .

Up to this point the fibers have been considered to be isotropic. The next step was to consider the anisotropy of the fibers by introducing the transverse and shear moduli for the fiber as parameters for curve-fitting [29], as shown in Eqs. 2.29 and 2.30. That is, the fiber is assumed to be orthotropic and transversely isotropic. In the same way, the Halpin-Tsai parameters may be varied independently until a good fit can be obtained.

The final values obtained were:

a.- Transverse modulus for the fiber:

T300 graphite filament:  $2 \times 10^6$  psi (13.78 GPa)

Kevlar filament: 1 x 10 6 psi (6.89 GPa)

b.- Shear modulus for the fiber:

T300 graphite filament:  $4 \times 10^6$  psi (27.56 GPa)

Kevlar filament: 2 x 10 6 psi (13.78 GPa)

c.- Halpin-Tsai parameters:

Graphite/epoxy: 2 for transverse modulus and

4 for shear modulus

Kevlar/Epoxy: 2 for transverse modulus and

1 for shear modulus.

It has to be mentioned that when fiber anisotropy was introduced, the transverse loss modulus of the fiber was obtained by the product of the transverse storage moduli of the fiber and the extensional fiber loss factor. More research has to be done in this area since the transverse loss factor for the fiber should be different from the extensional loss factor.

Figs. 6.10 through 6.18 show the experimental values with the best predicted curves. It is interesting to note that the stiffness is maximum at 0° and continuously goes down, rapidly up to about 30° and then slowly up to 90°. The loss factor increases up to a maximum for an optimum fiber direction of approximately 30° then decreases slowly with increasing fiber direction. The loss modulus has a maximum between 10° to 15°. These results verify the predictions from [25] in the cases of graphite/epoxy composite, continuous and discontinuous fiber. For Kevlar/epoxy specimens, the same trend for the curves were obtained, but the results showed a higher damping and lower stiffness than expected. These results may be related to the poor machinability of the specimens that produced delamination along the edges, and, an unexpected variation on the cure cycle for the laminate used for the specimens of 75° and 90° fiber direction.

These results showed that the influence of off-axis fiber orientation on damping were greater than that of the fiber length.

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These results can be better visualized with the tridimensional plots of  $E_X^*$  / $E_M^*$ , loss factor and  $E_X^*$  / $E_M^*$  as functions of fiber direction and fiber aspect ratio for graphite/epoxy, continuous fiber that are shown in Figs. 6.19, 6.20 and 6.21, respectively.

## 7.0 CONCLUSIONS

This chapter summarizes the principal conclusions obtained during this research. They are separated into three sections: conclusions on the fabrication of the specimens, conclusions on the experimental techniques, and finally, conclusions on the analytical model and the results.

## 7.1 CONCLUSIONS ON FABRICATION OF THE SPECIMENS

- 1.- The technique used for cutting the pre-preg tape worked well. For graphite/epoxy and Kevlar/epoxy, fiber lengths as low as 1/16" (1.588 mm) were obtained using a utility knife with stainless steel blades. Boron/epoxy pre-preg tape was fractured with a guillotine device with carbide blade.
- 2.- The autoclave style-press cure yielded quality plates comparable to those obtained from an autoclave. The density, fiber and void volume fraction, and even loss factor and storage modulus were almost the same for both plates.
- 3.- Machining of graphite/epoxy and boron/epoxy were done without difficulties, using abrasive and diamond cut-off wheels, respectively. Machining of Kevlar/epoxy specimens had complications since delamination was present at the specimen edges. The use of an electric band saw at high

- speed with the saw blade mounted backward gave the best results with this material.
- 4.- The digestion test was the step with more complications. First of all, there was not enough literature on the subject, so a great amount of time was expended in experimentation. For graphite/epoxy, the ASTM Standard D 3171-76, Procedure A was followed without problems. procedure with modifications was used for Kevlar/epoxy. It was noted that Kevlar filaments reacted with the acid and increased their weight, so, the fiber weight in the composite had to be corrected. complications were experienced with the boron/epoxy composite, since its digestion took over 16 hours for completion.

## 7.2 CONCLUSIONS ON THE EXPERIMENTAL TECHNIQUE

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- a.- The impulse technique using the electromagnetic hammer gave reliable loss factor measurements. The reason is because of the nearly constant input that minimizes nonlinearities and extraneous noises. The data scatter was reduced.
- 2.- The impulse technique is faster than the previously used techniques because of the use of the response spectrum. This eliminates the need for a large number of ensemble averages, as was the case when the transfer function was used for loss factor determination.

- 3.- Air damping, support damping and friction losses from the apparatus were minimized since the impulse from the electromagnetic hammer kept the maximum amplitude below the thickness of the specimen. This conclusion was also verified by the good results obtained for calibration tests done with aluminum specimens. This was a severe test for the apparatus, since damping of composite materials is much larger than that for aluminum.
- 4.- The results obtained from the flexural and extensional vibration tests showed good agreement since overlap was obtained.

# 7.3 CONCLUSIONS ON THE ANALYTICAL MODEL AND EXPERIMENTAL RESULTS

- 1.- The analytical model is based on a single fiber. The experimental results showed that there may be an influence of fiber bundle effects, since an "effective fiber length" had to be introduced as curve-fitting parameter. Fiber interaction may also contribute to this disagreement between experiment and theory.
- 2.- In the original analytical model, the fiber and matrix were assumed to be isotropic. The experimental results for off-axis composites showed that the fibers are anisotropic. The transverse modulus and shear modulus for the fiber had to be used as curve-fitting parameters.

- 3.- Experimental results showed that fibers do contribute to damping. The fiber loss factor had to be treated as a curve-fitting parameter, since data on loss factors of fibers are not available.
- 4.- The analytical model does not consider interface effects, which may be partly responsible for the "fiber damping".
- 5.- As predicted in [6], discontinuous fiber composite materials give optimum damping for very low fiber aspect ratios (based on a single fiber) not attainable in this research. These aspect ratios are in the range of whiskers and microfibers, however.
- 6.- The results showed, as predicted in [25], that even better damping can be obtained with off-axis fiber orientation, at the expense of reduced stiffness.
- 7.- The results showed that the Kevlar/epoxy composite has the best damping properties, while the boron/epoxy composite has the best stiffness properties.

Much work has to be done, and there are still questions that have to be answered. For example, the analytical model should account for fiber interaction and interface effects. In any case, the results obtained from this research verify the excellent properties and unlimited versatility of fiber reinforced composite materials.

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TABLE 3.1

List of layup materials used in the vacuum mold

Material	Trade Name
Bleeder cloth (polyester breather)	Air Weave N-4
Release fabric (porous teflon-coated	Release Ease 234 TFP
fiberglass, 0.003" thick)	
Vacuum bag (nylon, 0.002" thick)	Wrightlon 7400.002
Rubber Dam	Air Pad
Releasing agent	Frekote 33 <sup>2</sup>
Prepreg graphite/epoxy	Fiberite Hy-E1034C <sup>3</sup>
Prepreg Kevlar/epoxy	Fiberite Hy-E1734A2 <sup>3</sup>
Prepreg boron/epoxy	avco 55054

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4 AVCO Specialty Materials Division, 2 Industrial Ave., Lowell, MA 01851

Airtech International, 2452 E. Del Amo Blvd., P.O. Box 6207. Carson, CA 90749

Frekote Inc., 170 W. Spanish River Blvd., Boca Raton, FL 33431
Fiberite Corporation, 501 W. 3rd St., Winona, MN 55987

Cure cycle for Fiberite Hy-E1034C (graphite/epoxy) and Hy-E1734A2 (Kevlar/epoxy) prepregs. Resin system: Fiberite 934.

- 1.- Apply vacuum to mold-bag assembly
- 2.- Heat to 245 °F with a rate of 3 -5°F/minute
- 3.- Hold 15 minutes at 245 °F.
- 4.- Apply 100 psi pressure.
- 5.- Hold 45 minutes.

- 6.- Heat to 350 °F at a rate of 3 -5°F/minute.
- 7.- Hold at 350 °F for 2 hours.
- 8.- Cool to 150 °F under vacuum and pressure for graphite/epoxy (release vacuum and pressure before cooling for Kevlar/epoxy).

# Cure cycle for AVCO 5505 boron/epoxy prepreg

- 1.- Apply vacuum to mold-bag assembly.
- 2.- Heat to 245 °F at a rate of 4 -6°F/minute.
- 3.- Hold 30 minutes at 245 °F.
- 4.- Apply 50 psi pressure.
- 5.- Heat to 350 °F at a rate of 4 -6°F/minute.
- 6.- Hold at 350 °F for 2 hours.
- 7.- Cool to room temperature under vacuum and pressure.
- 8,- Oven post-cure at 375 °F for 4 hours.

TABLE 3.4

Comparison of volume fractions and density between autoclave and autoclave style press cure for graphite/epoxy composites.

	Autoclave	style	press cure	Auto	oclave Cu	ıre
	Density	v <sub>f</sub>	v.,	Density	<sup>v</sup> f	$\mathbf{v}_{v}$
	(g/cc)	(%)	(%)	(g/cc)	(%)	(%)
			<del></del>			
1	1.588	65.9	0.12	1.588	67.0	0.28
2	1.587	66.3	0.15	1.565	62.0	0.38
3	1.573	65.1	0.84	1.593	67.1	0.03
4	1.590	66.7	0.03	1.589	68.3	0.64
AVG.	1.585	66.0	0.28	1.584	66.1	0.33

Note.- Graphite/epoxy fabricated from Fiberite Hy-E1034C prepreg.

 $v_f$  = fiber volume fraction

 $v_v = void volume fraction$ 

Cure cycle for Fiberite 934 Resin 350 °F (177 °C) Cure

- 1.- Let resin warm overnight to room temperature.
- 2.- Heat in oven at 80 °C for 1 hour, then pour into vacuum flask.
- 3.- Apply vacuum while maintaining 80 °C temperature until resin has outgassed.
- 4.- Preheat mold to 125 °C, then pour resin into mold.

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- 5.- Cure at 100 °C for 16 hours, then at 177 °C for 2 hours.
- 6.- Removed cured resin plate from mold and allow to cool to room temperature.

Cure cycle for AVCO 5505 Resin 350 °F (177 °C) Cure

- 1.- Let resin warm overnight to room temperature.
- 2.- Heat in oven at 125 °C for 30 minutes, then pour into vacuum flask.
- 3.- Apply vacuum while maintaining 95 °C temperature until resin has outgassed.
- 4.- Preheat mold to 150 °C, then pour resin into mold.
- 5.- Cure at 100 °C for 16 hours, then at 177 °C for 2 hours.
- 6.~ Removed cured resin plate from mold and allow to cool to room temperature.

# Density determination procedure

American Society for Testing and Materials (ASTN) Standard Test Methods for Specific Gravity and Density of Plastics by Displacement, D-792-66.

# Procedure:

- 1.- Zero analytical balance.
- 2.- Weigh specimen in air to nearest mg.
- 3.- Fill 600 ml beaker with distilled water such that clip will be totally submerged.
- 4.- Weigh glass rod/clip assembly with clip submerged in water (do not allow it to touch side).
- 5.- Prewet composite specimen then place it in clip. Slowly lower specimen in water as to avoid air bubbles. Use wire to remove air bubbles.
- 6.- Weigh glass rod/clip/composite assembly.
- 7.- Record room temperature and pressure.
- 8.- Use tables to determine  $\rho_{w}$  and  $\rho_{a}$ .
- 9.- Make the corresponding calculations.

Fiber and void volume fraction determination for graphite/epoxy and Kevlar/epoxy composite specimen.

American Society for Testing and Materials (ASTM) Standard Test Methods for Fiber Content of Reinforced Resin Composites, D 3171-76.

### Procedure A:

- 1.- Determine density by ASTM Standard D 792-66.
- 2.- Zero analytical balance.
- 3.- Weigh specimen and crucible to nearest mg.
- 4.- Place composite specimen in beaker containing 40 ml of 70% nitric acid.
- 5.- After 5 hours digestion at 75 °C (167 °F) filter complete contents into crucible.
- 6.- To obtain proper vacuum, let drain tube fill with water by pinching the end.
- 7.- Rinse beaker into crucible with 20 ml nitric acid and filter a second time.
- 8.- Rinse 3 times with distilled water and twice with acetone. Rinse beaker each time.
- 9.- Place crucible on screen in core oven for 1 hour at 100 °C (212 °F).
- 10.- Cool crucible in desiccator, then weigh.
- 11.- Make the corresponding calculations.

Fiber and void volume fraction determination for boron/epoxy composite specimen.

### Procedure:

- -1.- Determine density by ASTM Standard D 792-66.
- 2.- Zero analytical balance.
- 3.- Weigh specimen and crucible to nearest mg.
- 4.- Heat 40 ml of concentrated sulfuric acid up to 204 C (400 F) using a hot plate.
- 5.- Place the specimen in beaker an digest it for 16 hours.
- 6.- Remove beaker and heat an additional 40 ml of sulfuric acid for rinsing.
- 7.- Partially fill the vacuum flask with cold water. Pour the contents of the beaker into the crucible and filter the excess acid.
- 8.- Rinse beaker with the hot sulfuric acid and pour into the crucible for a second rinse.
- 9.- Let the crucible to cool off for 10 minutes and drain the contents of the vacuum flask. Rinse 3 times with hot distilled water and twice with acetone. Rinse beaker each time.
- 11.- Cool crucible in dessicator, then weigh.
- 12.- Make the corresponding calculations.

TABLE 3.10

Description of specimens used for calibration tests.

	Len	gth	Thick	ness	Densi	ty
Specimen	in	mm	in	mm	lb/in <sup>3</sup>	g/cc
Aluminum (cantilever beam)	10.00	254.00	0.125	3.175	0.098	2.718
Aluminum (Free-free beam)	20.00	508.00	0.125	3.175	0.098	2.718
3M SP-322 Graphite/epoxy (Cantilever)	8.00	203.20	0.0655	1.664	0.0536	1.487
PPG-SMC R65 E-glass/polyester (Cantilever beam)	7.75	196.85	0.135	3.429	0.067	1.858
PPG-SMC R65 E-glass/polyester (Free-free beam)	16.00	406.4	0.135	3.429	0.067	1,858
Unreinforced epoxy (Cantilever beam)	8.25	209.55	0.125	3.175	0.0456	1,265

Note. The aluminum foil targets for the composite materials and unreinforced epoxy cantilever beam specimens are located at 7.375 in (187.3 mm).

Width for all specimens: 0.75 in (19.05 mm)

Resin for 3M SP-322 graphite/epoxy.

TABLE 4.1

Description of discontinuous aligned specimens tested with flexural vibration

No	Fiber	Length	Wi	dth	Thick	ness	Dens	sity	Vf
	(in)	(mm)	(in)	(mm)	(in)	(mm)	(lb/in)	(g/cc)	(%)
_			F	IBERITE S	34 RESIN	CASTIN	G		
6	N,	/ A	0.744	18.898	0.123	3.134	0.0469	1.298	N/A
				AVCO 550	5 RESIN	CASTING			
6	N,	/ A	0.753	19.126	0.119	3.023	0.0445	1.232	N/A
				GRAPI	HITE / EP	OXY			
6	7.5	190.5	0.746	18.948	0.062	1.575	0.0571	1.584	66
6	1/2	12.7	0.76	19.304	0.0588	1.494	0.0571	1.584	66.
6	1/4	6.35	0.733	18.618	0.06	1.523	0.0572	1.586	66.
6	1/8	3.175	0.749	19.025	0.061	1.554	0.0568	1.574	64.
6	1/16	1.588	0.753	19.126	0.0625	1.588	0.0567	1.574	63.
				KEV	LAR / EPO	XY			
5	7.5	190.5	0.755	19.177	0.0528	1.341	0.0492	1.362	65.
5	1/2	12.7	0.755	19.177	0.0542	1.377	0.0487	1.348	68.
6	1/4	6.35	0.74	18.796	0.0521	1.323	0.0477	1.323	66.
6	1/8	3.175	0.749	19.025	0.053	1.346	0.049	1.357	67.
6	1/16	1.588	0.746	18.948	0.0553	1.405	0.0498	1.382	63
				ВО	RON / EPO	XY			
6	7.5	190.5	0.745	18.923	0.0459	1.167	0.0743	2.061	<b>6</b> 3.
6	2	50.8	0.746	18.923	0.0481	1.221	0.0732	2.031	62.
6	1	25.4	0.749	19.012	0.0551	1.399	0.0724	2.009	60.
6	1/2	12.7	0.748	18.991	0.0532	1.351	0.073	2.024	60.
6	1/4	6.35	0.751	19.075	0.0487	1.237	0.0751	2.084	62
6	1/8	3.175	0.75	19.05	0.0504	1.279	0.0734	2.037	60.
6	1/16	1,588	0.751	19.075	0.0541	1.373	0.0704	1.954	56

Note. The length for all the specimens is 7.5 in (190.5 mm)Vf = Fiber volume fraction

Description of off-axis fiber specimens tested with extensional vibration

TABLE 4.2

No	Fiber	Le	ength	Wie	dth	Thick	ness	Mass	1	Mas	s 2
	angle	(in)	(mm)	(in)	(mm)	(in)	(mm)	(1b)	(kg)	(1b)	(kg)
				FIBERI	ГЕ 934 R	ESIN C	ASTING				
1	N/A	8.344	211.938	0.745	18.923	0.124	3.150	0.627	0.285	0.6316	0.287
1	N/A		141.3		18.948	0.126		0.627	0.285	0.6316	
1	N/A		47.625		18.948	0.126		0.627	0.285	0.6316	
1	N/A	0.875	22.225	0.746	18.948	0.126	3.2	0.627	0.285	0.6316	0.287
lot	e The	densi	ty for re	sin ca	sting is	0.046	9 lb/in	<sup>3</sup> (1.298	g/cc)		
			GR	APHITE	/ EPOXY	CONTI	NUOUS F	IBER			
6	Q	8.375	212.725	0.746	18.948	0.062	1.575	1.2615	0.573	1.2686	0.577
5	2.5	8.556	217.322	0.741	18.821	0.060	1.524	1.2615	0.573	1.2686	0.57
5	7.5	8.25	209.550	0.729	18.517	0.062	1.575	1.2615	0.573	1.2686	
6	12.5	7.896	200.558	0.73	18.542	0.060	1.524	0.946	0.43	0.9514	0.43
5	20	7.422	188.519	0.728	18.491	0.060	1.524	0.627	0.285	0.6316	0.28
6	45	5.037	127.94	0.731	18.567	0.060	1.524	0.305	0.139	0.3076	0.14
6	60	4.060	103.124	0.729	18.517	0.060	1.524	0.305	0.139	0.3076	0.14
6	75	3.058	77.673	0.729	18.517	0.061	1.549	0.305	0.139	0.3076	0.14
6	90	3.037	77.140	0.729	18.517	0.061	1.549	0.305	0.139	0.3076	0.14
Not			.0573 lb/			.)					
			ume fract								
	Voi	d volu	me fracti	on 0.2	%						
			GRAPHITE	/ EPO	XY DISCO	NTINUO	US FIBE	R			
6	0	8.125	206.375	0.733	18.618	0.060	1.524		0.573	1.2686	0.57
5	2.5		236.398	0.73	18.542		1.575	0.946	0.43	0.9514	
5	7.5		225.120		18.517		1.549	0.946	0.43	0.9514	
6	12.5		176.098		18.567		1.549	0.946	0.43	0.9514	
5	20		150.825	0.729	18.517	0.063	1.600	0.627	0.285	0.6316	0.28
6	45		115.240		18.339		1.549	0.305	0.139	0.3076	
6	<b>6</b> 0	3.542			18.593		1.575	0.305	0.139	0.3076	
5	75	3.063	77.788	0.725	18.415	0.062	1.575	0.305	0.139	0.3076	0.14

Note.- The fiber length for these specimens is 1/4 in (6.35 mm), Density  $0.0573~lb/in^3$  (1.589 g/cc), and Fiber volume fraction 67.8 % Void volume fraction 0.67 %

77.648 0.724 18.390 0.063 1.600

0.305

0.139

0.3076 0.14

90

3.057

TABLE 4.2 (continuation)

Description of off-axis fiber specimens tested with extensional vibration

No	Fiber	Le	ength	Wic	ith	Thick	ness	Mass	1	Mass	s 2
	angle	(in)	(mm)	(in)	(mm)	(in)	(mm)	(lb)	(kg)	(lb)	(kg)
			к	EVLAR ,	/ EPOXY	CONTIN	UOUS FI	BER			
4	0	8.375	212.725	0.755	19.177	0.053	1.346	1.2615	0.573	1.2686	0.577
4	2.5	7.914	201.016	0.733	18.618	0.053	1.346	1.2615	0.573	1.2686	0.577
4	7.5	7.899	200.635	0.734	18.644	0.054	1.372	0.946	0.43	0.9514	0.432
5	12.5	6.938	176.225	0.742	18.847	0.053	1.346	0.627	0.285	0.6316	0.287
3	20	4.464	113.386	0.753	19.126	0.054	1.372	0.627	0.285	0.6316	0.287
6	45	3.072	78.029	0.75	19.05	0.054	1.372	0.305	0.139	0.3076	0.14
4	60	2.089	53.061	0.742	18.847	0.054	1.372	0.305	0.139	0.3076	0.14
4	75	1.553	39.446	0.75	19.05	0.057	1.448	0.305	0.139	0.3076	0.14
4	90	1.442	36.627	0.762	19.355	0.055	1.397	0.305	0.139	0.3076	0.14

Note.~ Density 0.05 lb/in<sup>3</sup>(1.387 g/cc) Fiber volume fraction 71.5 % Void volume fraction 1.38 %

ABLE 5.1

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Experimental data for the first three modes in flexural vibration for graphite/epoxy specimens

	- 1						E	ċ	>	*
r Der	(mm)	300	(HZ)	Mpsi	GPa	Kpsi	MPa	0	<b>%</b>	<del>%</del>
		-	40 145	16 88	116.3	37.08	255.48	2.197		
	2000101400	۰ -	376.35	16.83	115.96	43.32	298.47	2.574	99	0.284
	8000	ı m	1053.23	16.81	115.82	47.75	329	2.84		
		į						300		
		-	54.125	15.17		45.43	313.01	2.992	,	6
172	12.7	α.	340.83	15.32	105.55	58.54	403.34	3.821	66.2	O.4
1		m	951.5	15.25		65.25	449.57	4.279		
								100		
		-	54.194	14.86	102.39	47.91	330.1	3.224	,	•
177	7 35	۰ ۸	368,688	14.76	101.70	74.27	511.72	5.035	9.99	0.38
	) )	i w	951.13	14.84	102.25	88.48	609.63	5.962		
		-	53.354	13.76	94.81	57.74	397.83	4.	;	(
1/8	3 175	~	334.235	13.76	94.81	80.84	556.99	5.875	04.	0.56
2		<b>(*</b> 7)	934.73	13.73	94.6	101.97	702.57	7.427		
1		-	48.397	10.69	73.65	9.69	479.54	6.511	;	,
1/16	1 588	۰ ،	308.815	11.07	76.27	97.75		8.83	63.2	0.20
2		. ~	870.85	11.22	77.31	111.99		9.981		
		•								

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POOL DESCRIPTION ASSESSMENT REPORTED FOR STANDARD FOR STANDARD ASSESSMENT ASS

Experimental data for the first three modes in flexural vibration for Kevlar/epoxy specimens

≯ <u>6€</u>	2.38	3.16	5.3	3.1	0.743
> €	65.8	68.1	66.7	67.1 3.1	63.0
60.	1.115 1.341 1.500	1.18 1.355 1.252	1.335 1.468 1.536	1.371 1.594 1.708	1.547 1.836 2.069
MPa	770.3 887.43 969.42	677.29 795.11 732.41	823,36 866.07 899.83	717.25 843.34 909.48	673.84 789.59 910.86
Kps i	111.8 128.8 140.7	98.3 115.4 106.30	119.50 125.70 130.60	104, 10 122, 40 132, 00	97.8 114.6 132.20
GPa	67.15 66.81 65.90	58.12 58.83 58.55	58.36 57.88 57.74	52.35 52.76 53.40	43.61 43.17 44.3
Mpsi	9.746	8.436 8.538 8.498	8.47 8.40 8.38	7.598	6.33 6.265 6.429
f (HZ)	41.54 259.84 722.48	40.5 255.33 713.3	39.15 244.380 683.55	37.19 234.5 659.23	35.35 220.32 624.75
Mode	- 0 m	- 0 m	-26	32	- 25
Fiber length in (mm)	nuous	12.7	6.35	3.175	1.588
Fiber	Cont innous	1/2	1/4	1/8	1/16

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TABLE 5.3

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Experimental data for the first three modes in flexural vibration for boron/epoxy specimens

Fiber		Mode	4	   		_		6	<b>*</b> ^	>
Ē	( IIII )		(HZ)	Мрѕі	c <b>P</b> a	Kpsi	MPa	<b>1</b> 0	<u>%</u>	ક્લે
		-	55.613	33.87	233.36	73.94	509.45	2.182		
Cont	Continuous	7	343.00	32.81	226.06	245.36	1690.53	7.466	63. ₹	0.20
		m	961.845	32.91	226.75	111.80	770.3	3.400		
		-	56.872	32.21	221.93	100.49	692.38	3.134		
~	50.8	2	353.801	31.75	218.76	181.27	1248.95	5.815	62.9	0.67
I	) 	m	992.353	31.85	219.45	129.03	889.05	4.589		
		-	58.38	25.58	176.25	91.30	629.06	3.416		
-	25.4	~	355.922	24.45	168.46	158.71	1093.51	6.614	4.09	0,64
		m	1012.66	24.97	172.04	125.11	862.01	5.086		
		-	55.421	24.94	171.84	97.79	673.77	3.967		
1/2	12.7	~	346.221	24.81	170.94	111.82	770.44	4.364	60.2	0
		٣	967.16	24.67	169.98	118.02	813,16	4.643		
		-	50.888	25.74	177.35	111.68	769.48	4.368		
1/4	6.35	2	313.849	24.93	171.77	133.90	922.57	5.420	62.0	0
		٣	879.14	24.94	171.84	141.37	974.04	5.722		
		-	46.24	19.47	134.15	130.4	858.46	6.746		
1/8	3.175	7	289.785	19.47	134.15	178.3	1228.49	9.225	60.3	0
		က	812.661	19.52	134.49	187.78	1293.80	9.783		
		-	37.926	10.90	75.10	131.15	903.62	12.05		
1/16	1.588	~	238.917	11.01	75.86	159.70	1100.33	14.522	56.0	0.3
		က	670.834	11.07	76.27	182.82	1259.63	16.535		
							-			

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TABLE 5.4

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Experimental results of the first resonant frequency for graphite/epoxy specimens, as used for comparison with the analytical model

Data scatter	(8/q)	E' /E'	E" /E"	1
Upper limit Mean value Lower limit	27174.00	30.38 29.39 28.02	3.46 2.314 1.49	3.432 2.197 1.403
Upper limit Mean value Lover limit	1811.59	30.51 26.41 22.67	3.16 2.835 2.42	3.681 2.995 2.259
Upper limit Mean value Lover limit	905.80	30.12 25.87 20.46	2.99	l
Upper limit Mean value Lower limit	452.90	25.81 23.96 19.83	3.604 3.13	1
Upper limit Mean vaiue Lower limit	226.45	20.20 18.61 15.28	4.344 3.57	7.488 6.511 5.843

Note.- Average frequency for the first mode: 54.00 Hz Average fiber volume fraction: 65.4% Average void volume fraction: 0.36% Fiber diameter: 0.000276" (0.00701 mm)

TABLE 5.

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Experimental results of the first resonant frequency for Kevlar/epoxy specimens, as used for comparison with the analytical model

	(0/1)	E. /E.	.3/3	Loss factor
upper timit Mean value Lower limit	15957.45	17.88 17.60 15.38	7.53	1, 155
Upper limit Mean value Lower limit	1063.83	14.71	6.62	1.350 1.180 0.9164
Upper limit Mean value Lower limit	531.91	16.06 14.77 12.90	8.26 8.05 6.85	1.335
Upper limit Mean value Lower limit	265.96	13.25	7.01	1.37)
Upper limit Mean value Lower limit	132.98	11.04	7.04 6.59 6.01	1.547

Note.- Average frequency for the first mode: 38.75 Hz Average fiber volume fraction: 66.1% Average void volume fraction: 3.02% Fiber diameter: 0.00047" (0.011938 mm)

TABLE 5.6

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Experimental results of the first resonant frequency for boron/epoxy specimens, as used for comparison with the analytical model

Data	(8/4)	E' /E'	E" /E"	Loss factor 103
Upper limit Mean value Lower limit	1875.00	61.30 59.32 58.23	5.22 4.67 4.21	2.182 1.99
Upper limit Mean value Lower limit	500.00	56.41 56.41 51.45	7.49 6.35 5.20	4.032 3.134 3.013
Upper limit Mean value Lower limit	250.00	48.49 44.80 40.09	10.37 5.77 3.96	3.416 2.349
Upper limit Mean value Lower limit	125.00	47.53 43.68 37.57	7.66 6.18 5.09	3.967 3.967 3.296
Upper limit Mean value Lower limit	62.50	49.21 45.08 41.59	8.66 7.06 6.50	5.768 4.368 3.819
Upper limit Mean value Lower limit	31.25	37.20 34.10 27.22	9.42 8.24 7.44	6.746 5.924
Upper limit Mean value Lover limit	15.625	21.44 19.09 17.99	9.28 8.29 7.23	13.390 12.050 11.140

Note.- Average frequency for the first mode: 51.62 HZ Average fiber volume fraction: 60.7% Average void volume fraction: 0.30% Fiber diameter: 0.004" (0.1016 mm)

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TABLE 6.

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Experimental results of the first resonant frequency for graphite/epoxy continuous fiber specimens tested with extensional vibration, as used for comparison with the analytical model.

Oata scatter	angle (deg)	(HZ)	Mpsi	GPa	Kpsi	мРа	Loss factor 103	E'/E'	£"/£"
	0.0	1256.94	18.555 127.84	127.84	54.77 377.00		3.701 2.947 2.045	30.02 29.22 28.55	2.40 1.86 1.26
1	2.5	1170.24	17.052 117.49	117.49	85.49	589.00	5.027 3.755	26.85 25.61	3.54 2.91 2.26
	7.5	1094, 74	14.176	97.67	112.85 778.00	778.00	9.732 8.006 6.223	22.32 21.01	4.60 3.84 3.11
	12.5	1091.93	10.438	11.92	139.68 962.00	962.00	13.412	16.44 15.26	4.75 4.35
	20.0	1114.27	6.804	46.88	124.98	124.98 861.00	19.390 18.398 17.660	10.71	4.25 4.16
	45.0	1161.70	2.432	16.76	53.76	53.76 370.00	22.048 20.930	3.83	1.83
	60.0	1108.70	1.799	12.40	37.98	37.98 262.00	23.680 21.113 19.280	2.83 2.74	1,29
	75.0	1199.92	1.563	10.77	30.26	30.26 208.00	24.360 19.410 17.620	2.46 2.36	1.03
	90.0	1174.18	1.489	10.26	28.33	28.33 195.00	21.200 19.060 17.250	2.34 2.22	1.01 0.96 0.92

Note.- Average frequency for the first mode: 1152.5 Hz Average fiber volume fraction: 67.5% Average void volume fraction: 0.20% Fiber aspect ratio: 10870

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Experimental results of the first resonant frequency for graphite/epoxy discontinuous fiber specimens tested with extensional vibration, as used for comparison with the analytical model.

Data	angle (deg)	f (HZ)	Mpsi	CP <sub>3</sub>	Kpsi	MPa	Loss factor	£,/E,	£"/£"
Upper limit Mean value Lower limit	0.0	1189.92	16.295 112.27	112.27	74.18 511.00	511.00	1	26.32 25.51 24.72	3.32 2.49 1.70
Upper limit Mean value lover limit	2.5	1259.45	15.990 110.17	110.17	78.49	541.00	5.757 4.915 3.391	25.57 25.04 24.69	3.05 2.64 1.83
Upper limit Mean value Lower limit	7.5	1180.63	13.602	93.72	114.28	114.28 787.00	10.340 8.431 7.393	21.84 21.30 20.17	3.84
Upper limit Mean value Lower limit	12.5	1189.19 10.700 73.72	10.700	73.72	127.00	127.00 875.00	11,872 10,780	16.88 16.75 16.53	5.24 4.27 3.90
Upper limit Mean value Lower limit	20.0	1236.77	604.9	44.16	112.54	112.54 775.00	19.580 17.584 16.630	10.03	3.78 3.57
Upper limit Mean value Lover limit	45.0	1233.80	2.457	16.93	57.44	57.44 396.00	24.800 23.388 22.520	3.85 3.75	2.02 1.93 1.86
Upper limit Mean value Lower limit	60.09	1192.43	1.756	12.10	37.86	37.86 261.00	25.750 21.592 19.920	2.75 2.68	1.27
Upper limit Mean value Lower limit	75.0	1204.11	1.545	1.545 10.65	29.34	29.34 202.00	19.470 18.996 18.150	2.53	0.99
Upper limit Mean value Lower limit	90.0	1169.05	1.435	9.89	27.21	27.21 187.00	19.610 18.974 17.560	2.25 2.19	0.94 0.91 0.88

Note.- Average frequency for the first mode: 1206.15 Hz Average fiber volume fraction: 67.8% Average void volume fraction: 0.67% Fiber aspect ratio: 905.8

TABLE 6.3

THE CONTROL OF CONTROL CONTROL

Experimental results of the first resonant frequency for Kevlar/epoxy continuous fiber specimens tested with extensional vibration, as used for comparison with the analytical model.

Data	ang le	f (HZ)	Mpsi	GPa	Kpsi GPa	GPa	Loss factor	£'/E'	E"/E"
Upper limit Mean value Lower limit	0.0	920.7	11.550	79.58	11.550 79.58 148.50	1.023		18.82 18.60 18.10	5.34
Upper limit Mean value Lower limit	2.5	2.5 872.5	10.130	69.80	145.60	1.003	1.618 1.444 1.163	16.31	5.24 4.46 4.75
Upper limit Mean value Lover limit	7.5	7.5 866.27	7.331	50.51	50.51 126.66	0.872	1.729	11.80	4.55
Upper limit Mean value Lower limit	12.5	947.95	5.137	35.39	91.56 0.631	0.631	1.859	8.27	3.29
Upper limit Mean value Lower limit	20.0	908.40	2.912	2.912 20.06	83.07 0.572	0.572	3.450 2.897 1.960	4.69	2.99 2.25
Upper timit Mean value Lower timit	45.0	844.81	0.852	5.87	29.93	0.206	4.963 3.544 2.569	1.37	1.08
Upper limit Mean value Lower limit	60.09	923.16	n69°0	4.78	22.55	0.155	3.246	1.12	0.81
Upper limit Mean value Lover limit	75.0	896.87	0.451	3.11	15.65	15.65 0.108	3.521	0.73 0.66	0.56
Upper limit Mean value Lower limit	90.0	875.28	0.415	2.86	21.41	21.41 0.148	4.423 3.114	0.67	0.77

Note.- Average frequency for the first mode: 895.1 Hz Average fiber volume fraction: 71.5% Average void volume fraction: 1.38% Fiber aspect ratio: 3191.5 175557

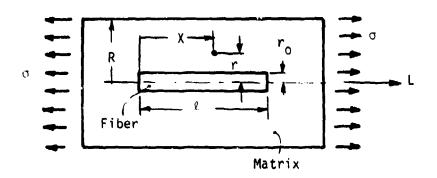


Figure 2.1 Representative Volume Element for Single Aligned Fiber

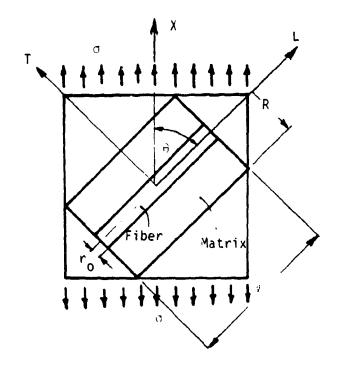


Figure 2.2 Representative Volume Element for a Single Off-Axis Fiber

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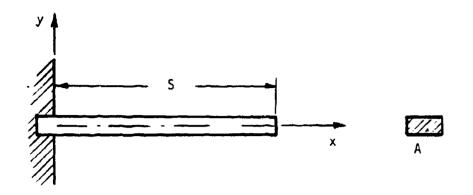


Figure 2.3 Cantilever beam specimen

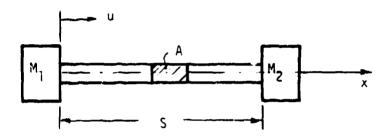


Figure 2.4 Mass-mass beam specimen

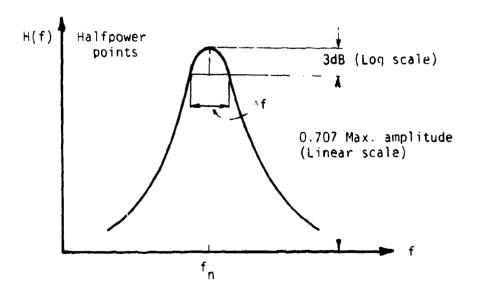


Figure 2.5 Transfer function vs frequency

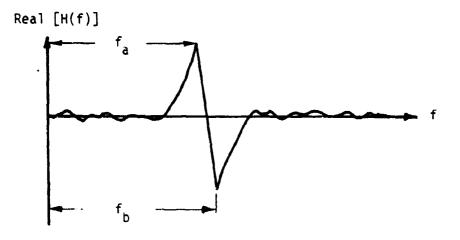


Figure 2.6 Real component of the transfer function vs frequency

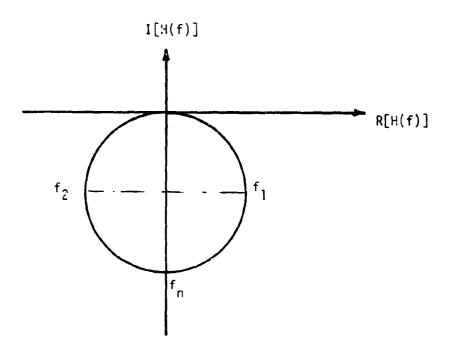


Figure 2.7 Nyquist plot

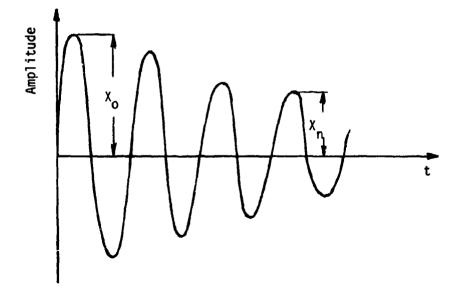


Figure 2.8 Free vibration decay

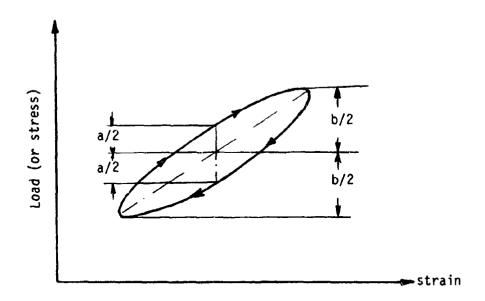


Figure 2.9 Hysteresis loop in a force-displacement curve



Figure 3.1. Slicing graphite/epoxy prepreg tape to produce discontinuous fibers.



Figure 3.2. Transferring single sliced ply to the uncured laminate.

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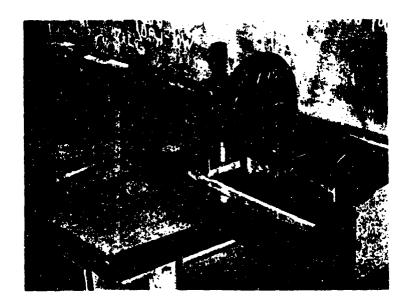


Figure 3.3. Set-up for cutting boron/epoxy prepreg tape.

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# AUTOCLAVE-STYLE PRESS CURE

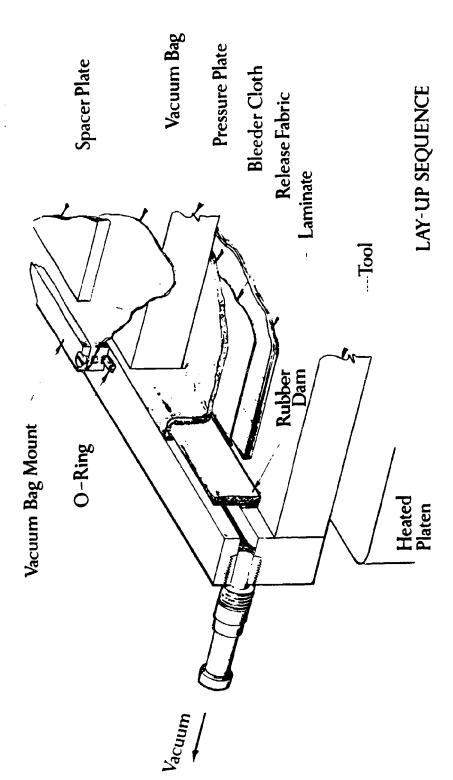
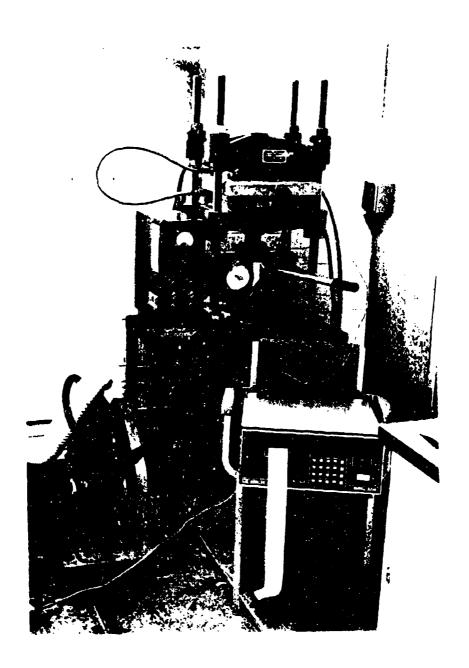


Figure 3.4. Cutaway view of autoclave-style press cure system.



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Figure 3.5. Set-up for autoclave sytle press cure.

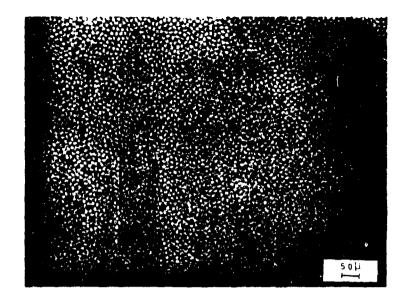


Figure 3.6. Graphite/epoxy composite produced by autoclave (scale 200:1).

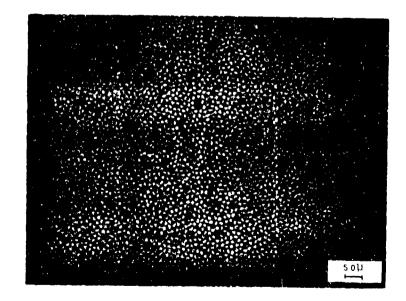


Figure 3.7. Graphite/epoxy composite produced by autoclave style press cure (scale 200:1).

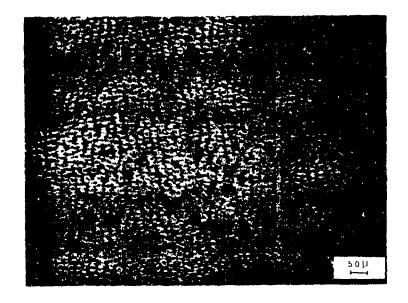


Figure 3.8. Kevlar/epoxy composite produced by autoclave style press cure (scale 200:1).

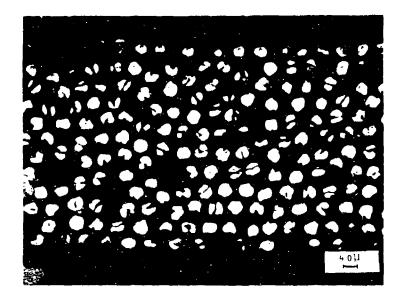


Figure 3.9. Boron/epoxy composite produced by autoclave style press cure (scale 40:1).

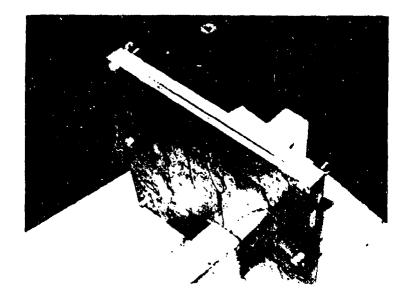


Figure 3.10. Mold to cast near resin plates.

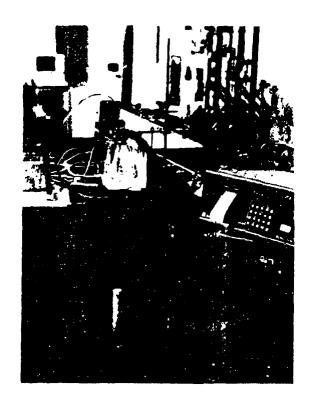


Figure 3.11. Equipment required to produce neat resin plates.

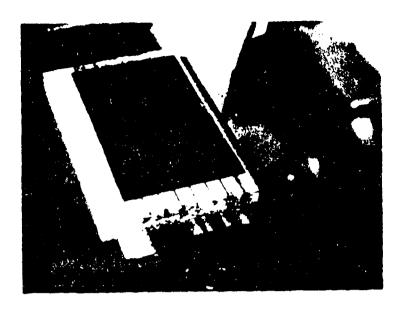


Figure 3.12. Precision reciprocating grinder used to machine graphite/epoxy and boron/epoxy composite specimens.

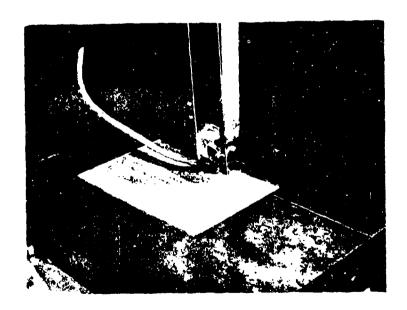


Figure 3.13. Electric band saw used to machine Kevlar/epoxy composite specimens.

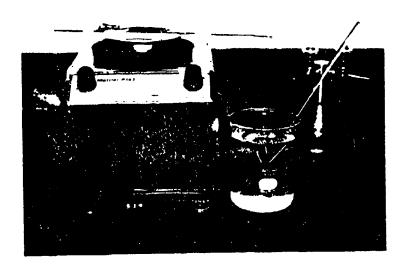
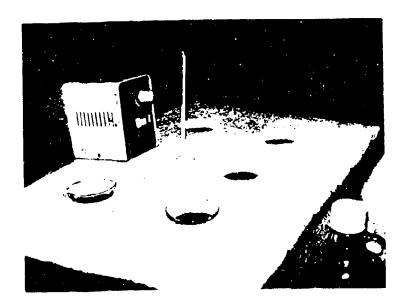


Figure 3.14. Equipment required to measure composite density.



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Figure 3.15. Water bath used for volume fraction determination of graphite/epoxy and Kevlar/epoxy composites.

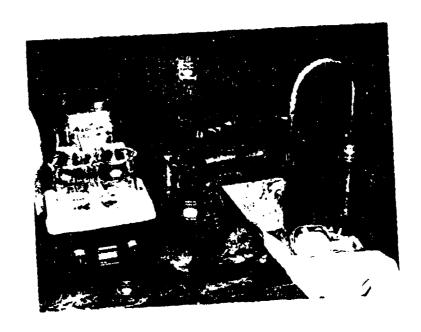
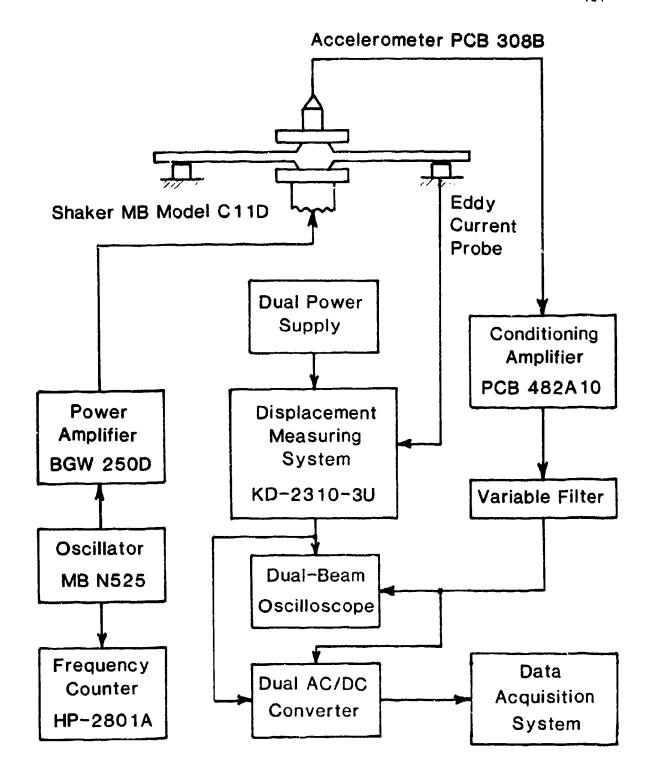


Figure 3.16. Hot plate used for the digestion of boron/epoxy composite.



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Figure 3.17. Block diagram of instrumentation for the forced sinusoidal vibration technique.

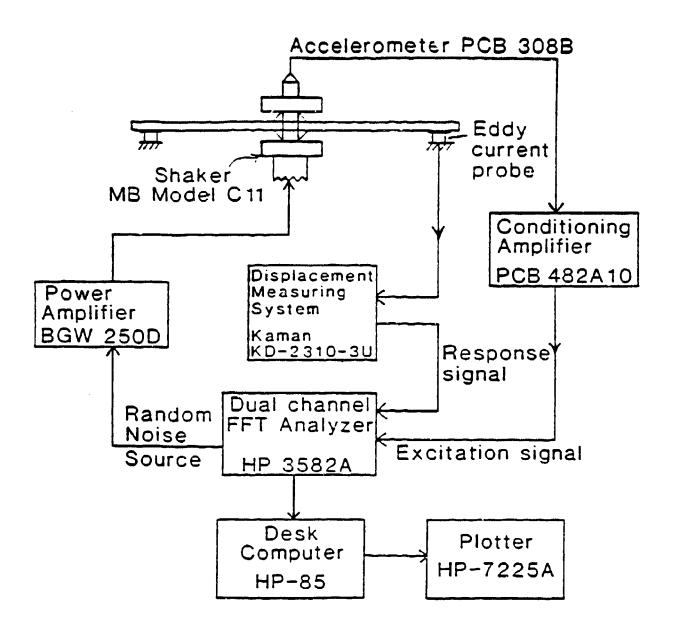
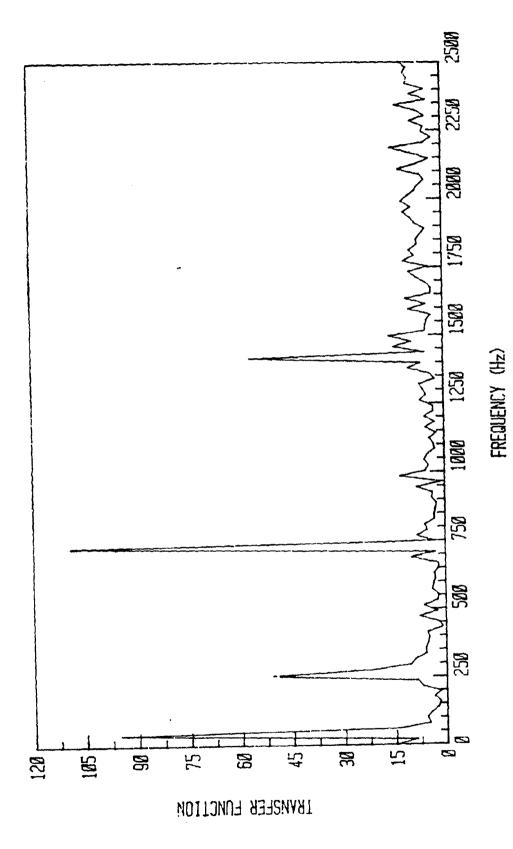


Figure 3.18. Block diagram of instrumentation for the random technique.

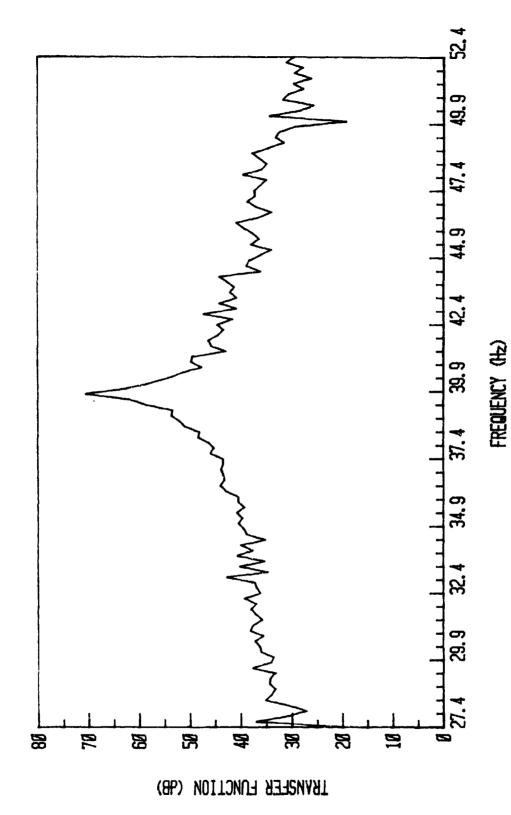
Typical transfer function trace on screen of FFT Analyzer.

Figure 3.19.



Peak of transfer function at one mode using zoom feature on FFT Analyzer.

Figure 3.20.



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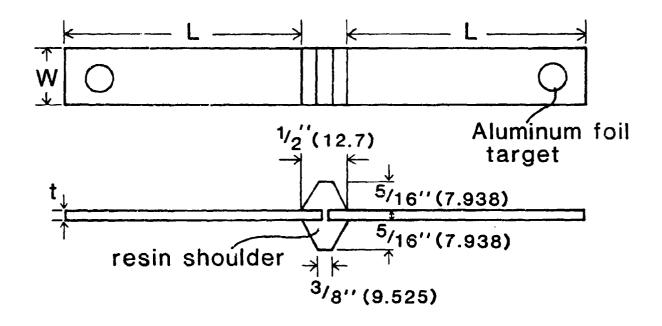


Figure 3.21. Double cantilever beam specimen.

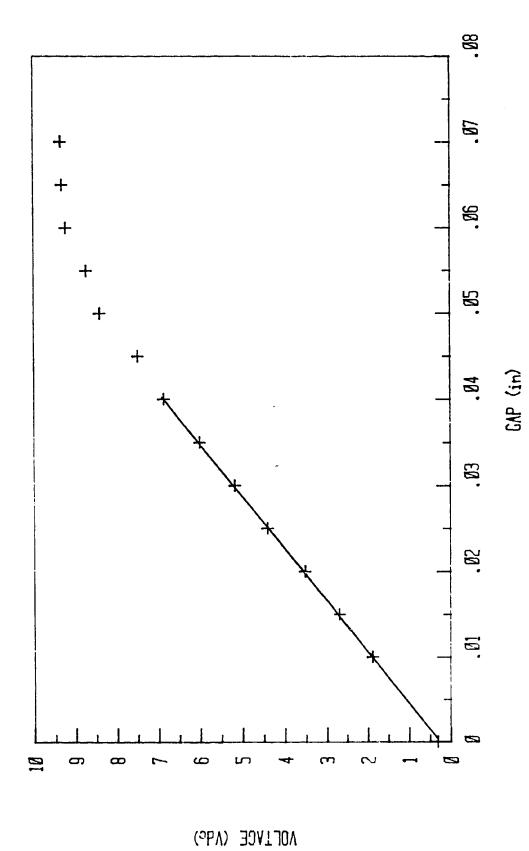


Figure 3.22. Calibration curve for a non contacting probe.

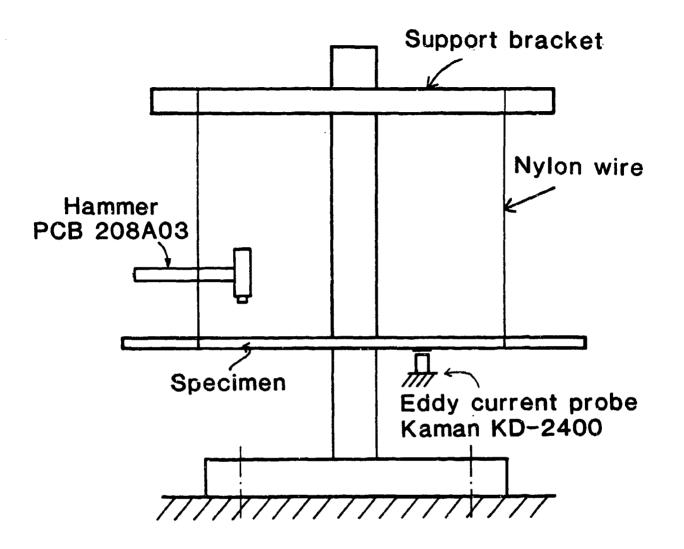


Figure 3.23. Free-free beam apparatus

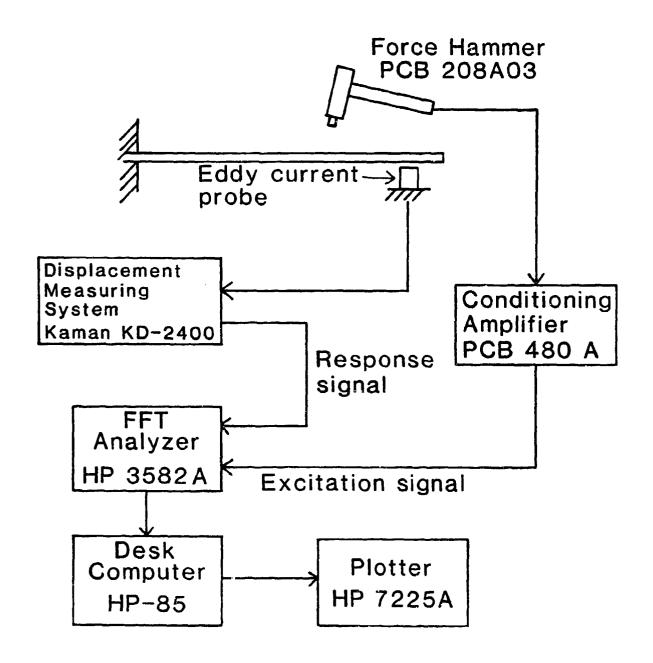
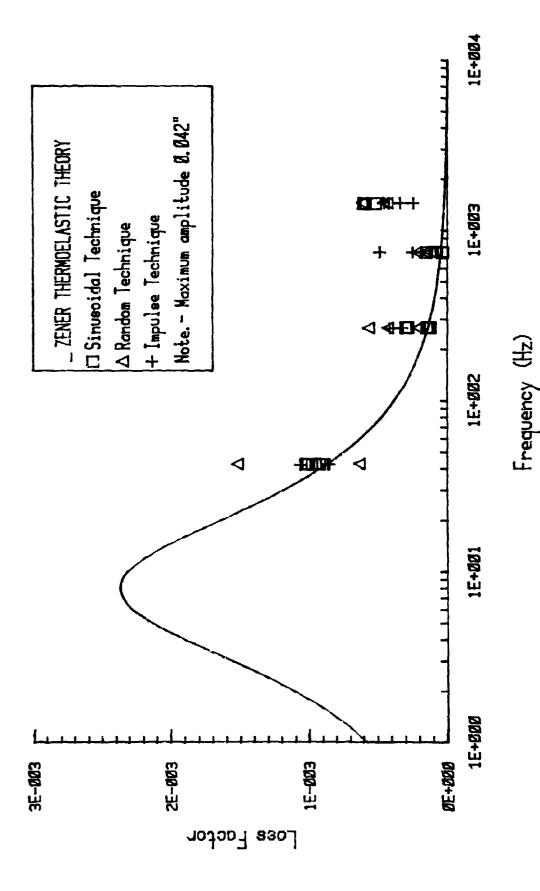


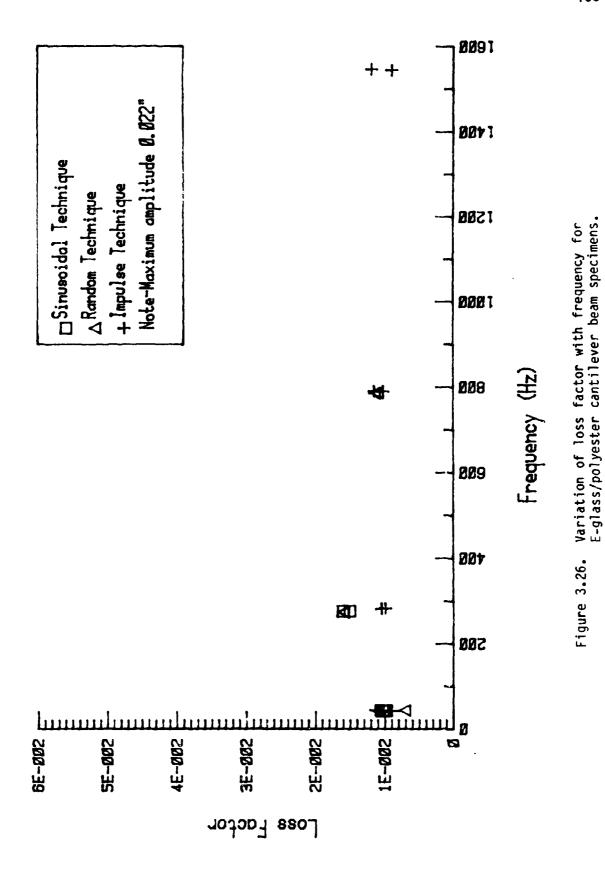
Figure 3.24. Block diagram of instrumentation for impulse technique.



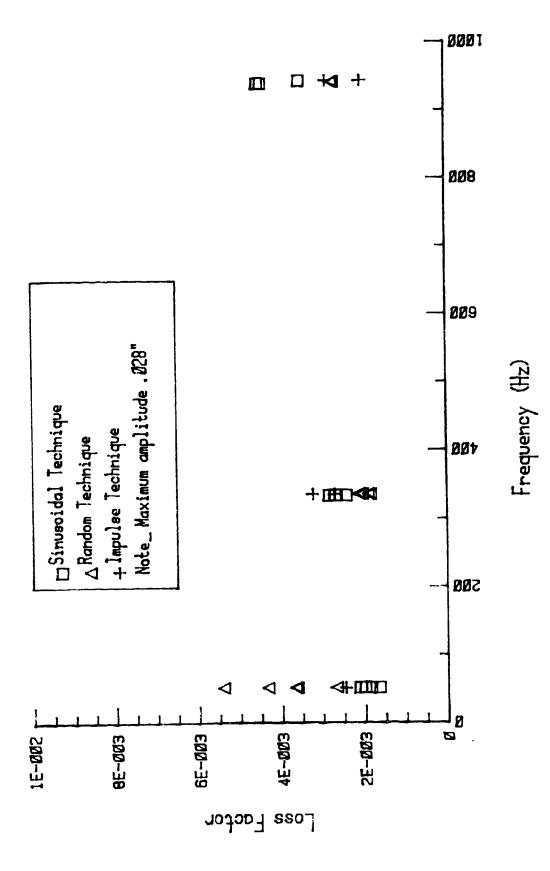
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Figure 3.25. Variation of loss factor with frequency for aluminum cantilever beam specimens.

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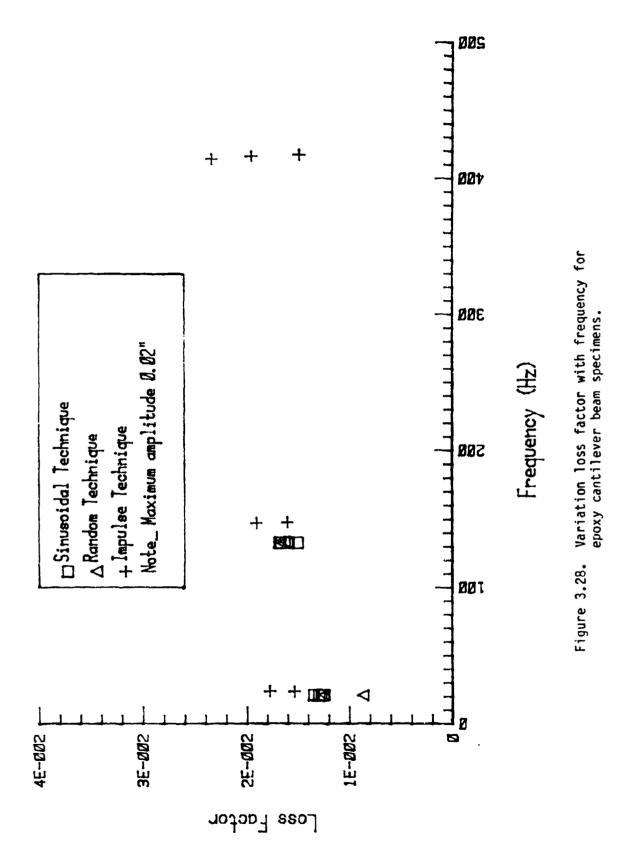


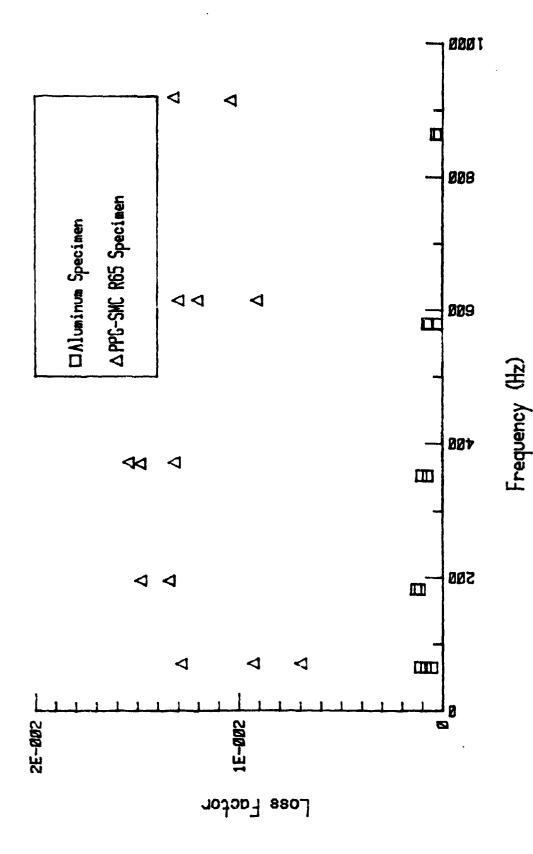
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Figure 3.27. Variation of loss factor with frequency for graphite/epoxy cantilever beam specimens.





Variation of loss factor with frequency for aluminum and E-glass/epoxy polyester free-free beam specimens. Figure 3.29.

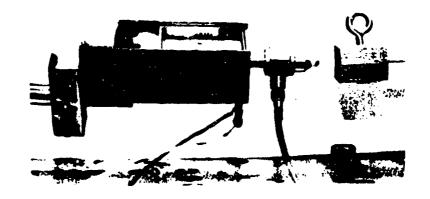
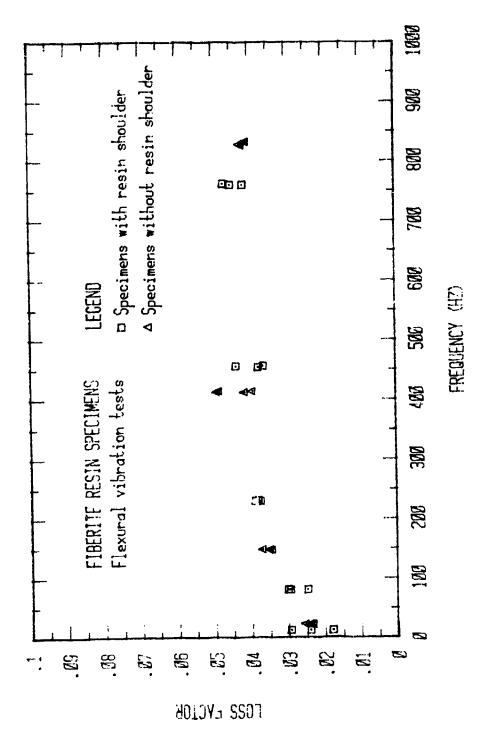
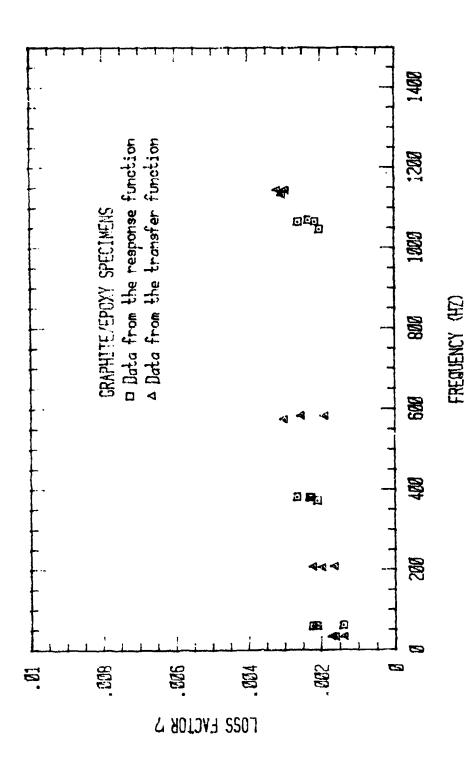


Figure 3.30. Solenoid-type electromagnetic exciter.



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Figure 3.31. Comparison of loss factors using the specimens with and without shoulders.



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Figure 3.32. Comparison of loss factors using the response and the transfer functions.

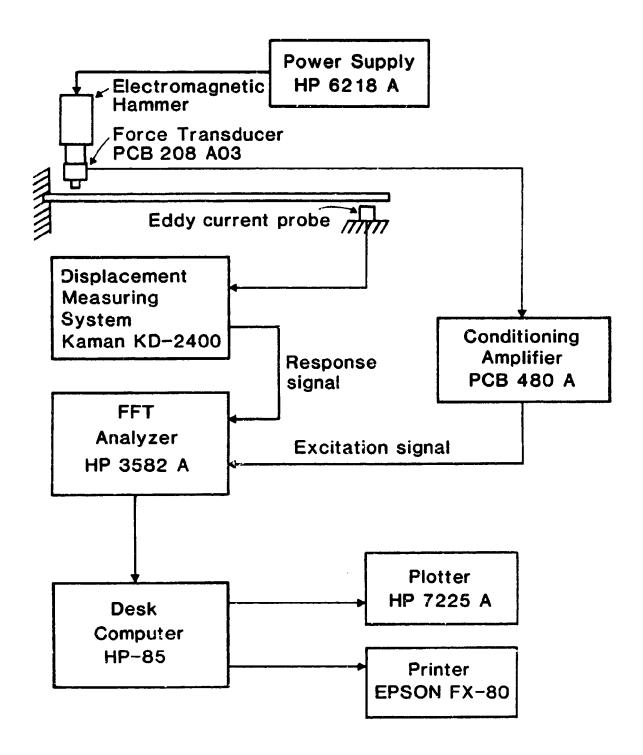


Figure 3.33. Block diagram of instrumentation for flexural vibration tests.

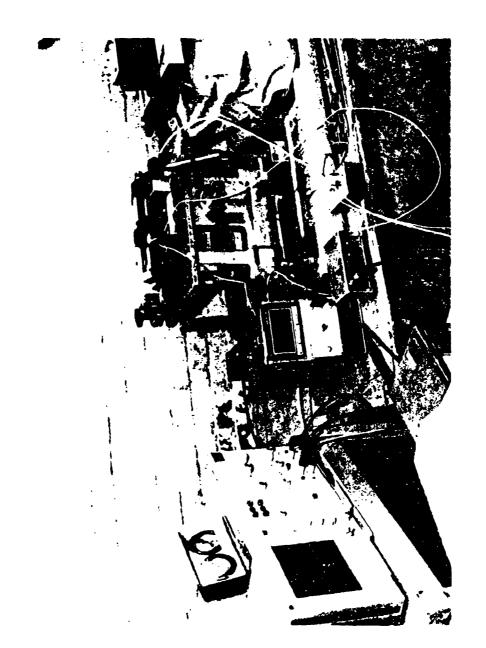


Figure 3.34. Equipment used with the flexural vibration tests.

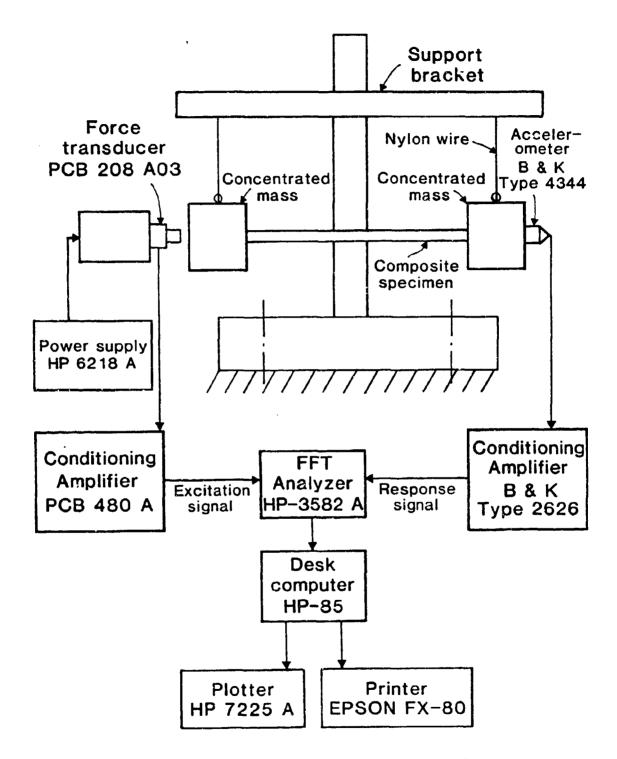


Figure 3.35. Block diagram of instrumentation for extensional vibration tests.

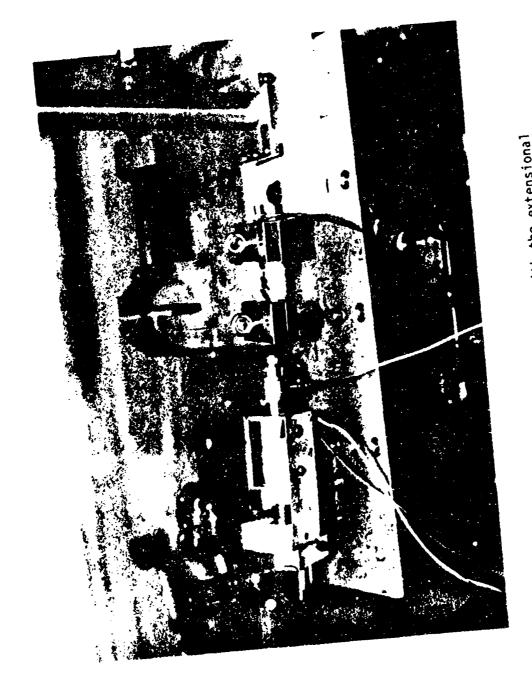


Figure 3.36. Equipment used with the extensional

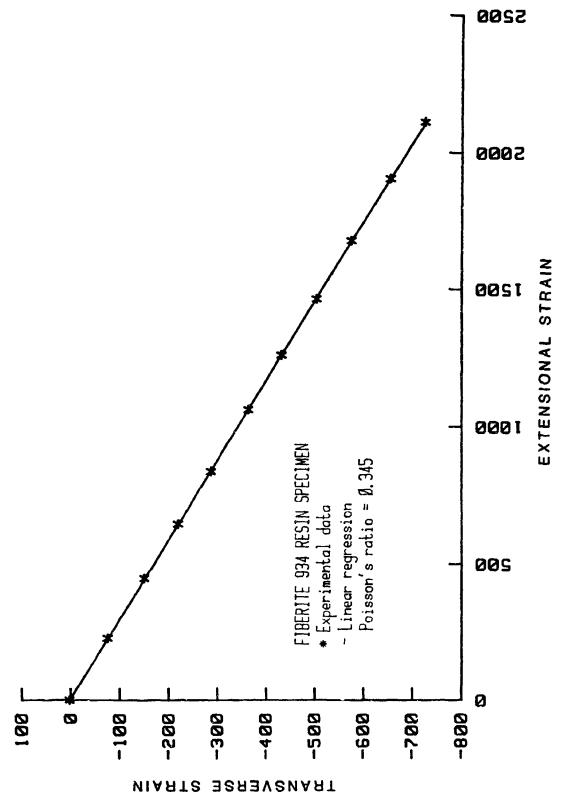
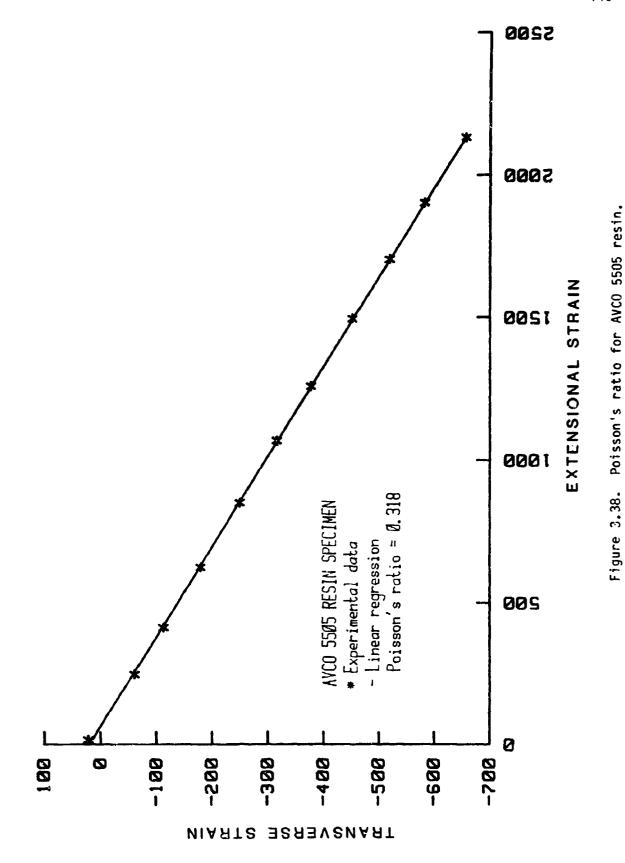
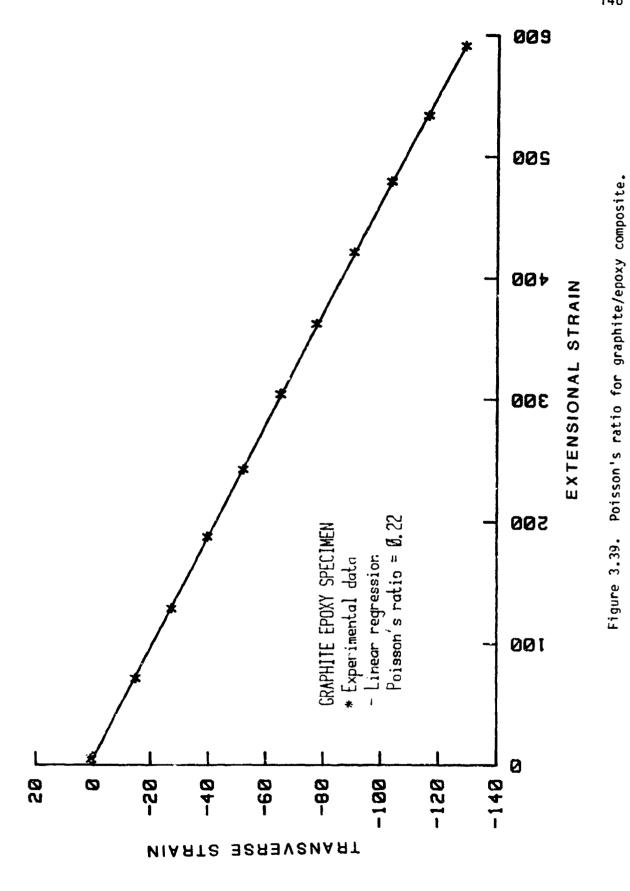
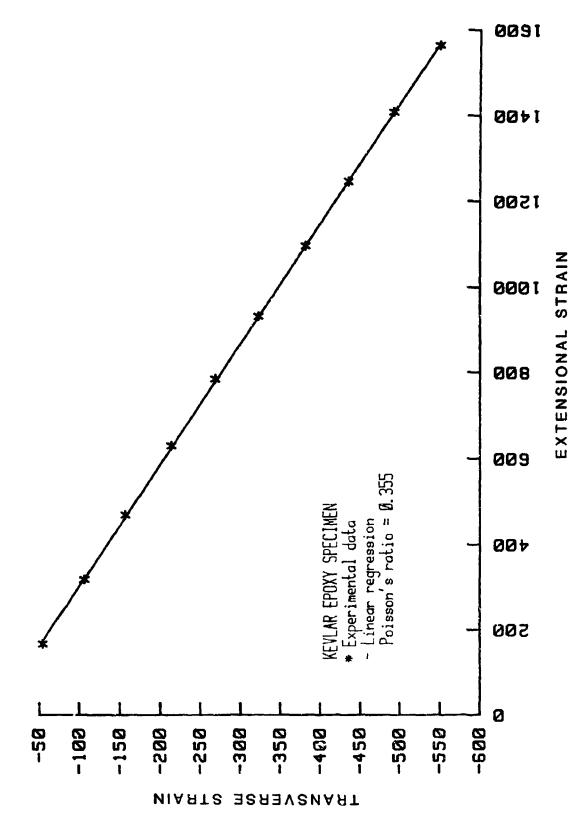


Figure 3.37. Poisson's ratio for Fiberite 934 resin.



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Figure 3.40. Poisson's ratio for Keylar/epoxy composite.

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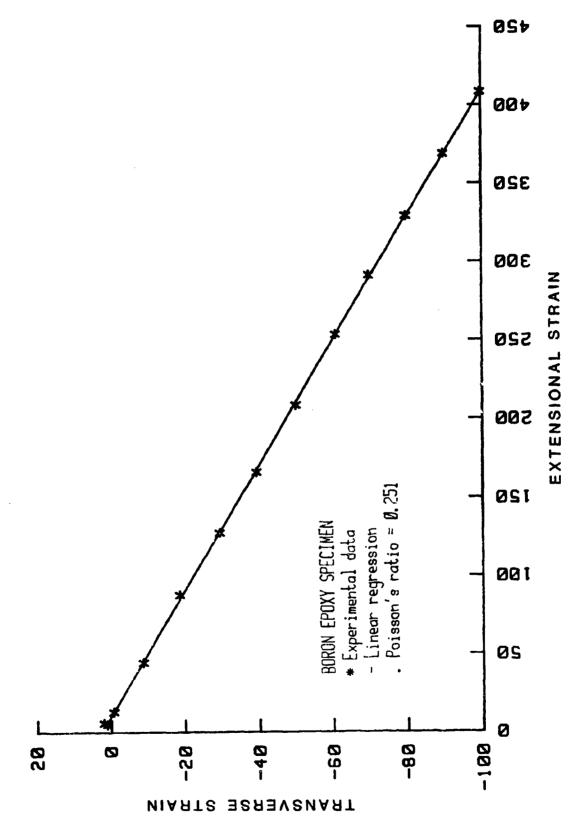
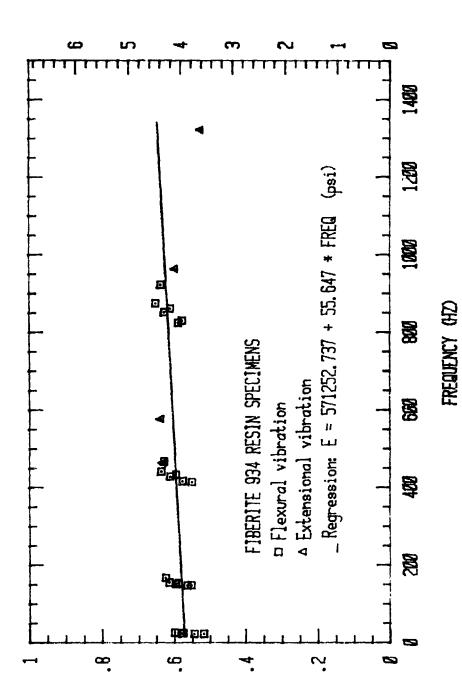
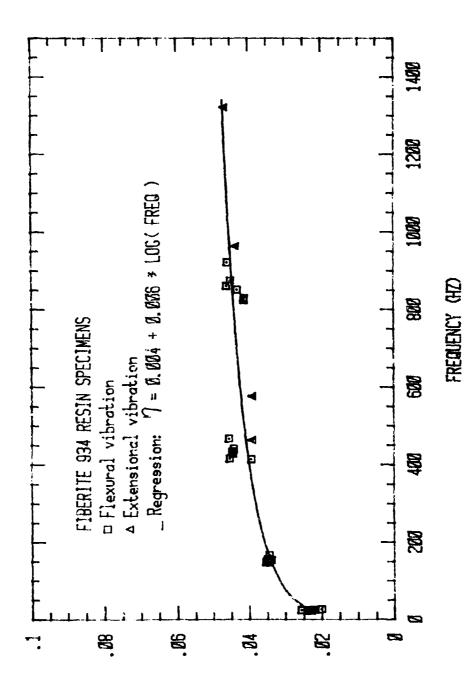


Figure 3.41. Poisson's ratio for boron/epoxy composite.



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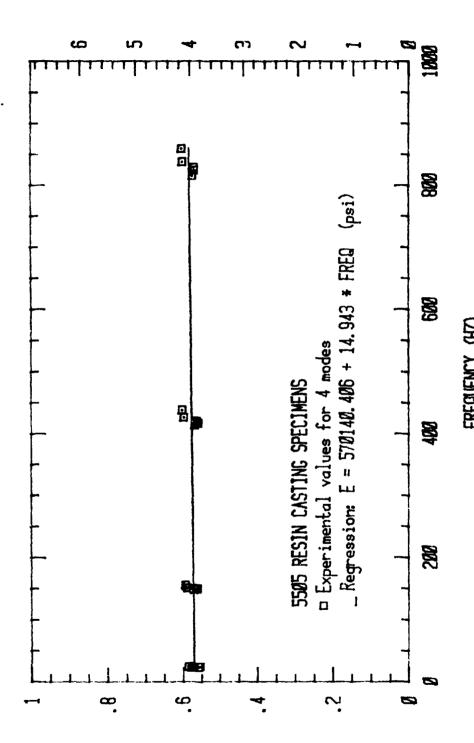
igure 4.1. Variation of storage moduli with frequency for Fiberite 934 resin casting.



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Figure 4.2. Variation of loss factor with frequency for Fiberite 934 resin casting.



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gure 4.3. Variation of storage modulus with frequency for AVCO 5505 resin casting.

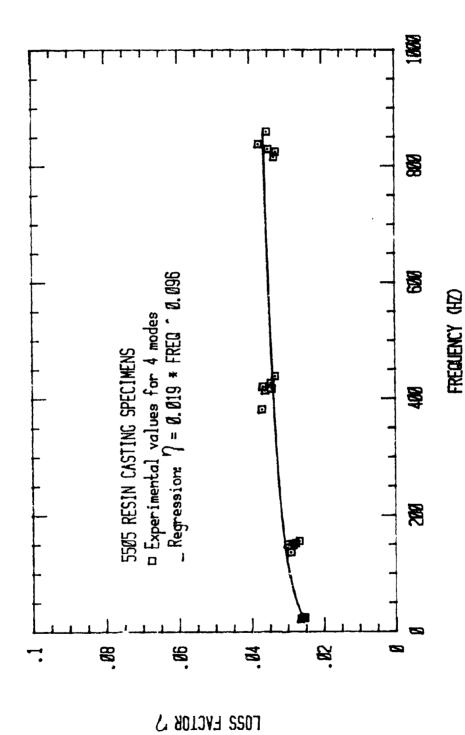


Figure 4.4. Variation of loss factor with frequency for AVCO 5505 resin casting.

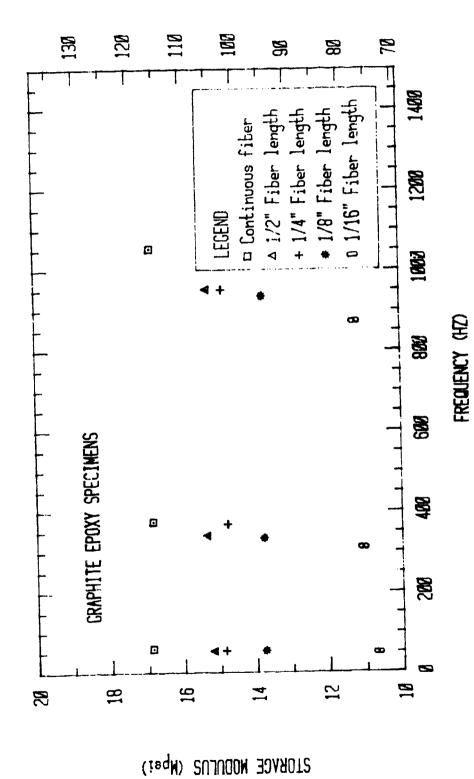
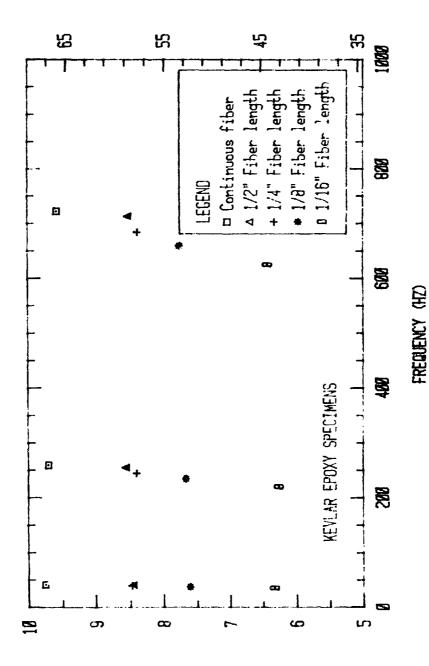


Figure 5.1. Variation of storage modulus with frequency for graphite/epoxy specimens.

Variation of storage modulus with frequency for Kevlar/epoxy specimens.

Figure 5.2.

## STORAGE MODULUS (GPa)



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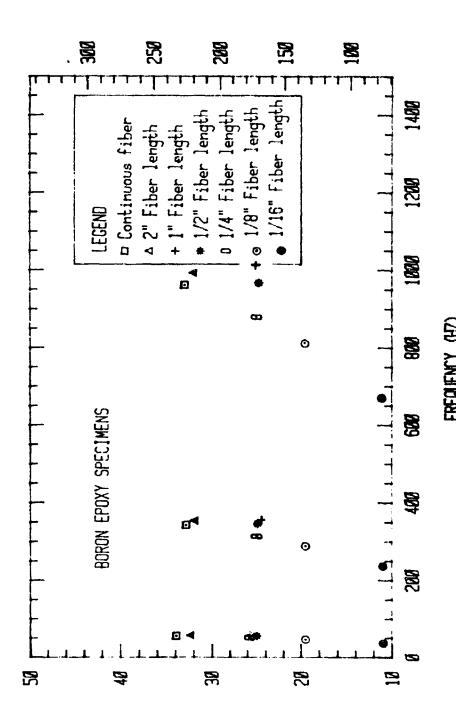


Figure 5.3. Variation of storage modulus with frequency for boron/epoxy specimens.

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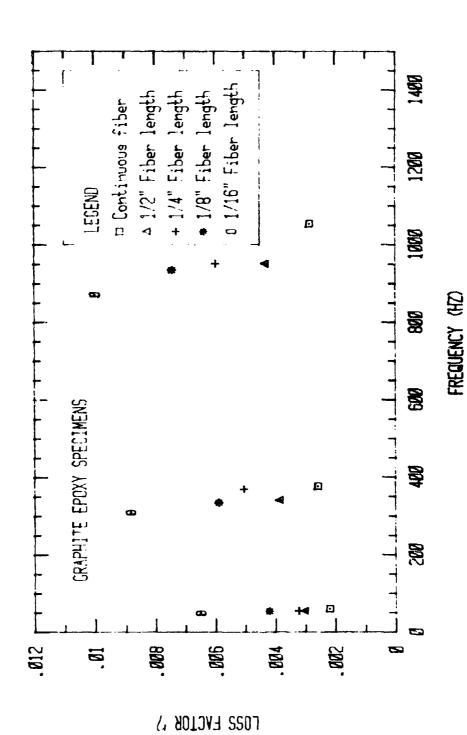


Figure 5.4. Variation of loss factor with frequency for graphite/epoxy specimens.

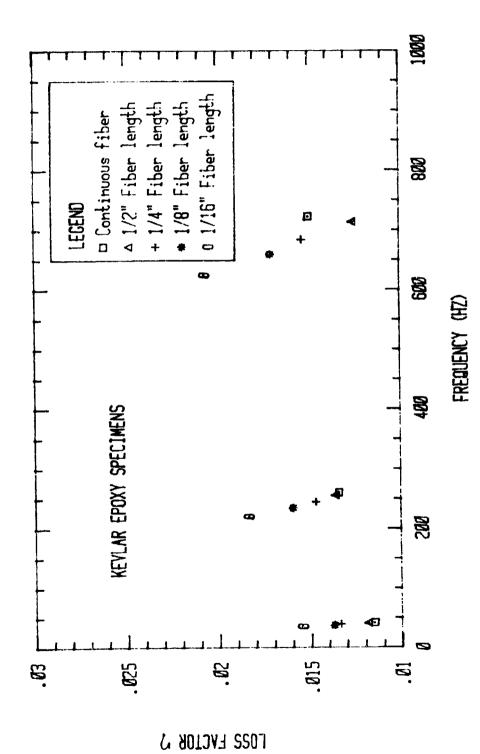


Figure 5.5. Variation of loss factor with frequency for Kevlar/epoxy specimens.

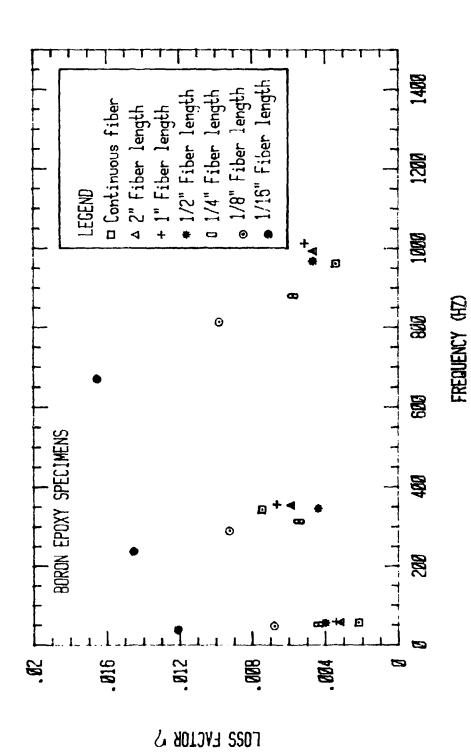
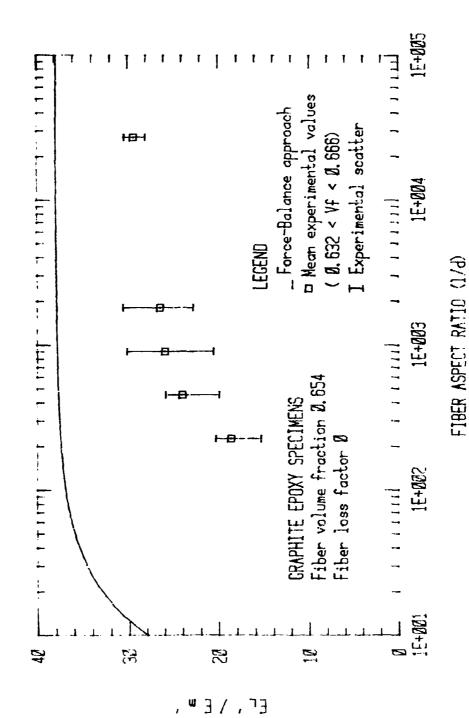


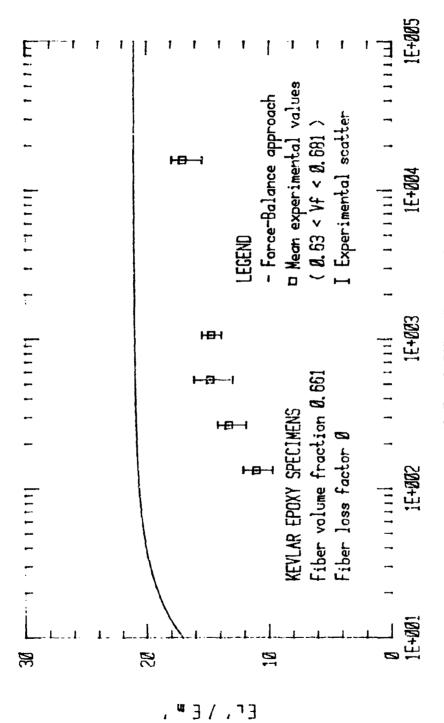
Figure 5.6. Variation of loss factor with frequency for boron/epoxy specimens.



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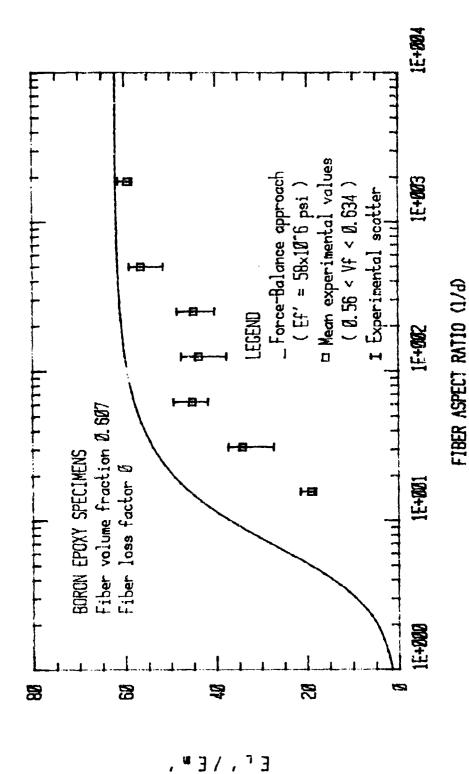
 $E_L^1/E_m^1$  vs. fiber aspect ratic for graphite/epoxy without curve fitting [ $E_f^1$  = 33.10 $^6$  psi (227.38 GPa), Z=1, f=54 Hz]. Figure 5.7.

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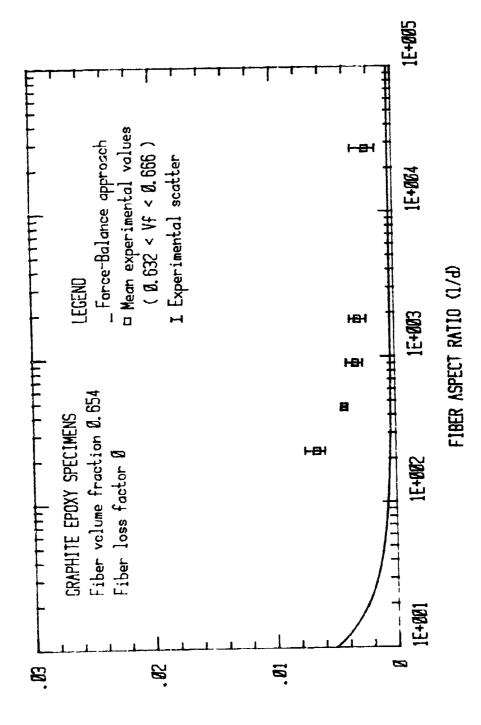
FIBER ASPECT RATIO (1/d)

 $E_{L}^{\prime}/E_{m}^{\prime}$  vs. fiber aspect racio for Kevlar/epoxy without curve fitting [ $E_f^{+}$  = 18 x 10<sup>6</sup> psi (124.02 GPa), Z=1, f=54 Hz]. Figure 5.8.



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 $E_1^{\prime}/E_m^{\prime}$  vs. fiber aspect ratio for boron/epoxy without curve fitting [ $E_f^*$  = 58 x  $10^6$  psi (399.62 GPa), Z=1, f=51.62 Hz]. Figure 5.9.



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Loss factor vs. fiber aspect ratic for graphite/epoxy without curve fitting [Ef =  $33 \times 10^6$  psi (227.38 GPa), Z=1, f=54 Hz]. Figure 5.10.

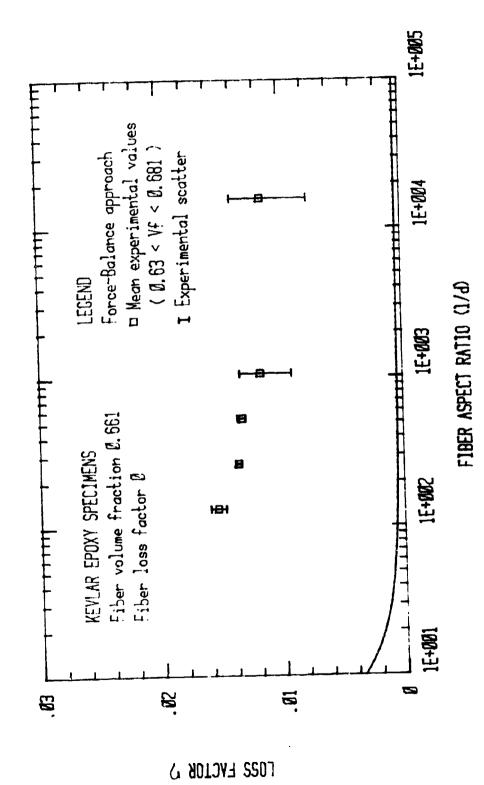
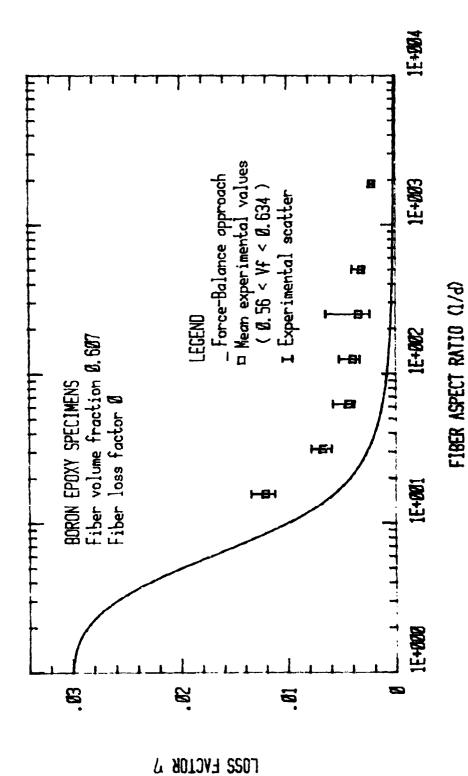
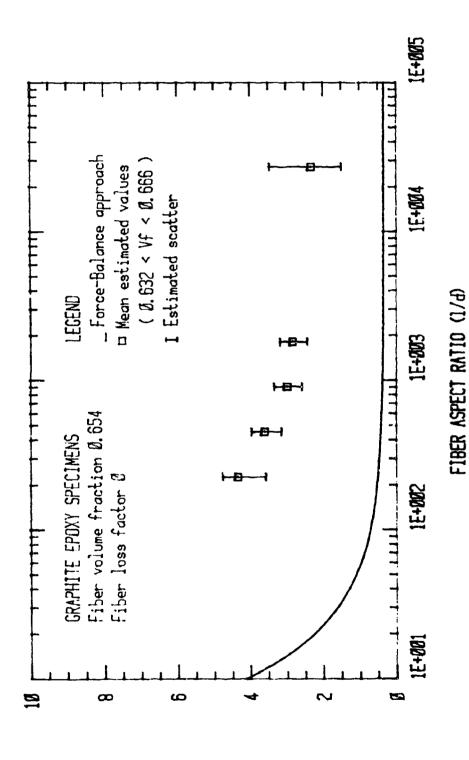


Figure 5.11. Loss factor vs. fiber aspect ratio for Kevlar/epoxy without curve fitting [E $_{
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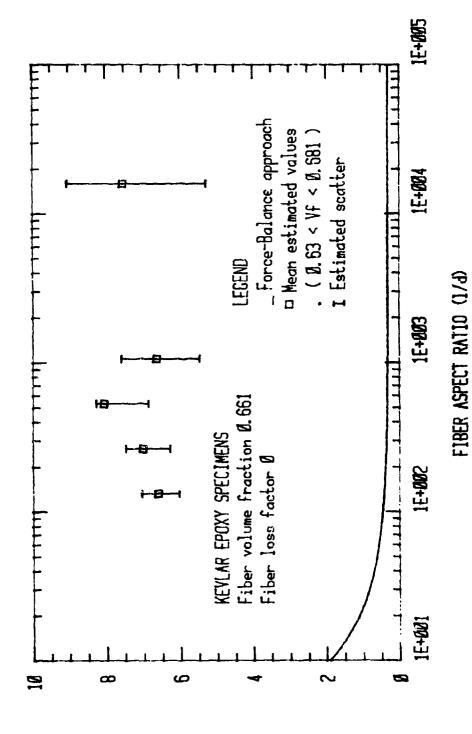
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Figure 5.12. Loss factor vs. fiber aspect ratio for boron/epoxy without curve fitting [E $_{\rm f}$  = 58 x 10 $^{\rm 6}$  psi (399.62 GPa), Z=1, f=51.62 Hz].



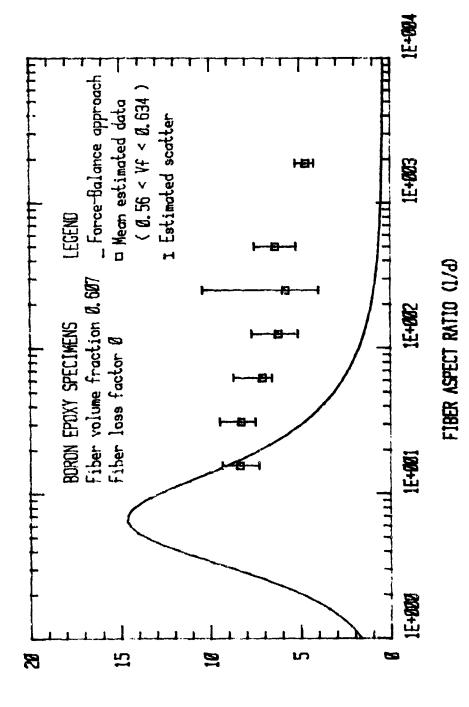
 $E_{\rm L}'/E_{\rm m}'$  vs. fiber aspect ratio for graphite/epoxy without curve fitting  $[E_{\rm f}'=33\times10^6$  psi '. .38 GPa), Z=1, f=54 Hz.] Figure 5.13.

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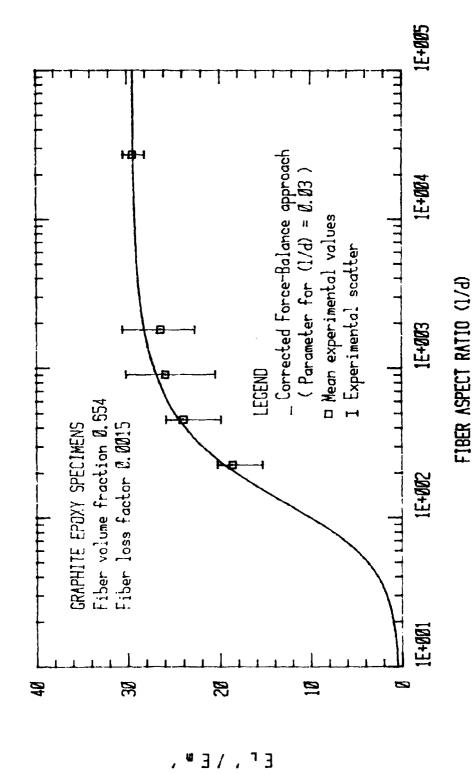


curve fitting  $[E_f^* = 18 \times 10^5 \text{ psi (124.02 GPa), Z=1, f=38.75 Hz]}$ . Figure 5.14.  $E_L^{\prime\prime}/E_m^{\prime\prime}$  vs. fiber aspect ratio for kevlar/epoxy without

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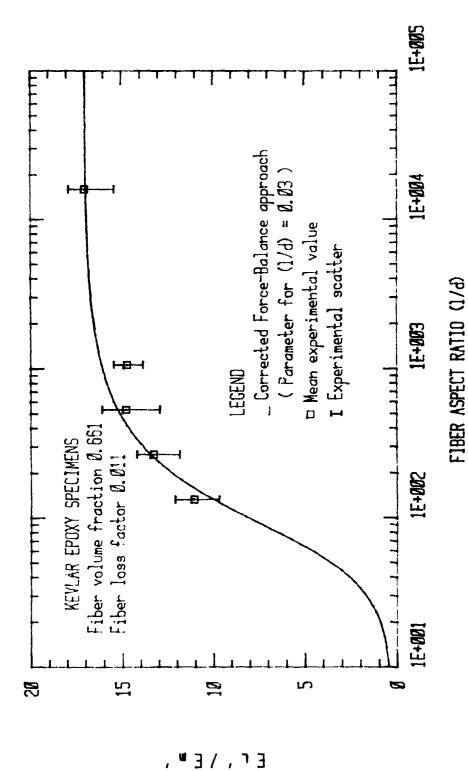


curve fitting [Ef =  $58 \times 10^6$  psi (399.62 GPa), Z=1, f=51.62 Hz]. Figure 5.15.  $E_L^{\prime}/E_m^{\prime}$  vs. fiber aspect ratio for boron/epoxy without



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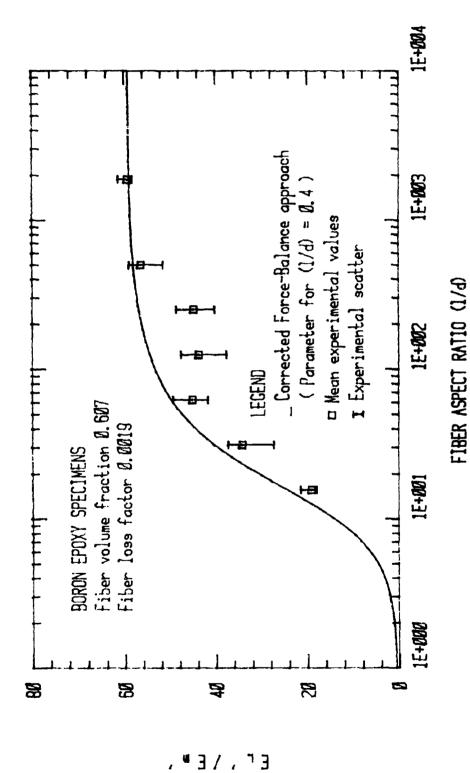
fitting [E $_{\rm f}$  = 25.51 x  $10^6$  psi (175.76 GPa), Z=0.03, f=54 Hz]. Figure 5.16.  $E_L^{\prime}/E_m^{\prime}$  vs. fiber aspect ratio for graphite/epoxy with curve



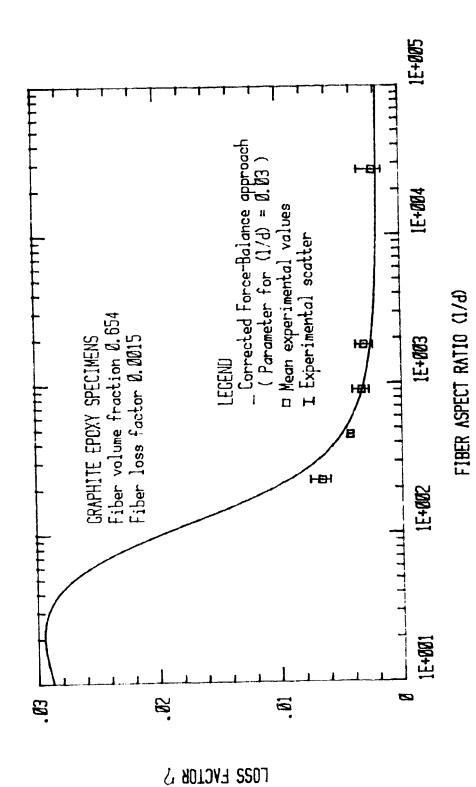
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fitting [ $E_f^{\perp}$  = 14.48 x 10<sup>6</sup> psi (99.767 GPa), Z=0.03, f=38.75 Hz]. Figure 5.17. E'/E" vs. fiber aspect ratio for Keviar/epoxy with curve



fitting [Ef = 55.43 x  $10^6$  psi (381.913 GPa), Z=0.4, f=51.62 Hz]. Figure 5.18.  $E_L^{\prime}/E_m^{\prime}$  vs. fiber aspect ratio for boron/epoxy with curve



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Figure 5.19. Loss factor vs. fiber aspect ratio for graphite/epoxy with curve fitting [Ef = 25.51 x  $10^5$  psi (175.76 GPa), Z=0.03, f=54 Hz].

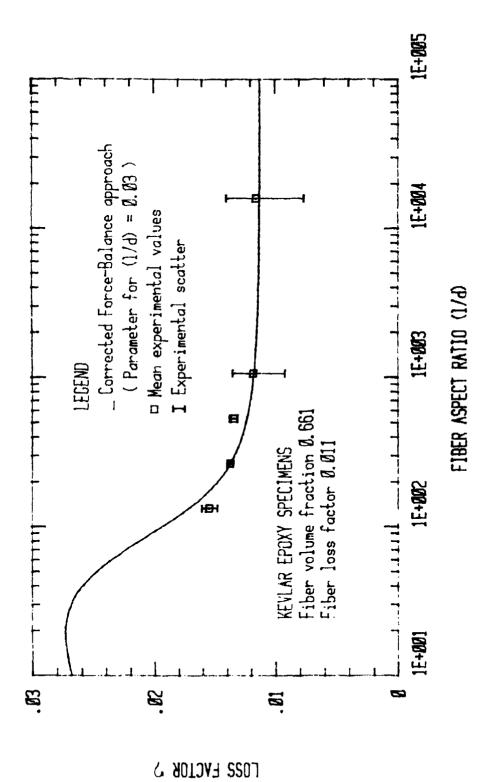
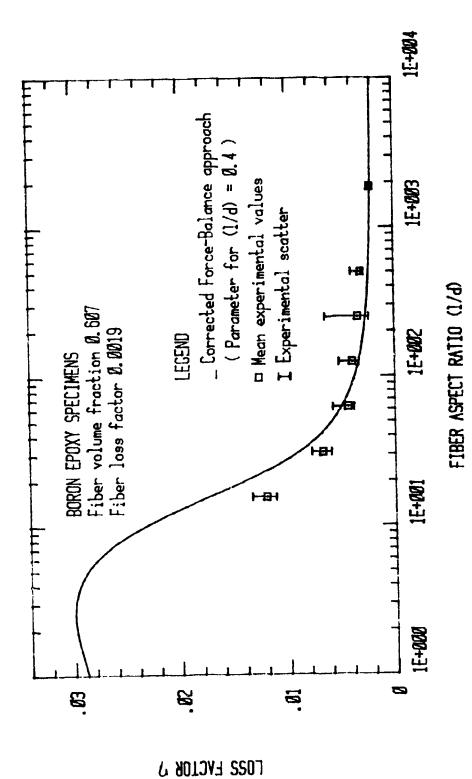
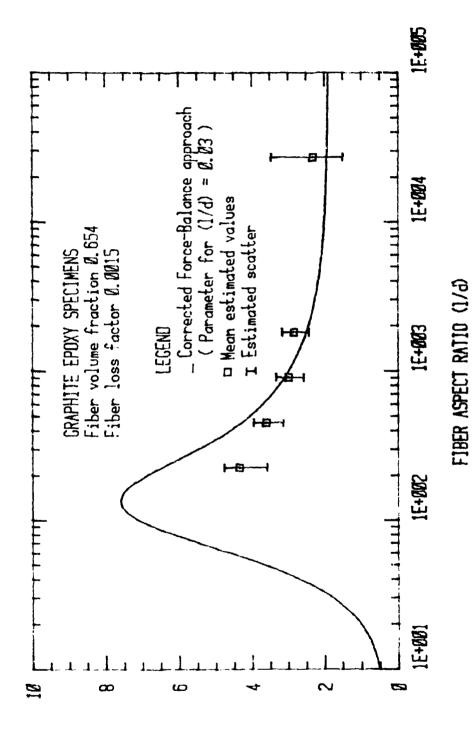


Figure 5.20. Loss factor vs. fiber aspect ratio for Kevlar/epoxy with curve fitting [E $_{\rm f}^{\prime}$  = 14.48 x  $10^{6}$  psi (99.767 GPa), Z=0.03, f=38.75 Hz].



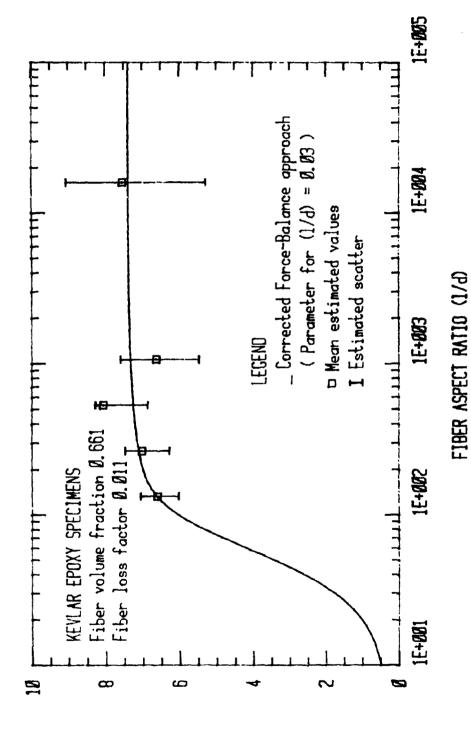
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Figure 5.21. Loss factor vs. fiber aspect ratio for boron/epoxy with curve fitting [E $_{\rm f}^*$  = 55.43 x  $_10^6$  psi (381.913 GPa), Z=0.4, f=51.62 Hz].



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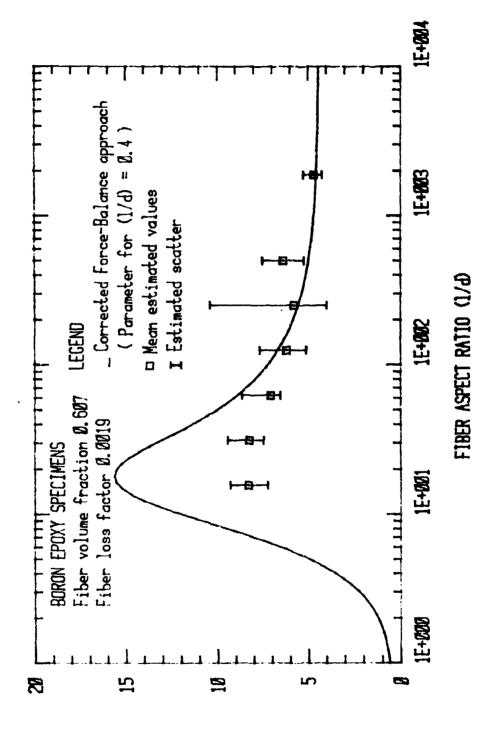
fitting [E $_1^{\prime}$  = 25.51 x 10<sup>6</sup> psi (175.76 GPa), Z=0.03, f=54 Hz]. Figure 5.22. E"/E" vs. fiber aspect ratio for graphite/epoxy with curve



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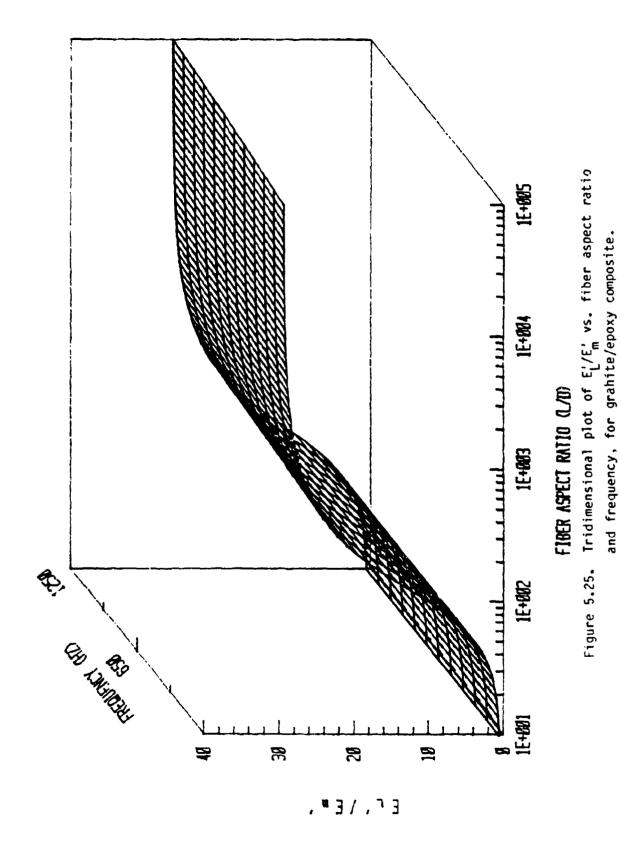
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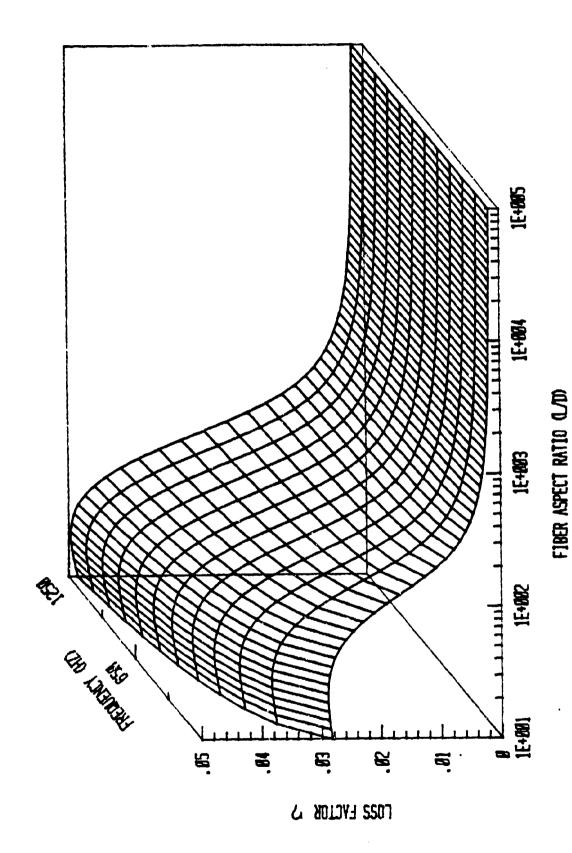
fitting [E' $_{\rm f}$  = 14.48 x 10<sup>6</sup> psi (99.767 GPa), Z=0.03, f=38.75 Hz]. Figure 5.23.  $E_L^{\prime}/E_m^{\prime}$  vs. fiber aspect ratio for Kevlar/epoxy with curve



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fitting [ $E_f^{\prime}$  = 55.43 x  $10^5$  psi (381.913 GPa), Z=0.4, f=51.62 Hz]. Figure 5.24.  $E_{\rm L}^{\prime}/E_{\rm m}^{\prime}$  vs. fiber aspect ratio for boron/epoxy with curve

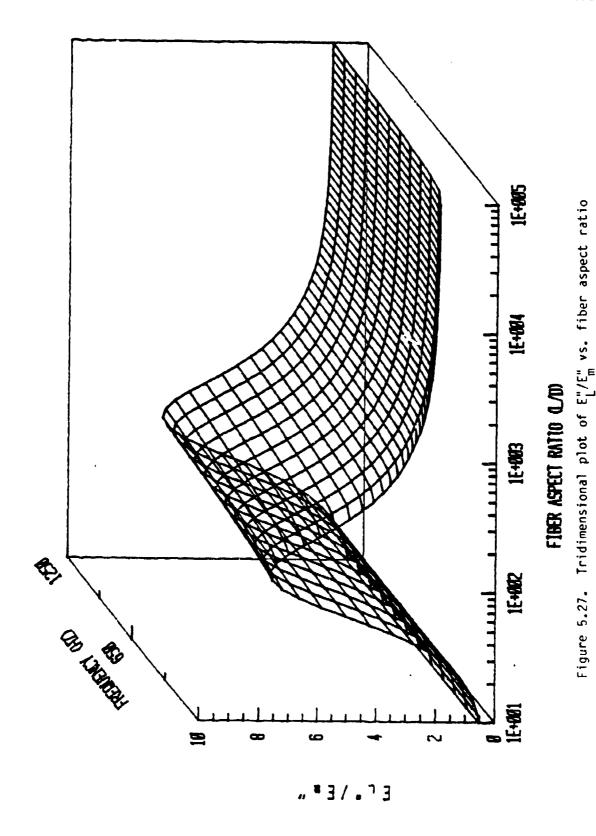




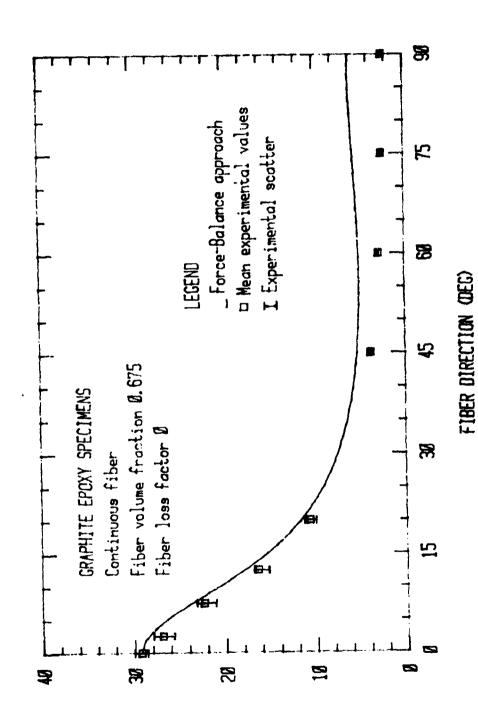
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Figure 5.26. Tridimensional plot of loss factor vs. fiber aspect ratio and frequency, for graphite/epoxy composite.

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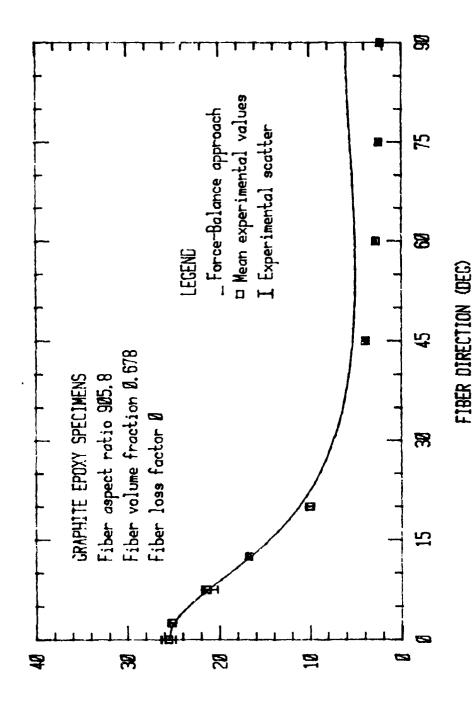


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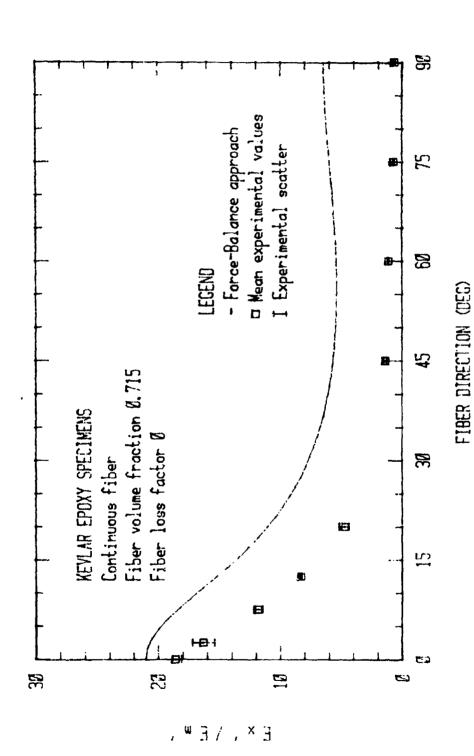
 $E_{\rm X}'/E_{\rm M}'$  vs. fiber direction for continuous graphite/epoxy, without curve fitting [isotropic fiber,  $E_{\rm f}'=33\times10^{\circ}$  psi (227.38 GPa),  $\varepsilon_1$ =2,  $\varepsilon_2$ =2, Z=1, f=1152.5 Hz]. Figure 6.1.



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 $E_x'/E_m'$  vs. fiber direction for discontinuous graphite/epoxy, without curve fitting [isotropic fiber,  $E_f'=33\times10^5$  psi (227.38 GPa),  $E_1=2$ ,  $E_2=1$ , Z=1, f=1206.15 Hz]. Figure 6.2.

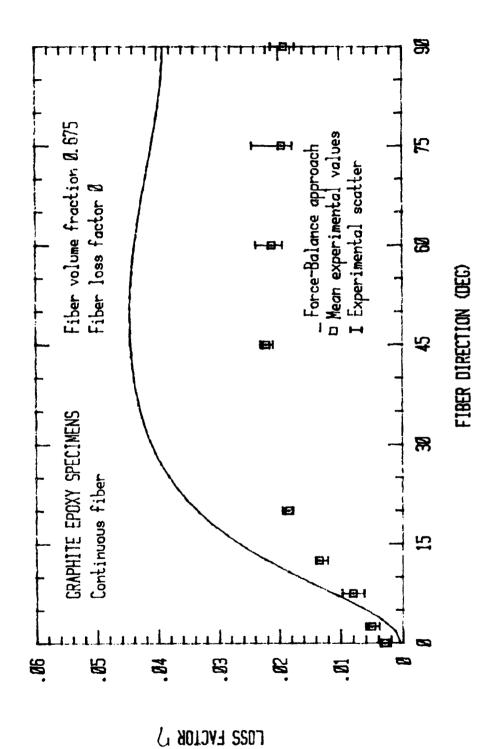


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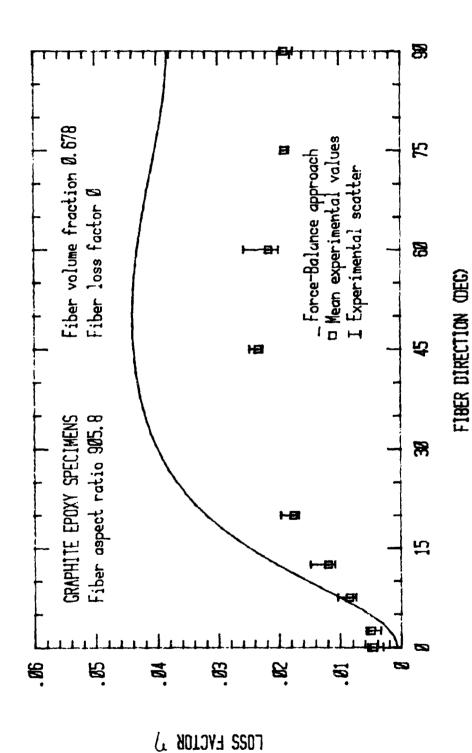
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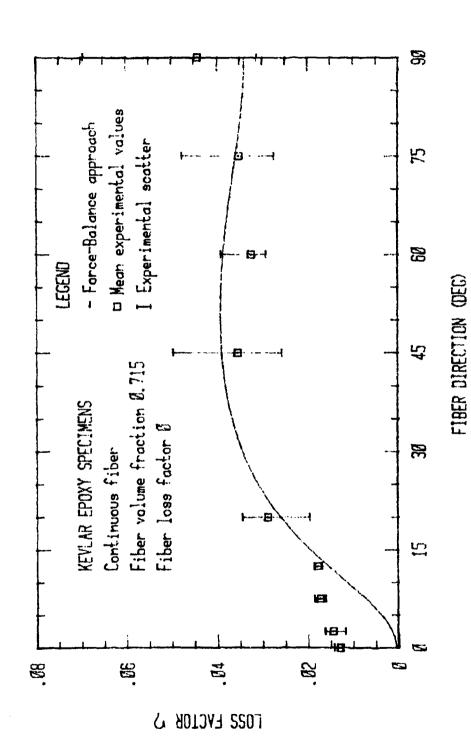


Loss factor vs. fiber direction for continuous graphite/epoxy, without curve fitting [isotropic fiber,  $E_f=33\times10^6$  psi (227.38 GPa),  $\varepsilon_1$ =2,  $\varepsilon_2$ =1, Z=1, f=1152.5 Hz]. Figure 6.4.



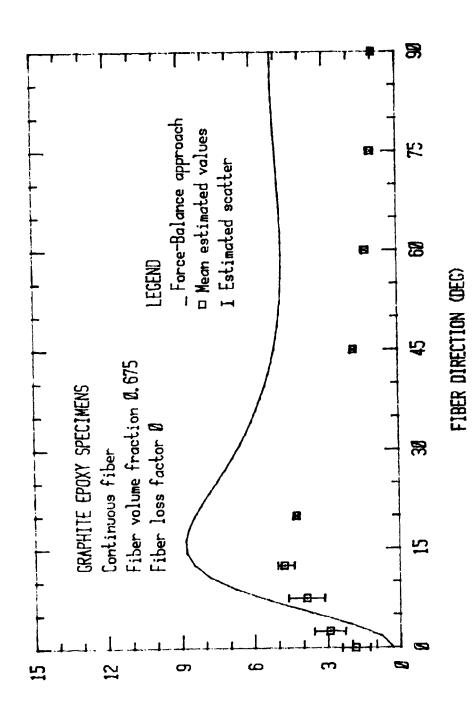
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Loss factor vs. fiber direction for discontinuous graphite/epoxy, without curve fitting [isotropic fiber,  $E_{\rm f}$  = 33 x  $10^5$  psi (227.38 GPa),  $\varepsilon_1$ =2,  $\xi_2$ =1, Z=1, f=1206.15 Hz]. Figure 6.5.



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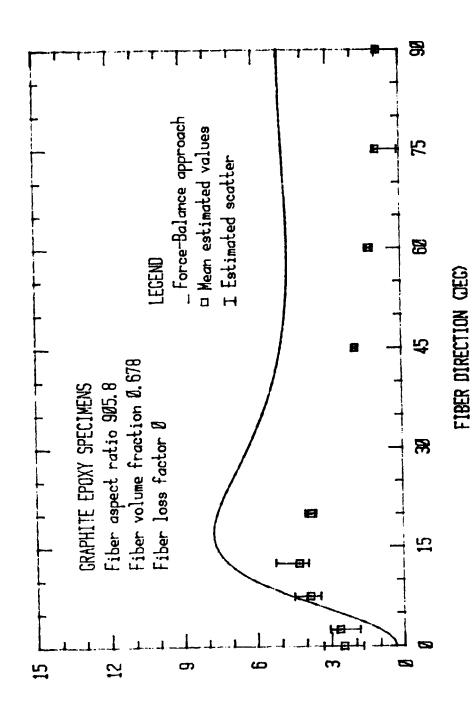
Loss factor vs. fiber direction for continuous Kevlar/epoxy, without curve fitting [isotropic fiber,  $E_f$  = 18 x 10 $^6$  psi (124.02 GPa),  $\varepsilon_1$ =2,  $\varepsilon_2$ =1, Z=1, f=895.1 Hz]. Figure 6.6.



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["sotropic fiber,  $E_f = 33 \times 10^6$  psi (227.38 GPa),  $\varepsilon_1$ =2,  $\varepsilon_2$ =1, Z=1, f=1152.5 Hz].  $E_{x}^{*}/E_{m}^{*}$  vs. fiber direction for continuous graphite/epoxy, without curve fitting Figure 6.7.

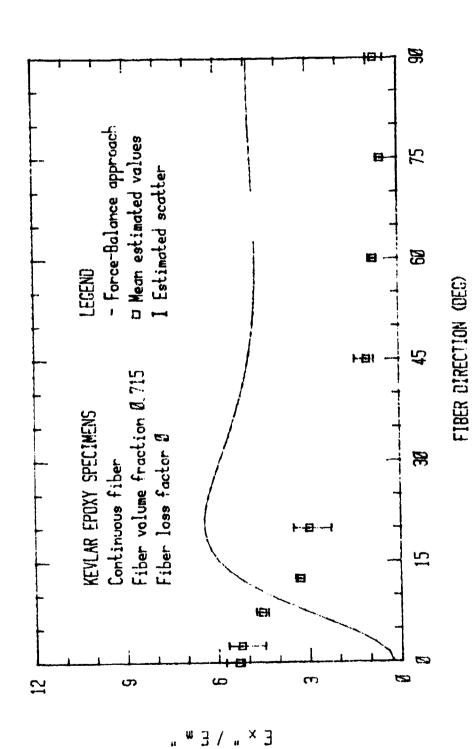


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 $E_{x}^{*}/E_{m}^{*}$  vs. fiber direction for discontinuous graphite/epoxy, without curve fitting [isotropic fiber,  $E_f^{+} = 33 \times 10^5$  psi (227.38 SPa),  $\xi_1^{-}$ 2,  $\xi_2^{-}$ 1, Z=1, f=1206.15 Hz]. Figure 6.8.



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[isotropic fiber,  $E_{\rm f}^{\rm f}$  = 18 x  $10^6$  psi (124.02 GPa),  $\varepsilon_1$ =2,  $\varepsilon_2$ =1, Z=1, f=895.1 Hz].  $E_{x}^{*}/E_{m}^{*}$  vs. fiber direction for continuous Kevlar/epoxy, without curve fitting Figure 6.9.

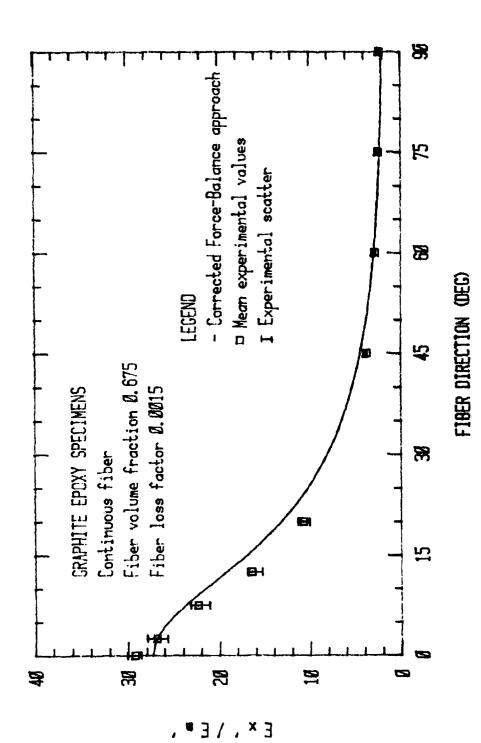
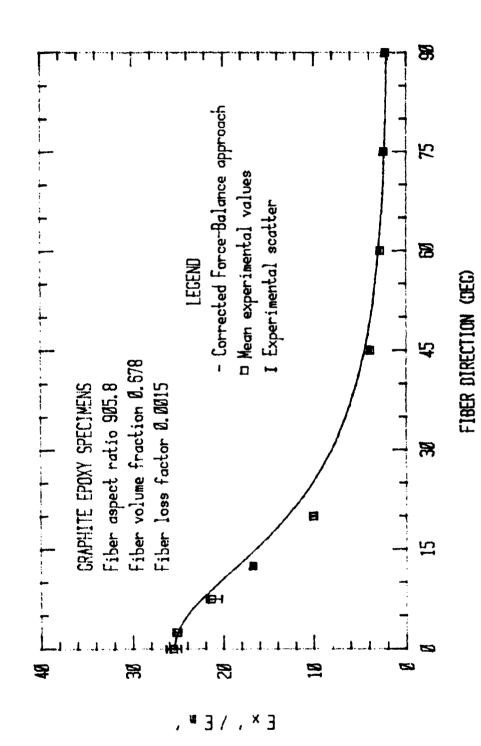
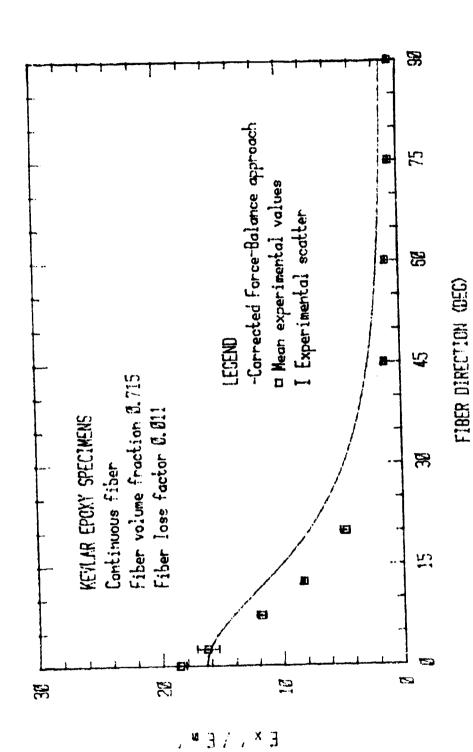


Figure 6.10,  $\mathbf{E}_{\mathbf{x}}/\mathbf{E}_{\mathbf{x}}$  vs. fiber direction for continuous graphite/epoxy, with curve fitting  $\{\mathbf{E}_{\mathbf{f}_{\mathbf{L}}}^{\dagger} = 25.51 \times 10^6 \text{ psi}$  (175.76 GPa),  $\mathbf{E}_{\mathbf{f}_{\mathbf{L}}}^{\dagger} = 2 \times 10^6 \text{ psi}$  (13.78 GPa),  $\mathbf{G}_{\mathbf{L}_{\mathbf{L}}}^{\dagger} = 4 \times 10^6 \text{ psi}$  (27.56 GPa),  $\mathbf{E}_{\mathbf{f}_{\mathbf{L}}}^{\dagger} = 2 \times 10^6 \text{ psi}$  (13.78 GPa),  $\mathbf{G}_{\mathbf{L}_{\mathbf{L}}}^{\dagger} = 4 \times 10^6 \text{ psi}$  (27.56 GPa),  $\mathbf{E}_{\mathbf{f}_{\mathbf{L}}}^{\dagger} = 2 \times 10^6 \text{ psi}$  (13.78 GPa),  $\mathbf{G}_{\mathbf{L}_{\mathbf{L}}}^{\dagger} = 4 \times 10^6 \text{ psi}$  (27.56 GPa),  $\mathbf{E}_{\mathbf{f}_{\mathbf{L}}}^{\dagger} = 2 \times 10^6 \text{ psi}$ 



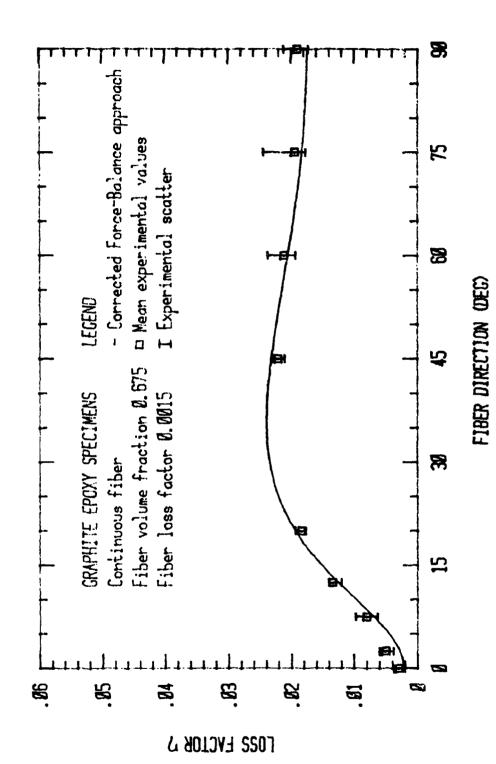
(175.76 GPa), Eir = 2 x 106 psi (13.78 GPa), G'flf 4 x 106 psi (27.56 GPa), E f2, E 34, Z=0.03, f=1206.15 Hz]. Figure 6.11.  $E_x/E_y$  vs. fiber direction for discontinuous graphite/epoxy with curve fitting  $[E_{fL}^*$  = 25.51 x  $10^6$  psi



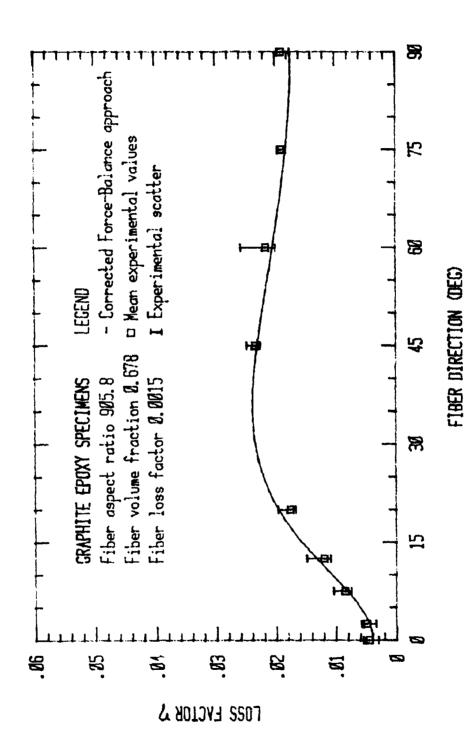
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[99.767 GPa]. Eir = 1 x 106 psi (6.89 GPa). Gilt = 2 x 106 psi (13.78 GPa), Ei=2. Er=1, Z-0.03, f=895.1 Hz]. Figure 6.12.  $E_{\chi}/E_{m}$  vs. fiber direction for continuous Kevlar/epoxy, with curve fitting  $[E_{\rm f}]_{\star}$  = 14.48 ×  $10^{6}$  psi

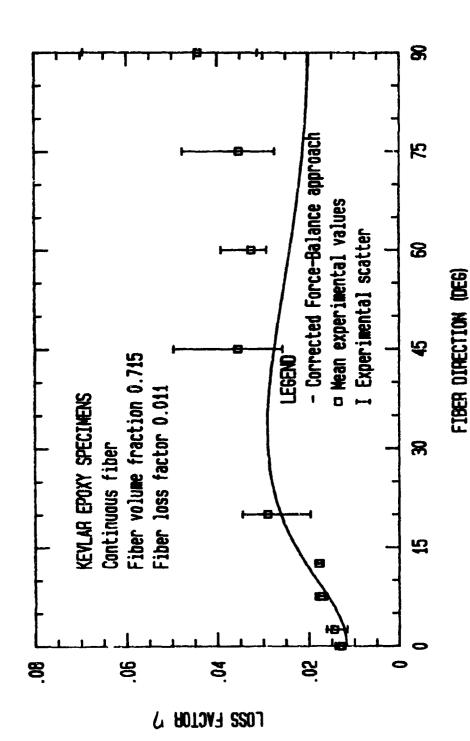
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(175.76 GPa).  $E_{ff} = 2 \times 10^6 \text{ psi}$  (13.78 GPa).  $G_{fLF} = 4 \times 10^6 \text{ psi}$  (27.56 GPa),  $E_{f} = 2$ ,  $E_{g} = 4$ , 2=0.03, f=1152.5 Hz]. Figure 6.13, loss factor vs. fiber direction for continuous graphite/epoxy, with curve fitting  $\{E_{ij}^{\dagger}=25.51 \times 10^6 \text{ pst}$ 



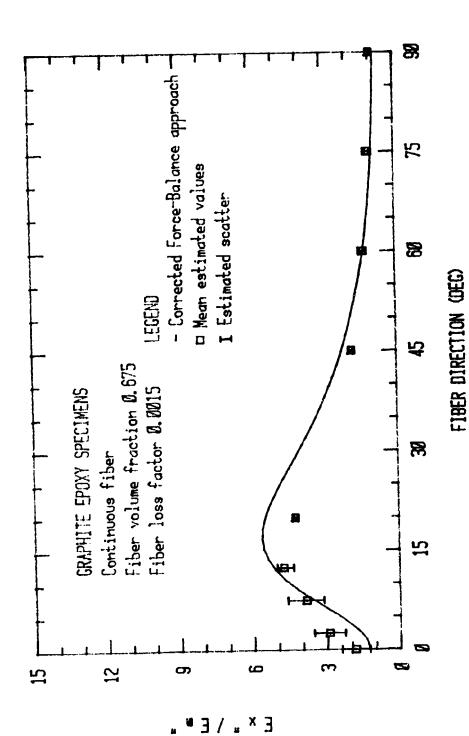
(175.76 GPa), E'T = 2 x 106 psi (13.78 GPa), Gilt = 4 x 106 psi (27.56 GPa), c1 2, c2 = 4, 2x0.03, f=1206.15 Hz]. Figure 6.14. Loss factor vs. fiber direction for discontinuous graphite/epoxy, with curve fitting  $[\mathsf{E}_{\mathsf{f}_1}^i]$ 



(99.767 GPa),  $E_{ff}^*=1 \times 10^6$  psi (6.89 GPa),  $G_{fLT}^*=2 \times 10^6$  psi (13.78 GPa),  $\epsilon_1$ =2,  $\epsilon_2$ =1, 2=0.03, f=895.1 Hz]. Figure 6.15. Loss factor vs. fiber direction for continuous Kevlar/epoxy, with curve fitting [Eft = 14.48 x 10<sup>6</sup> psi

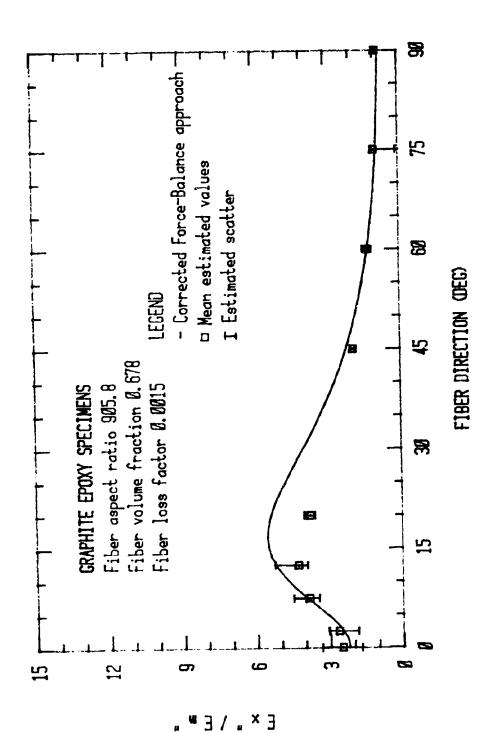
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Figure 6.16.  $E_x^\prime/E_x^{\mu}$  vs. fiber direction for continuous graphite/epoxy, with curve fitting  $f_{eff}^{\mu}$  = 25.51 x  $10^6$  psi (99.767 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa),  $G_{eff}^{\mu}$  = 2 x  $10^6$  psi (13.78 GPa),  $E_{eff}^{\mu}$  = 1 x  $10^6$  psi (6.89 GPa).



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(175.76 GPa), E'T = 2 x 106 psi (13.78 GFa), G'LT = 4 x 106 psi (27.56 GPa), E<sub>1</sub>=2, E<sub>2</sub>=4, Z=0.03, f=1206.15 Hz]. Figure 6.17.  $E_x/E_m$  vs. fiber direction for discontinuous graphite/epoxy with curve fitting  $[E_f]_{\rm r}=25.51 \times 10^6$  psi

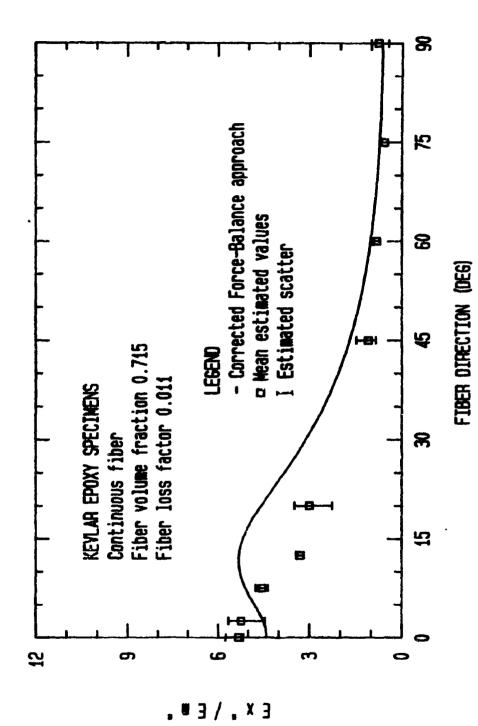


Figure 6.18.  $E_x'/E_m'$  vs. fiber direction for continuous Kevlar/epoxy, with curve fitting [E' = 14.48 x 10<sup>6</sup> psi (99.767 GPa),  $E_{iT}^* = 1 \times 10^6$  psi (6.89 GPa),  $G_{iLT}^* = 2 \times 10^6$ ,  $E_1^* = 2 \times 10^6$ ,  $E_2^* = 1$ , Z=0.03, f=895.1 Hz].

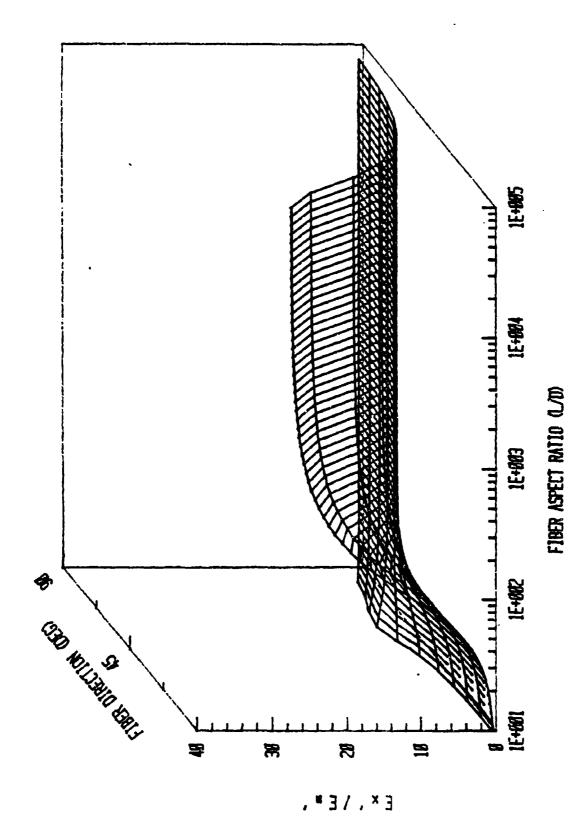
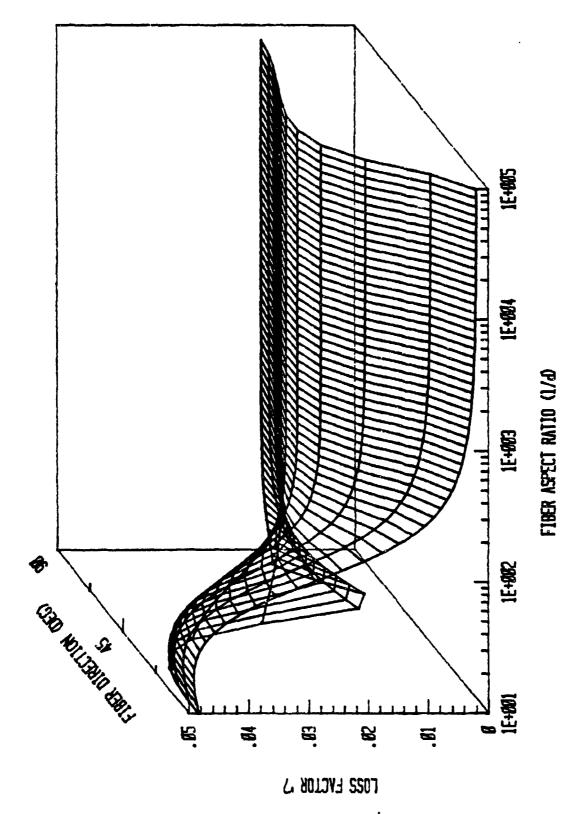
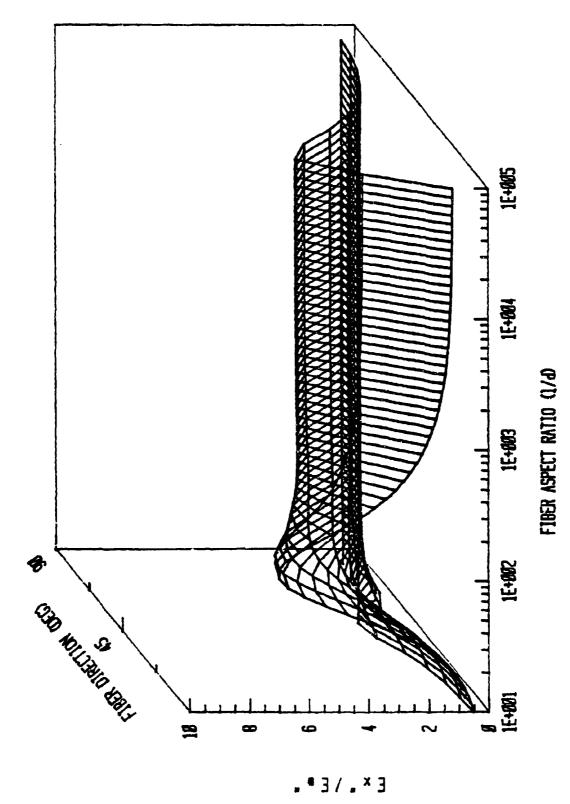


Figure 6.19. Tridimensional plot of  $E_{\rm X}^{\prime}/E_{\rm m}^{\prime}$  vs. fiber aspect ratio and fiber direction, for graphite/epoxy composite.



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Tridimensional plot of loss factor vs. fiber aspect ratio and fiber direction, for graphite/epoxy composite. Figure 6.20.



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Tridimensional plot of  $E_{x}^{\prime}/E_{m}^{\prime}$  vs. fiber aspect ratio and fiber direction, for graphite/epoxy composite. Figure **6.**21.